

Lecture 18

Goldstinos and Gravitinos

genera theory with W and gauge

$$\Psi = (\lambda^a, \psi_i)$$

$$M_{\text{fermion}} = \begin{pmatrix} 0 & \sqrt{2} g_a \langle \phi^* \rangle T^a{}^i \\ \sqrt{2} g_a \langle \phi^* \rangle T^a{}^j & \langle W^{ij} \rangle \end{pmatrix}$$

zero eigenvector $\propto (\frac{1}{\sqrt{2}} \langle D^a \rangle, \langle F_i \rangle,$

$$\tilde{\chi} = \frac{1}{\sqrt{W}} \begin{pmatrix} \frac{1}{\sqrt{2}} \langle D^a \rangle \lambda^a \\ \langle F_i \rangle \psi_i \end{pmatrix}$$

W gauge inv

$$\langle \phi^* \rangle T^a{}^i F_i = - \langle \phi^* \rangle T^a{}^i W_i{}^{*} = 0$$

$$V = \frac{1}{2} D^a D^a + F_i F_i{}^*$$

$$\frac{\partial V}{\partial \phi_j} = - g_a (\phi^* T^a{}^j) D^a - W_i{}^j \frac{\partial W^i}{\partial \phi_j}$$

$$\begin{aligned} \left\langle \frac{\partial V}{\partial \phi_j} \right\rangle &= - g_a \langle \phi^* \rangle T^a{}^j \langle D^a \rangle - \langle F_i \rangle \langle W^{ij} \rangle \\ &= 0 \quad \text{stability} \end{aligned}$$

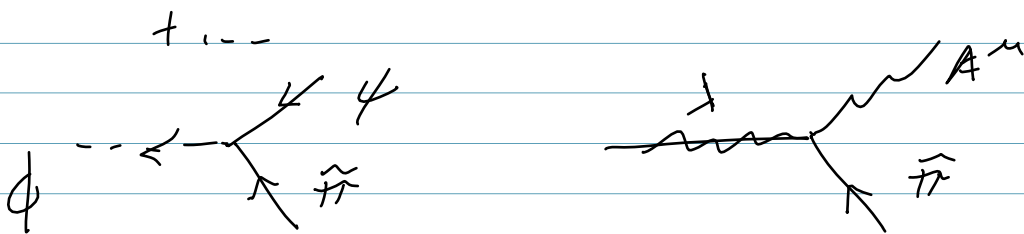
$$J_\alpha^\mu = (\sigma^\nu \bar{\sigma}^\mu \psi_i)_\alpha D_\nu \phi^{*i} - i (\sigma^\mu \psi^{*i})_\alpha W_i^{*\nu} \\ - \frac{1}{2\sqrt{2}} (\sigma^\nu \bar{\sigma}^\rho \sigma^\mu)^{\dagger a}_\alpha F_{\nu\rho} - \frac{i}{\sqrt{2}} g \phi^{*i} T^a \psi_i (\sigma^\mu)^{\dagger a}_\alpha$$

suppose $F_a \neq 0$ $D=0$ $\langle F \rangle^2 = \sum_a |\langle F_a \rangle|^2$

$$\partial_\mu J_\alpha^\mu = i \langle F_a \rangle (\sigma^\mu \partial_\mu \psi^{*a})_\alpha + \partial_\mu \tilde{j}_\alpha^\mu \\ = i \langle F \rangle (\sigma^\mu \partial_\mu \tilde{\pi}^{\dagger})_\alpha + \partial_\mu \tilde{j}_\alpha^\mu$$

$$\Delta_{\text{Goldstino}} = i \tilde{\pi}^{\dagger} \bar{\sigma}^\mu \partial_\mu \tilde{\pi} + \frac{1}{\langle F \rangle} \tilde{\pi} \partial_\mu \tilde{j}_\alpha^\mu$$

$$j_\alpha^\mu = (\sigma^\nu \bar{\sigma}^\mu \psi_i)_\alpha \partial_\nu \phi^{*i} - \frac{1}{2\sqrt{2}} \sigma^\nu \bar{\sigma}^\rho \sigma^\mu)^{\dagger a}_\alpha F_{\nu\rho} \\ + \dots$$



~~$\partial_\mu \tilde{j}_\alpha^\mu$~~ has two derivatives

$$\rightarrow \frac{\Delta m^2}{\langle F \rangle} \rightarrow \text{const as } \langle F \rangle \rightarrow 0$$

for reasonable $\Delta m^2 \sim (100 \text{ GeV})^2$
 $\langle F \rangle$ small enough

$\frac{\Delta m^2}{\langle F \rangle}$ interactions important

If include gravity \Rightarrow SUSY gravity

SUSY transformations must be local
 $E_\alpha \rightarrow E_\alpha(x) \Rightarrow$ supergravity

Spin 2 graviton \leftrightarrow Spin $3/2$ gravitino

inhomogeneous SUSY transformation
 $\delta \Omega_\mu^\alpha$

$$\delta \Omega_\mu^\alpha = -2\partial_\mu \epsilon^\alpha + \dots$$

analogous to gauge particle of SUSY

when SUSY is broken

gravitino eats the goldstino

(other SuperHiggs mechanism)

$2(3/2) + 1 = 4$ helicity states

$$m_{3/2} = \frac{\langle F \rangle}{M_{pl}} \quad \text{cf } m_W \propto g \langle H \rangle$$

gravity mediated: $m_{3/2} \sim m_{soft}$

$\frac{\Delta m_{soft}}{F} = \frac{m_{soft}}{F} \frac{F}{M_{pl}} \Rightarrow$ gravity strength interactions
not important for colliders

$= \frac{m_{soft}}{M_{pl}}$ can overdose universe

gauge mediation $m_{3/2} = \frac{F}{M_{pl}}$ $m_{soft} = \frac{\kappa}{4\pi} \frac{F}{M_{mess}}$

if $M_{mess} \ll M_{pl}$ $m_{3/2} \ll m_{soft}$

gravitino is the LSP

goldstinos have non-gravitational interaction,
gravitino components are irrelevant at colliders

$$\Gamma(\tilde{X} \rightarrow X \tilde{\pi}) = \frac{m_X^5}{16\pi \langle F \rangle^2} \left(1 - \frac{m_X^2}{m_{\tilde{X}}^2}\right)^4$$

$$m_{\tilde{X}} \approx 100 \text{ GeV} \gg m_X$$

if $\sqrt{\langle F \rangle} < 10^3 \text{ TeV}$

$$\Gamma > 2 \times 10^{-16} \text{ GeV}$$

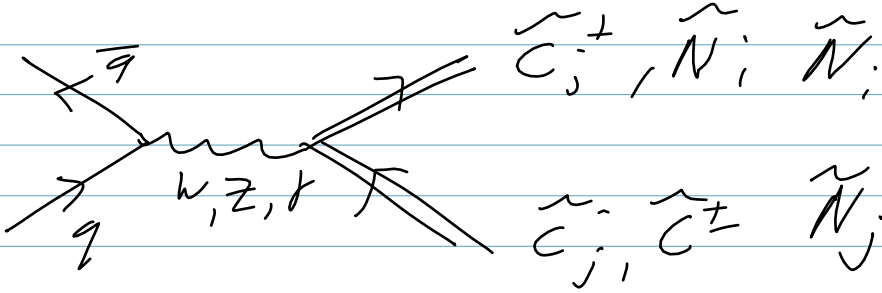
$$c\tau = \frac{c}{\Gamma} < 5 \times 10^{15} \text{ GeV}^{-1} \frac{1 \text{ s}}{1.5 \times 10^{24} \text{ GeV}^{-1}} \frac{3 \times 10^8 \text{ m/s}}{1}$$

$$< 1 \text{ m}$$

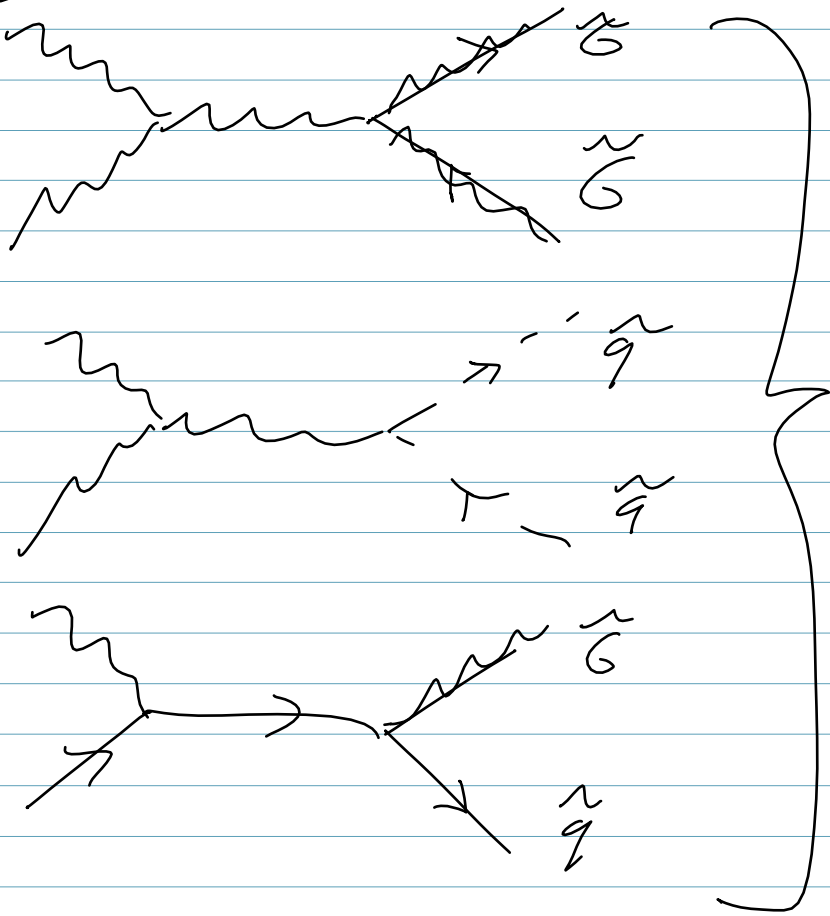
~~$\tilde{X} \rightarrow X + \tilde{\pi}$~~ could be observed

Finding SUSY

~~LHC~~

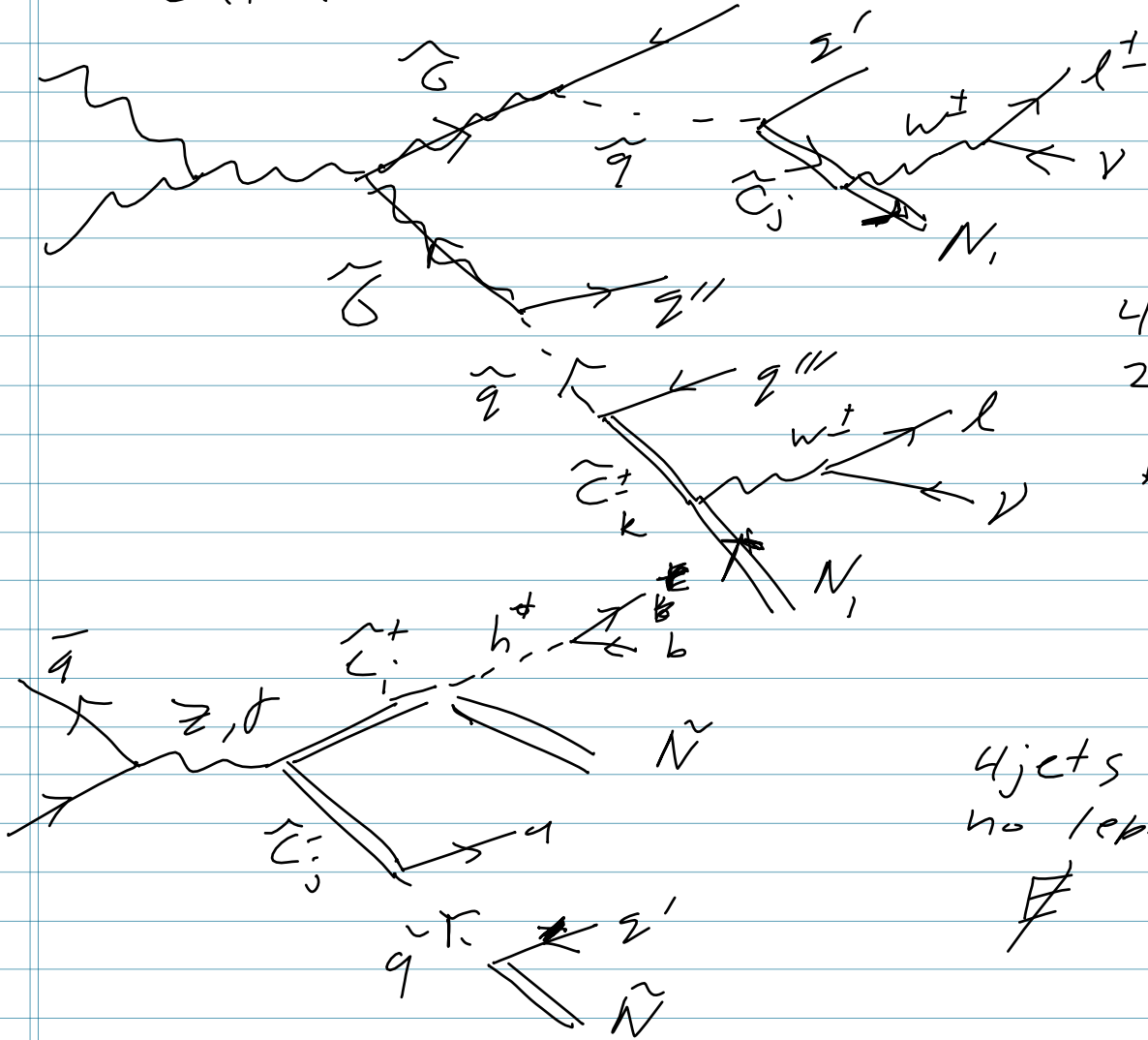


LHC



LHC
 pp

Neutralino LSP



4 jets
2 leptons
Same sign
50%

4 jets
no leptons

goldstino LSP, NLSP $U(1)_Y \rightarrow \tilde{N}, \tilde{L}_R$
decay length submicron to multi km $\sqrt{F} > 10^3 \text{ TeV}$

NLSP escape

\tilde{N} NLSP Super pair $\rightarrow X + \tilde{N} + \tilde{N}$
 $\hookrightarrow \gamma \tilde{\pi} \hookrightarrow \gamma \tilde{\pi}$

$\rightarrow X + 2\gamma + \cancel{\#}$ short decay length

\tilde{L}_R NLS Super pair $\rightarrow X + 2\gamma + \cancel{\#}$

Goldstino Theorem

$$Q^\alpha |0\rangle \neq 0 \iff \langle 0 | H | 0 \rangle \neq 0$$

$$\{ Q_\alpha, Q_\alpha^\dagger \} = 2 \sigma_{\alpha\dot{\alpha}}^\mu P_\mu$$

$$Q_\alpha^\dagger = \sqrt{2} \int d^3x J_\alpha^{\dagger 0}(\vec{x})$$

$$\int d^3x \langle 0 | \{ Q_\alpha, \sqrt{2} \int d^3x J_\alpha^{\dagger 0}(\vec{x}) \} | 0 \rangle = 2 \sigma_{\alpha\dot{\alpha}}^\mu \int d^3x \langle 0 | T_{\mu}^0(\vec{x}) | 0 \rangle$$

$$\frac{1}{\sqrt{2}} \langle 0 | \{ Q_\alpha, J_\alpha^{\dagger 0}(0) \} | 0 \rangle = 2 \sigma_{\alpha\dot{\alpha}}^\mu \langle 0 | T_{\mu}^0(0) | 0 \rangle$$

$$0 \neq \langle 0 | \{ \int d^3x J_\alpha^0(\vec{x}), J_\alpha^{\dagger 0}(0) \} | 0 \rangle \stackrel{E=3m}{=} \\ = \int d^3x \langle 0 | J_\alpha^0(\vec{x}) J_\alpha^{\dagger 0}(0) + J_\alpha^{\dagger 0}(0) J_\alpha^0(\vec{x}) | 0 \rangle \\ \neq 0 \Rightarrow \text{susy broken}$$

$$\begin{aligned}
& \int d^3x \langle 0 | J_\alpha^0(\vec{x}) J_\alpha^{+\nu}(0) + J_\alpha^{+\nu}(0) J_\alpha^0(\vec{x}) | 0 \rangle \\
&= \sum_n \int d^3x \langle 0 | e^{i\vec{p}_n \cdot \vec{x}} J_\alpha^0(0) e^{-i\vec{p}_n \cdot \vec{x}} | n \rangle \langle n | J_\alpha^{+\nu}(0) | 0 \rangle \\
&\quad + \langle 0 | J_\alpha^{+\nu}(0) e^{i\vec{p}_n \cdot \vec{x}} J_\alpha^0(0) e^{-i\vec{p}_n \cdot \vec{x}} | 0 \rangle \\
&= \sum_n \int d^3x \langle 0 | J_\alpha^0(0) e^{-i\vec{p}_n \cdot \vec{x}} | n \rangle \langle n | J_\alpha^{+\nu}(0) | 0 \rangle \\
&\quad + \langle 0 | J_\alpha^{+\nu}(0) | n \rangle \langle n | e^{i\vec{p}_n \cdot \vec{x}} J_\alpha^0(0) | 0 \rangle \\
&= \sum_n (2\pi)^3 \delta^{(3)}(\vec{p}_n) \left(\langle 0 | J_\alpha^0(0) | n \rangle \langle n | J_\alpha^{+\nu}(0) | 0 \rangle \right. \\
&\quad \left. + \langle 0 | J_\alpha^{+\nu}(0) | n \rangle \langle n | J_\alpha^0(0) | 0 \rangle \right) \\
&= \sum_n (2\pi)^3 \delta^{(3)}(\vec{p}_n) f(E_n) \quad \text{①}
\end{aligned}$$

$$\begin{aligned}
& \int d^4x \langle 0 | J_\alpha^0(x) J_\alpha^{+\nu}(0) + J_\alpha^{+\nu}(0) J_\alpha^0(x) | 0 \rangle \delta(t) \\
&= \int d^4x \partial_\rho \left(\langle 0 | J_\alpha^\rho(x) J_\alpha^{+\nu}(0) | 0 \rangle \Theta(t) \right. \\
&\quad \left. - \langle 0 | J_\alpha^{+\nu}(0) J_\alpha^\rho(x) | 0 \rangle \Theta(-t) \right) \\
&= \sum_n \int d^4x \partial_\rho \left(\langle 0 | J_\alpha^\rho(0) e^{-i\vec{p}_n \cdot x} | n \rangle \langle n | J_\alpha^{+\nu}(0) | 0 \rangle \Theta(t) \right. \\
&\quad \left. - \langle 0 | J_\alpha^{+\nu}(0) | n \rangle \langle n | e^{i\vec{p}_n \cdot x} J_\alpha^\rho(0) | 0 \rangle \Theta(-t) \right) \\
&= \sum_n \int d^4x (-i\vec{p}_n \cdot \rho) \left(e^{-i\vec{p}_n \cdot x} \langle 0 | J_\alpha^\rho(0) | n \rangle \langle n | J_\alpha^{+\nu}(0) | 0 \rangle \Theta(t) \right. \\
&\quad \left. + e^{i\vec{p}_n \cdot x} \langle 0 | J_\alpha^{+\nu}(0) | n \rangle \langle n | J_\alpha^\rho(0) | 0 \rangle \Theta(-t) \right) \\
&\quad \delta(t) \left(e^{-i\vec{p}_n \cdot x} \langle 0 | J_\alpha^0(0) | n \rangle \langle n | J_\alpha^{+\nu}(0) | 0 \rangle \right. \\
&\quad \left. + e^{i\vec{p}_n \cdot x} \langle 0 | J_\alpha^{+\nu}(0) | n \rangle \langle n | J_\alpha^0(0) | 0 \rangle \right)
\end{aligned}$$

$$\begin{aligned}
&= \sum_n (2\pi)^3 \delta^{(3)}(\vec{p}_n) \left(\int dt (-iE_n) \langle 0 | J_\alpha^0(0) | n \rangle \langle n | J_\alpha^{+\nu}(0) | 0 \rangle e^{-iE_n t} \right. \\
&\quad \left. + e^{iE_n t} \langle 0 | J_\alpha^{+\nu}(0) | n \rangle \langle n | J_\alpha^0(0) | 0 \rangle \right) \\
&\quad + \langle 0 | J_\alpha^0(0) | n \rangle \langle n | J_\alpha^+(0) | 0 \rangle \\
&\quad - \langle 0 | J_\alpha^{+\nu}(0) | n \rangle \langle n | J_\alpha^0(0) | 0 \rangle \\
&= \sum_n (2\pi)^3 \delta^{(3)}(\vec{p}_n) \left(-i \int_0^\infty dt e^{-iE_n t} E_n f(E_n) \right. \\
&\quad \left. + f(E_n) \right)
\end{aligned}$$

$$\int_0^\infty dt e^{-iE_n t} E_n f(E_n) = 0$$

$$f(E_n) \neq 0 \Rightarrow \text{broken susy}$$

$$f(E_n) \propto \delta(E_n)$$

\exists a fermionic state with $\vec{p}=0, E=0$