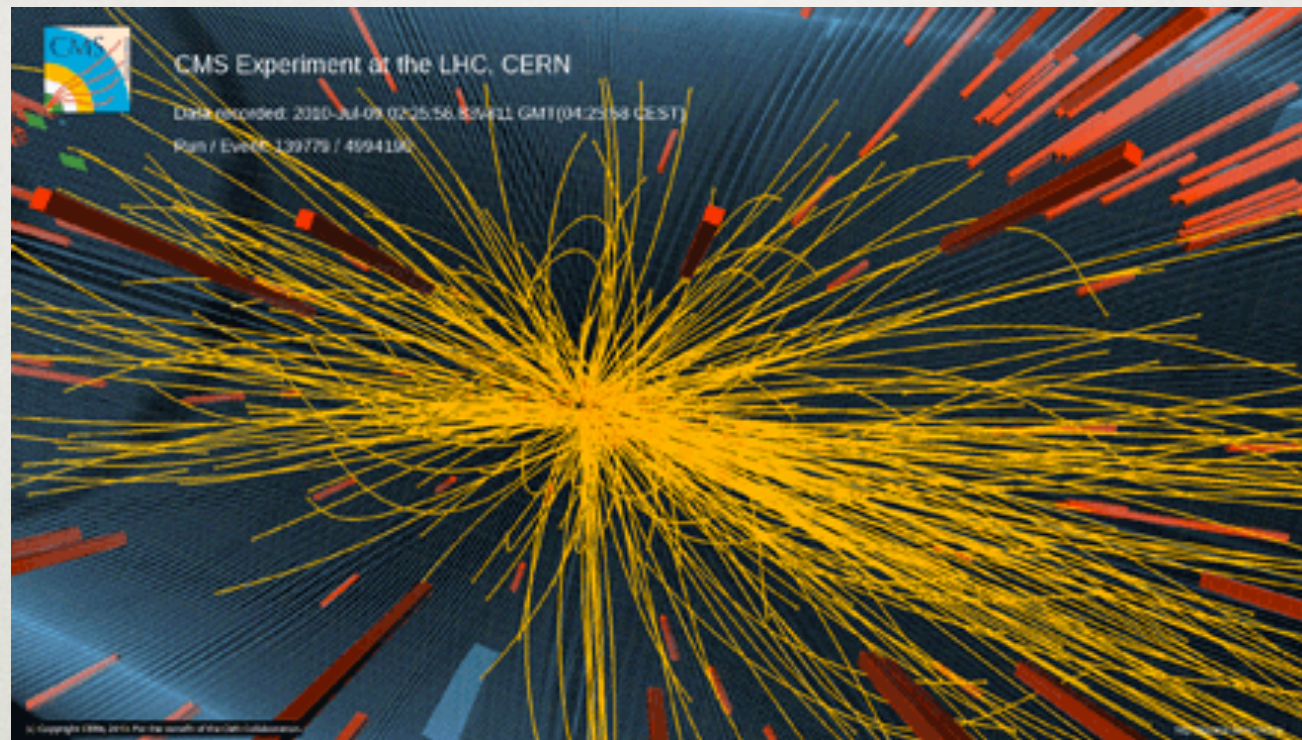


LOW-ENERGY SIGNS OF NEW PHYSICS AT HIGH-ENERGY COLLIDERS



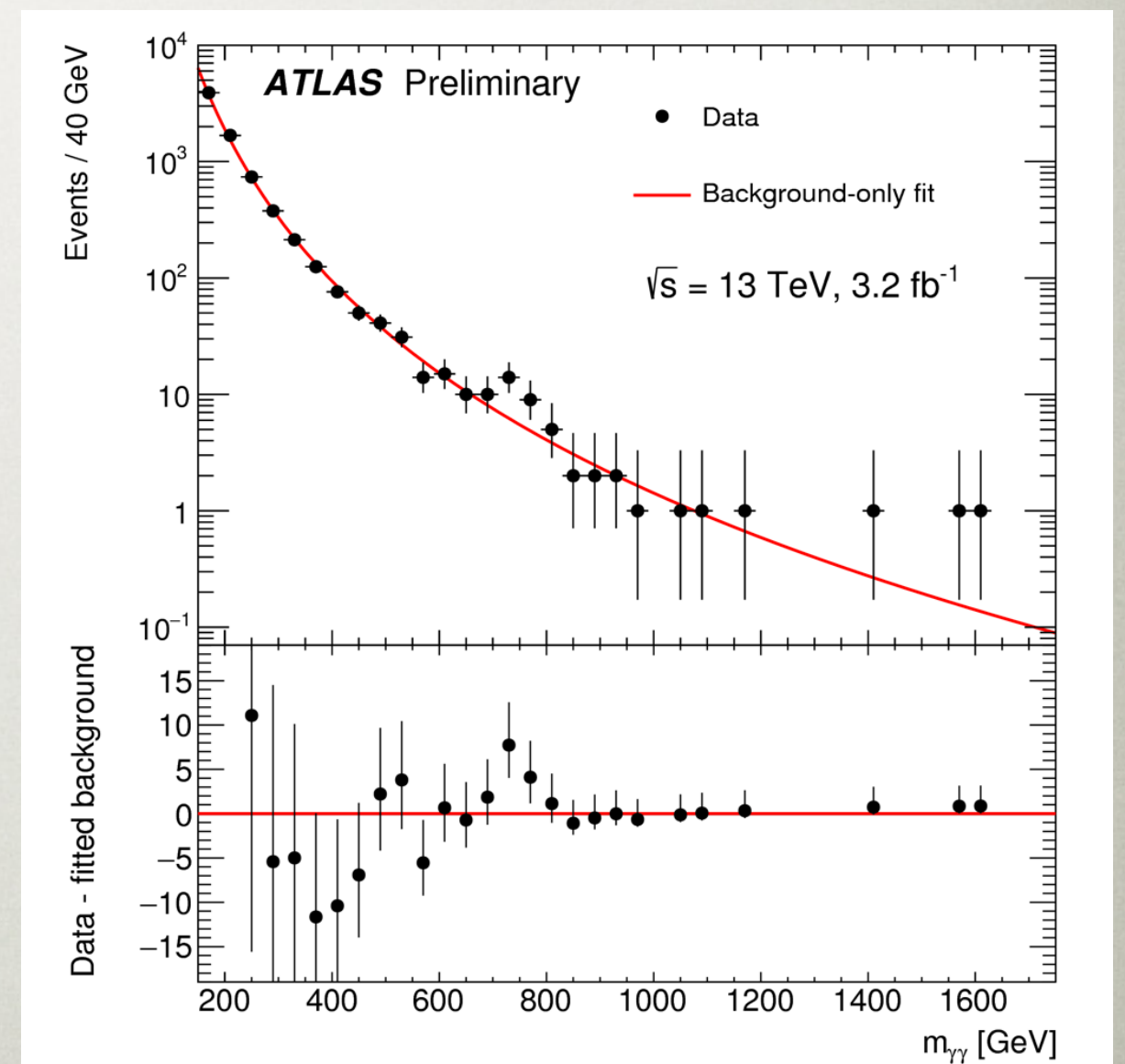
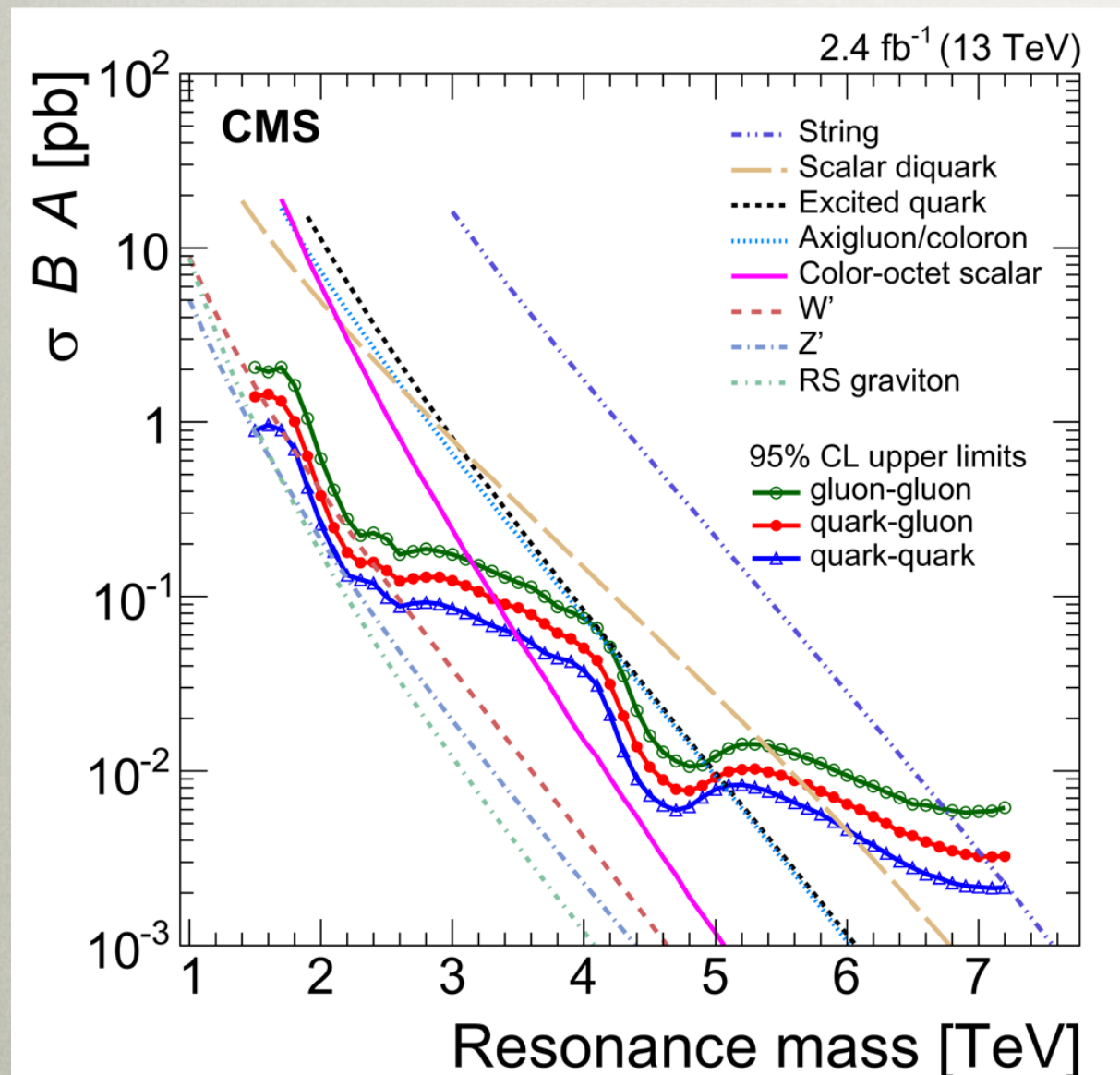
Brian Shuve -- SLAC

HEFTI Seminar, UC Davis

8 February 2016

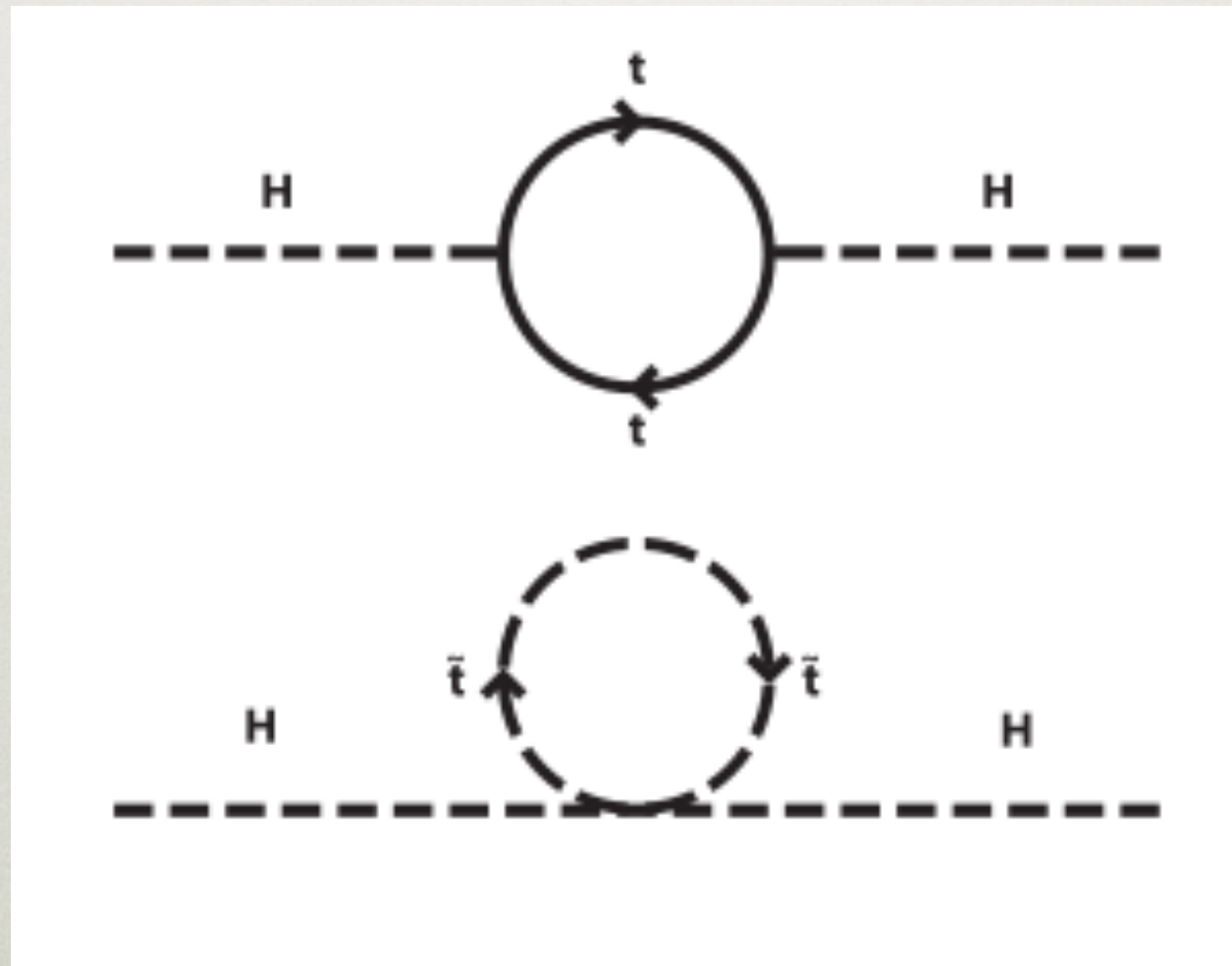
Energy Frontier @ LHC

- Exploring new energy regimes

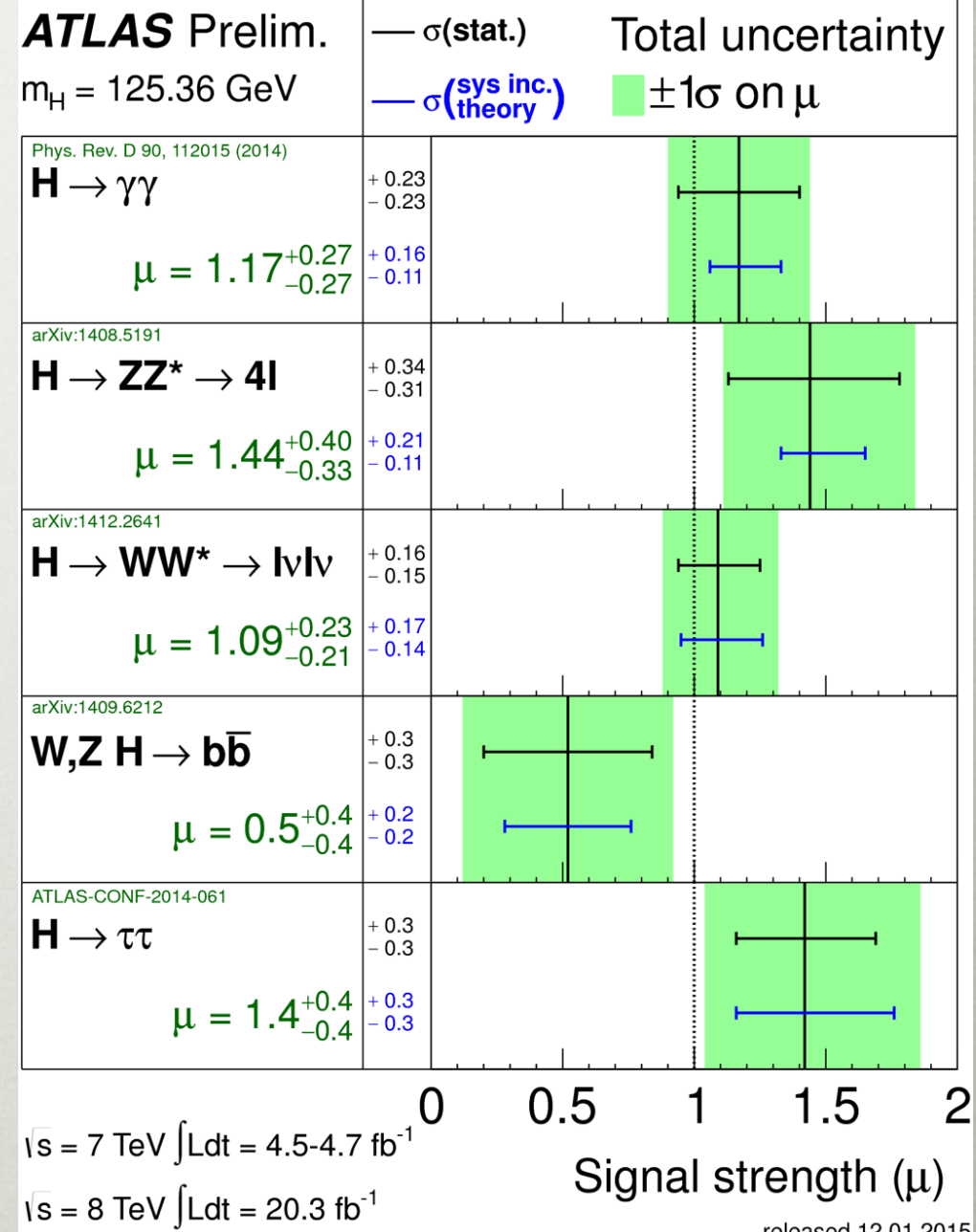
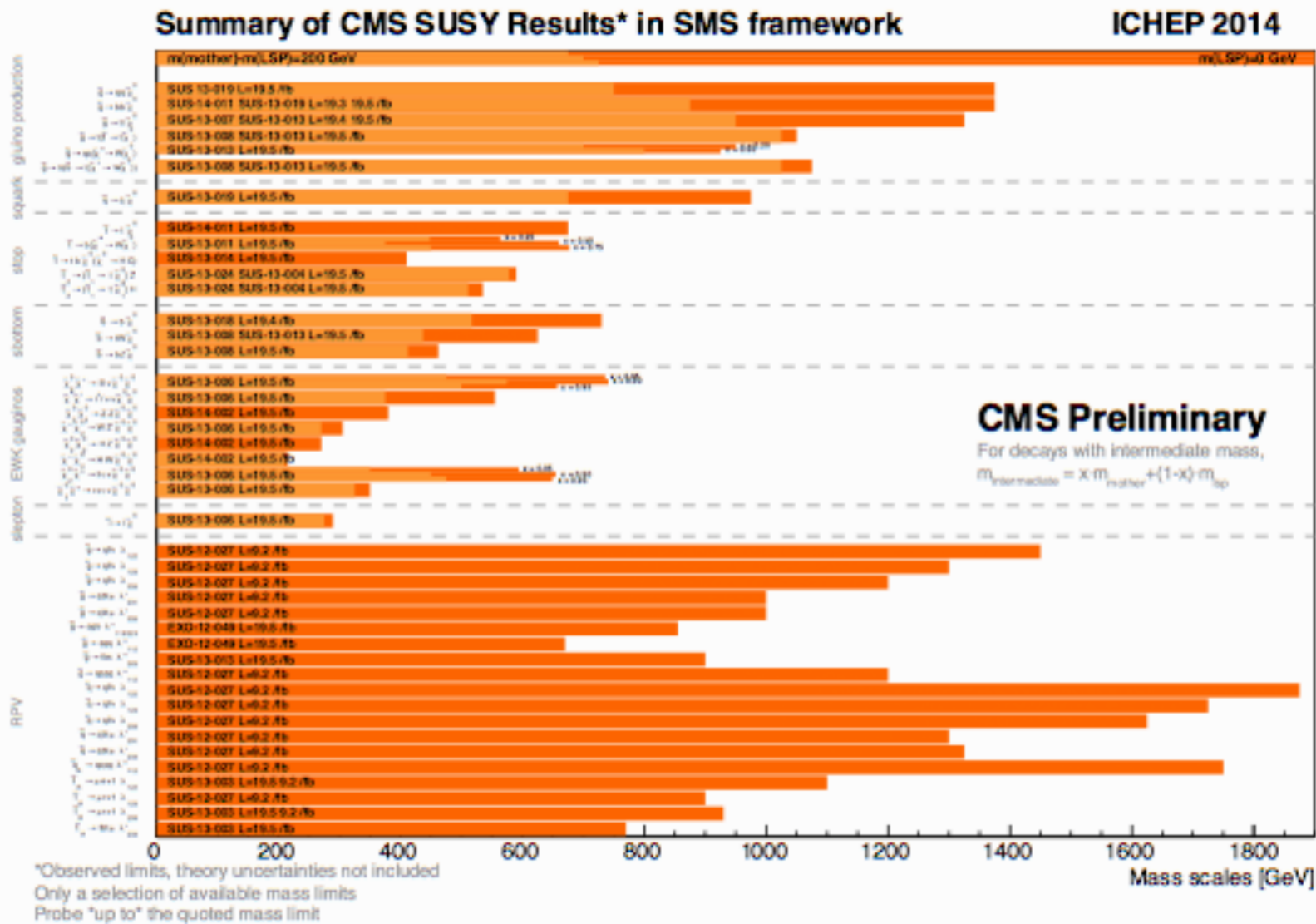


Energy Frontier @ LHC

- Expect new high-scale physics because of **naturalness**

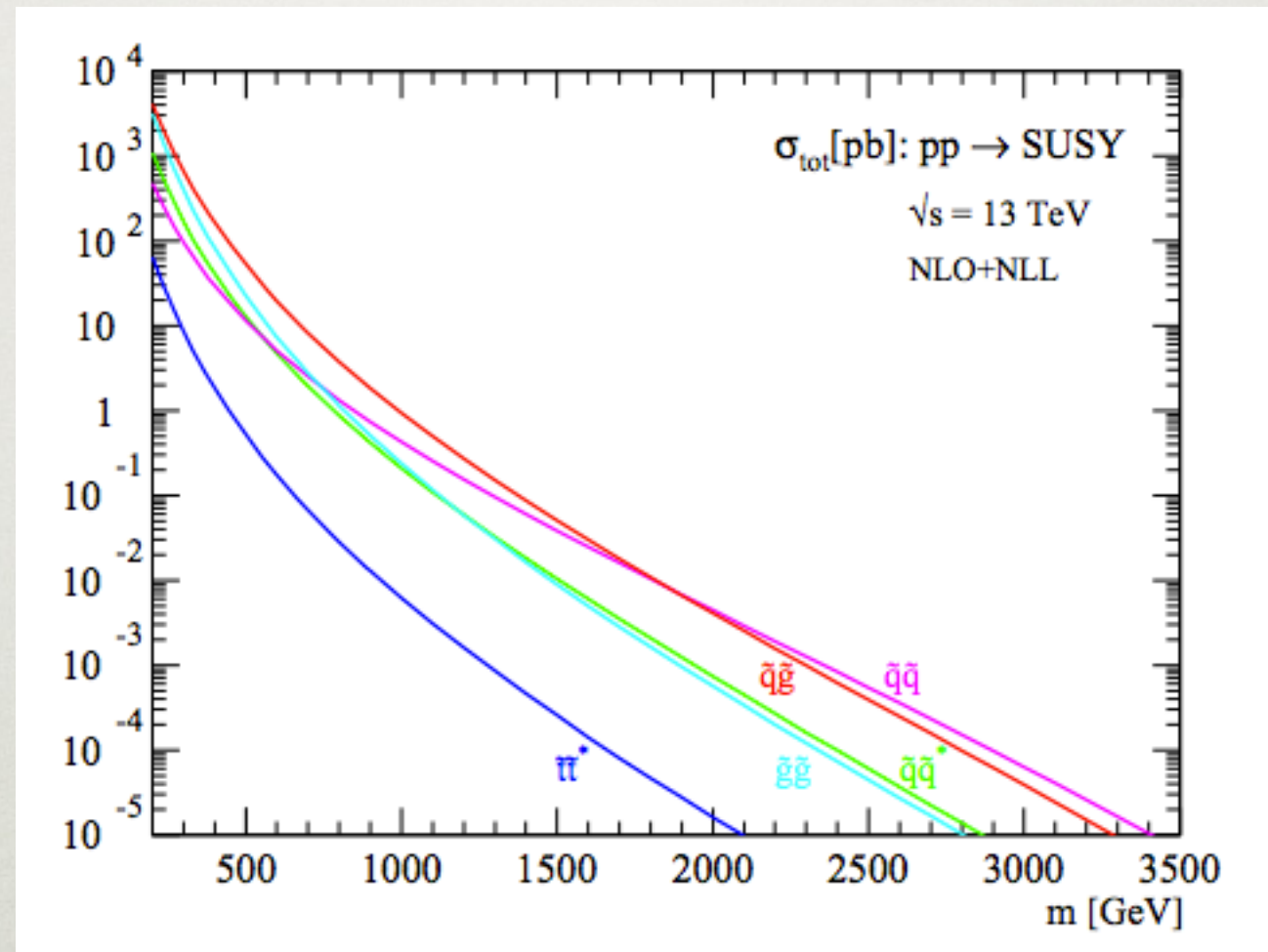


Energy Frontier @ LHC



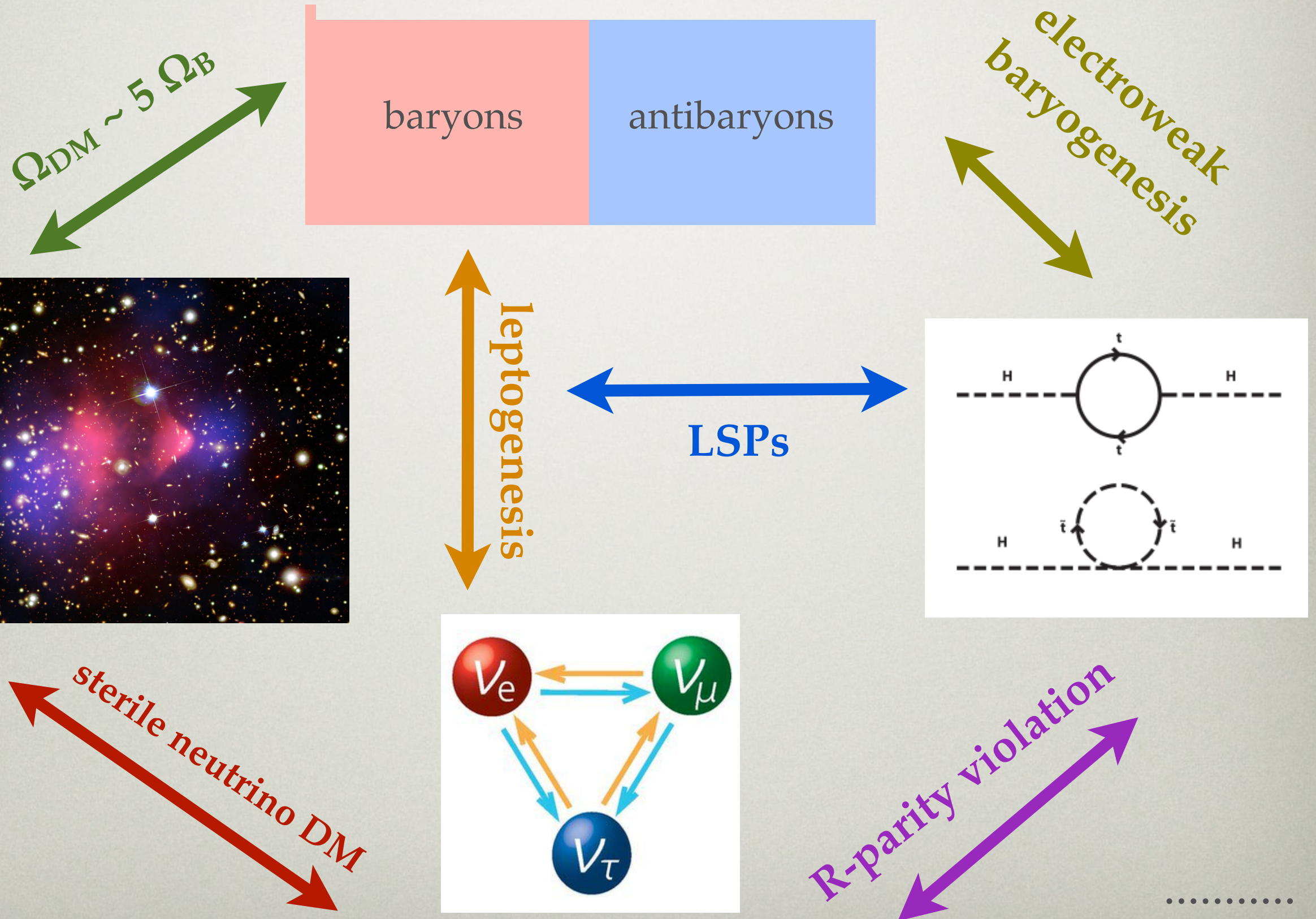
Energy Frontier @ LHC

- Should be easy to find new, high-mass physics



- ...but diminishing returns set in quickly

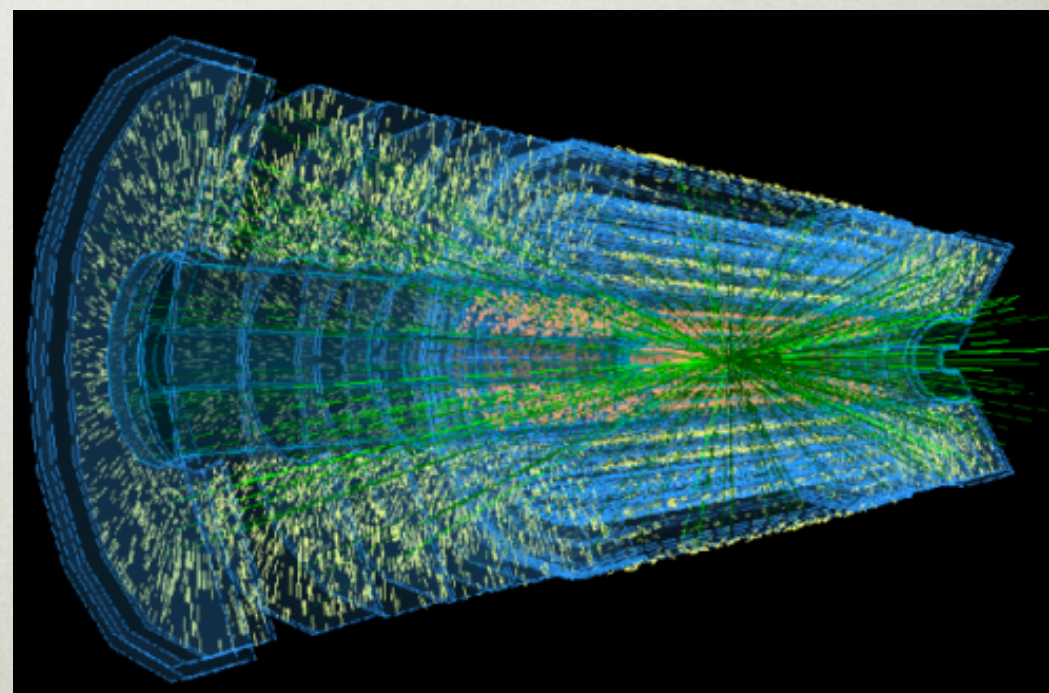
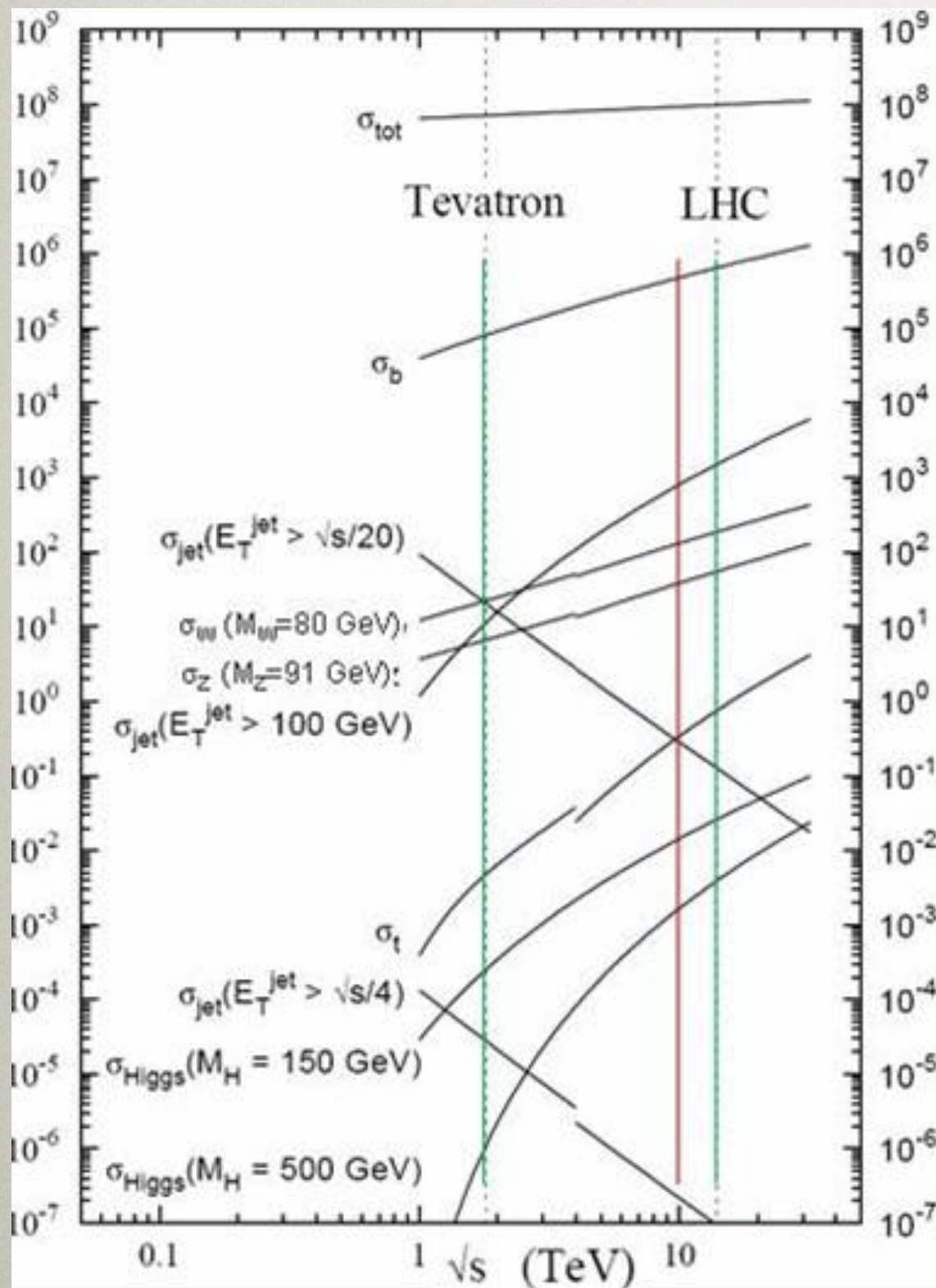
Unsolved Mysteries of the SM



The LHC is High Intensity

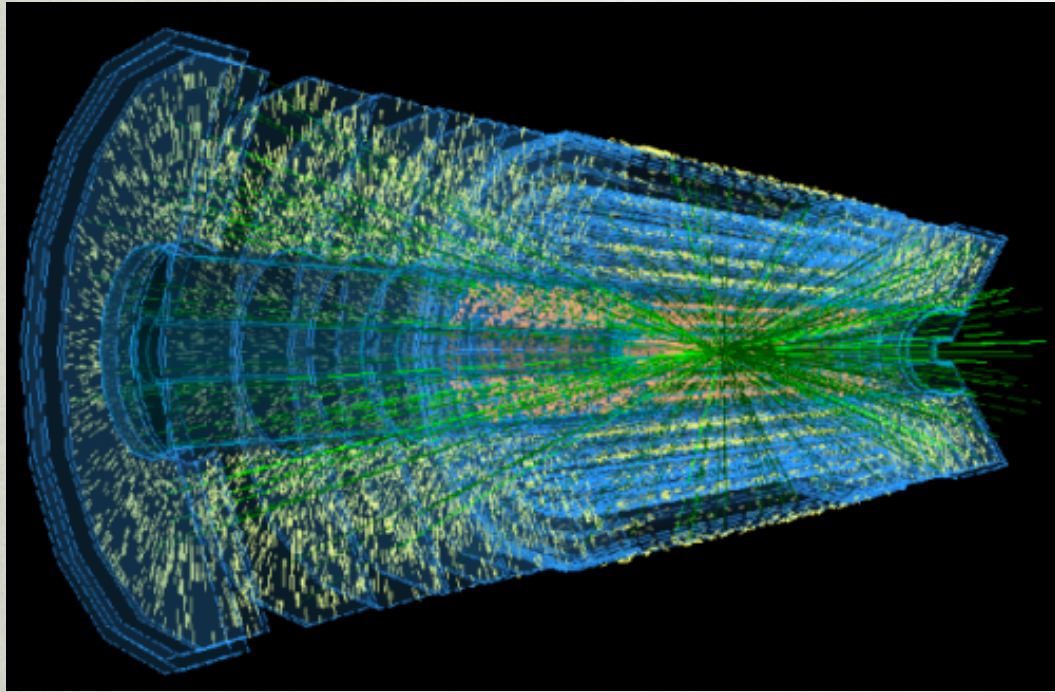
With 3000/fb @ 14 TeV:

- 700 billion W
- 100 billion Z
- 200 million H
- 2000 trillion B



From J. Stirling

Low-Energy Signatures



To see signatures of sub-weak masses, should be **spectacular**

Examples:

- long-lived particles
- multileptons

Low-Energy Signatures

Why long-lived?

$$c\tau(\pi^\pm) \sim 10 \text{ m}$$

$$c\tau(K_L^0) \sim 1 \text{ cm}$$

$$c\tau(D^\pm) \sim 0.1 \text{ mm}$$

$$c\tau(B^\pm) \sim 0.1 \text{ mm}$$

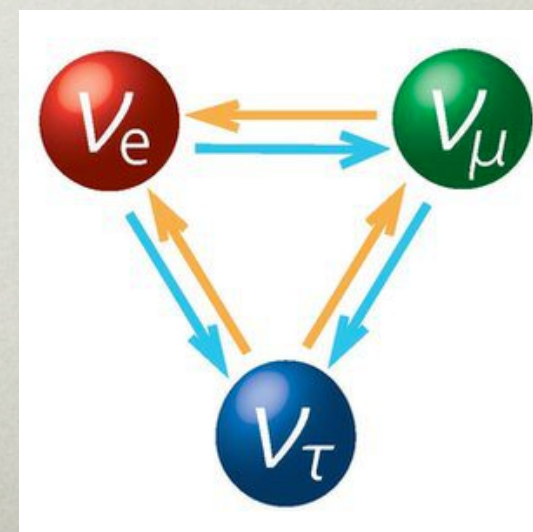
$$c\tau(J/\psi) \sim 1 \text{ pm}$$

- Small mixing angles
- Approximate symmetries
- Decays through heavy, off-shell states (W)

Why multileptons?

baryons

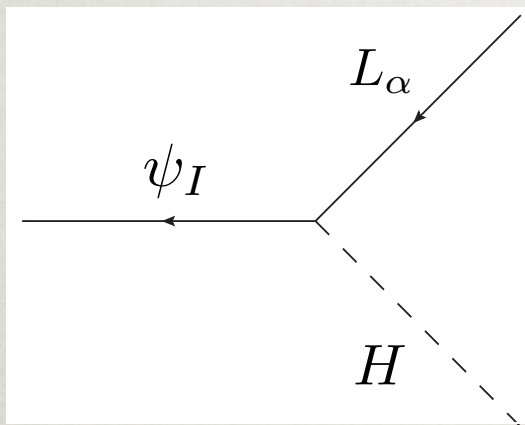
antibaryons



Hidden Sector Portals

- Dominant couplings of new singlets are via **renormalizable portals**

NEUTRINO PORTAL



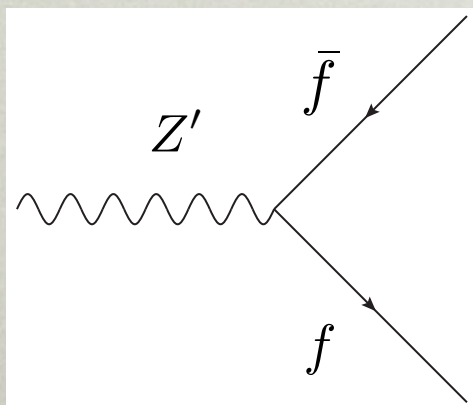
$$\mathcal{L}_\nu = y_{I\alpha} \bar{\psi}_I H L_\alpha$$

VECTOR PORTAL



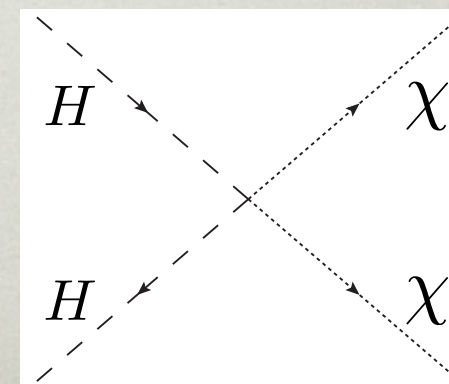
$$\mathcal{L}_{\text{vector}} = -\frac{\epsilon}{2} F'_{\mu\nu} F^{\mu\nu}$$

Z' PORTAL



$$\mathcal{L}_{Z'} = g_{Z'} \bar{f}_{\text{SM}} \gamma^\mu f_{\text{SM}} Z'_\mu$$

HIGGS PORTAL



$$\mathcal{L}_H = -\lambda |H|^2 |\chi|^2$$

Outline

- **Soft signatures of baryon and lepton number violation**
 - The neutrino portal and leptogenesis
 - The Z' portal: discovery of new B/L gauge forces

- **Soft signatures of dark matter**
 - The vector portal and inelastic dark matter

Motivations: Neutrino Masses

- Simplest UV completion: RH neutrinos (Type-I See-saw)

$$\mathcal{L}_{\nu\text{MSM}} = F_{\alpha I} L_{\alpha} \Phi N_I + \frac{M_I}{2} N_I^2 \quad (m_{\nu})_{\alpha\beta} = \langle \Phi \rangle^2 (F M_N^{-1} F^T)_{\alpha\beta}$$

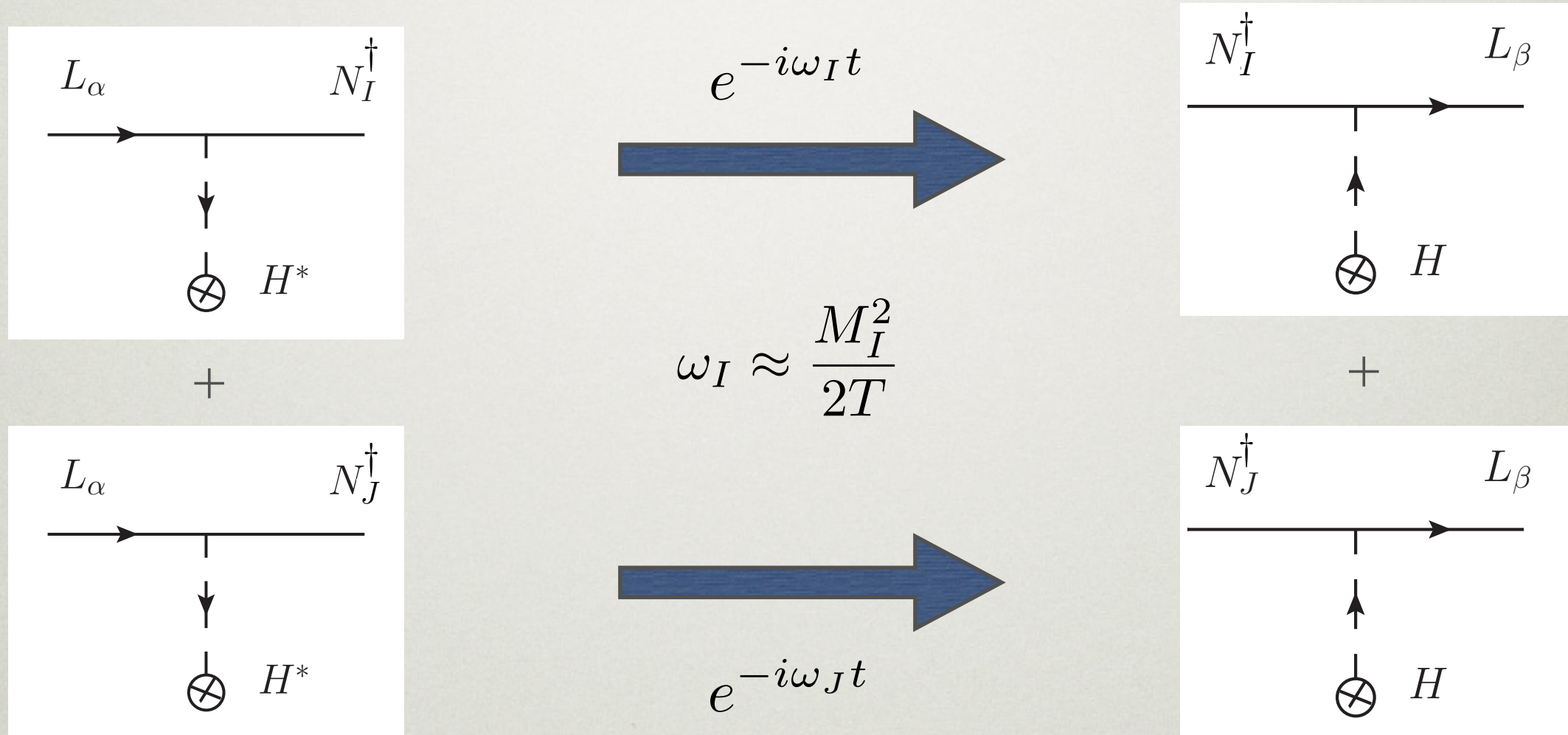
- For fixed $\langle \Phi \rangle^2$ and $m_{\nu} \sim 0.1$ eV, we have $M_N \sim \text{GeV} \left(\frac{F^2}{10^{-14}} \right)$
 - Larger couplings also possible with additional symmetries
- Lepton asymmetry can also be generated by the **decay** of a primordial abundance of N (Fukugita, Yanagida 1986)

$$\Gamma(N \rightarrow L\Phi) \neq \Gamma(N \rightarrow \bar{L}\Phi^*) \quad M_N \gtrsim 100 \text{ GeV}$$

Motivations: Leptogenesis

- Asymmetry can also be generated by the **production and oscillation** of N

Akhmedov, Rubakov, Smirnov, hep-ph/9803255; Asaka, Shaposhnikov, hep-ph/0505013
for more natural models & pheno, see BS, Yavin, 1401.2459



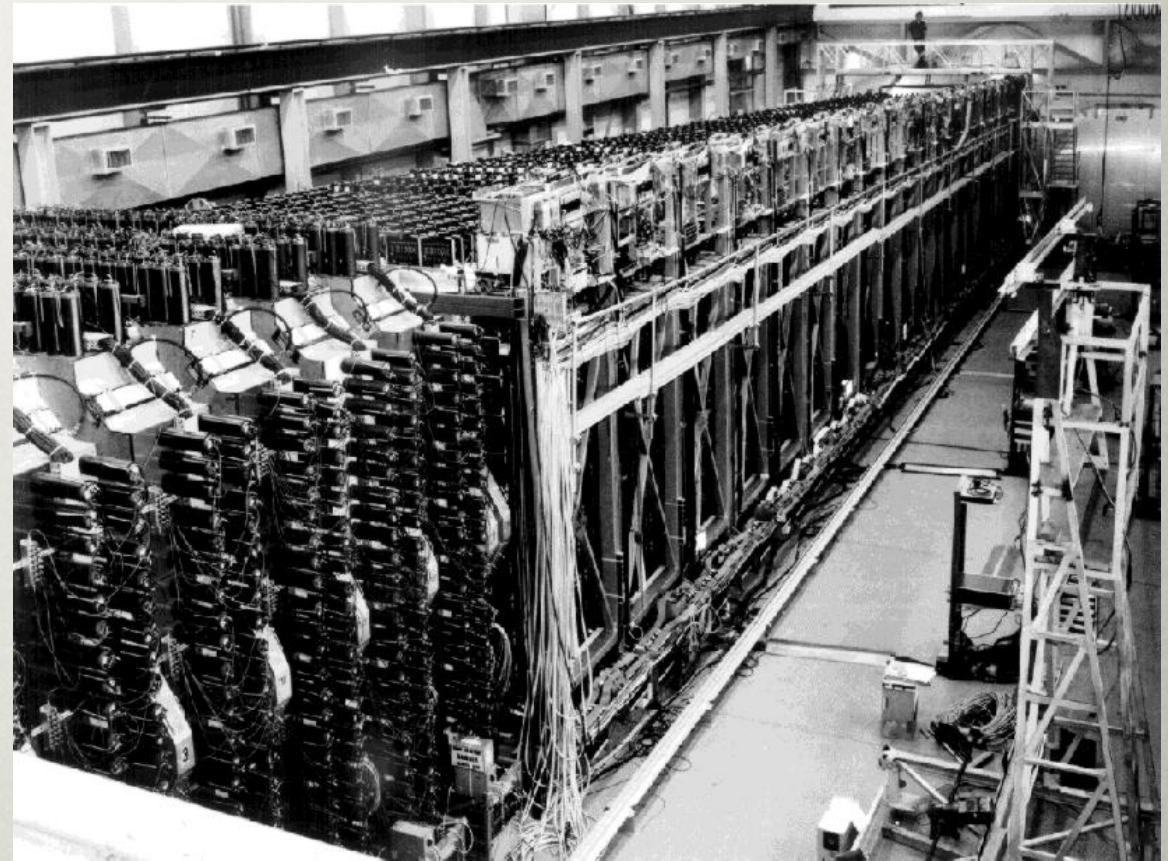
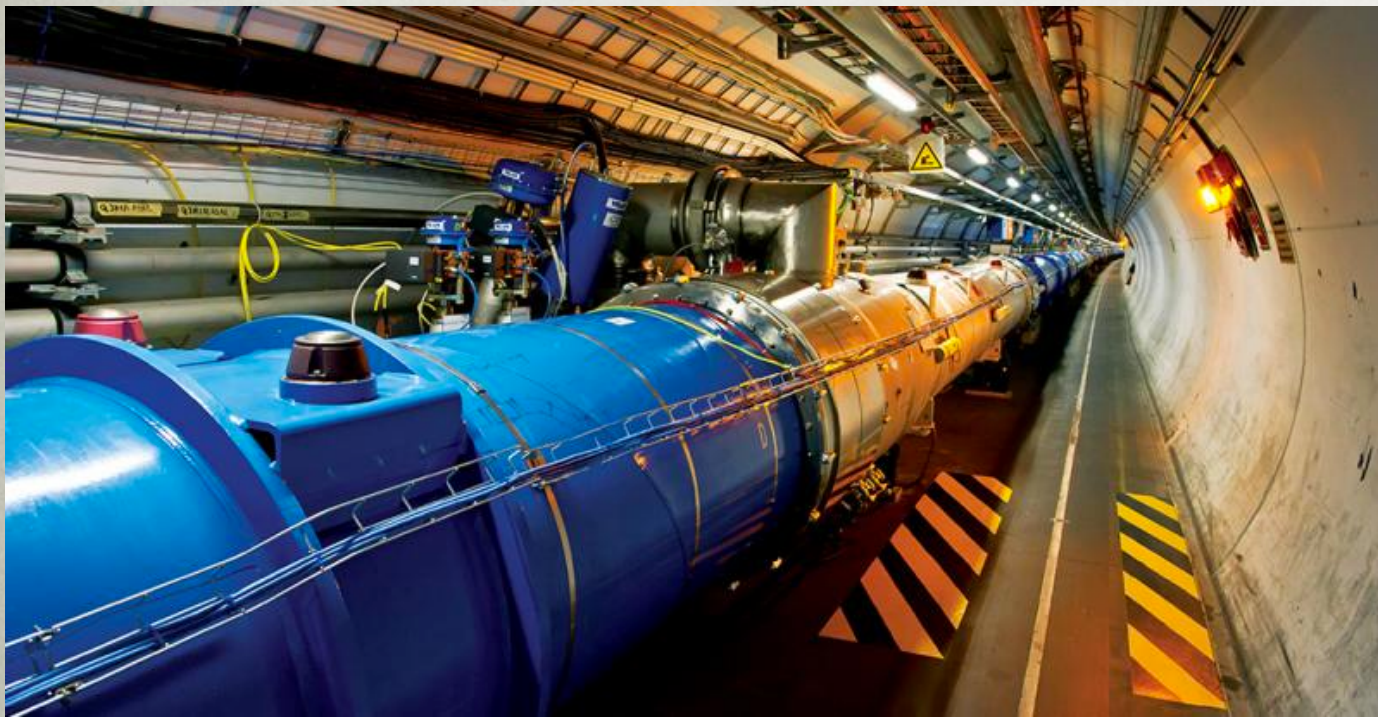
$$\Gamma(L_\alpha \rightarrow L_\beta) - \Gamma(\bar{L}_\alpha \rightarrow \bar{L}_\beta) \sim \text{Im}(F^4) \text{Im}[e^{-i(M_I^2 - M_J^2)t/T}] \text{ becomes large at } T \gg M$$

- Asymmetry lasts until after electroweak transition: $M_N \lesssim 100 \text{ GeV}$

Motivations: Leptogenesis

$$M_N \lesssim 100 \text{ GeV}$$

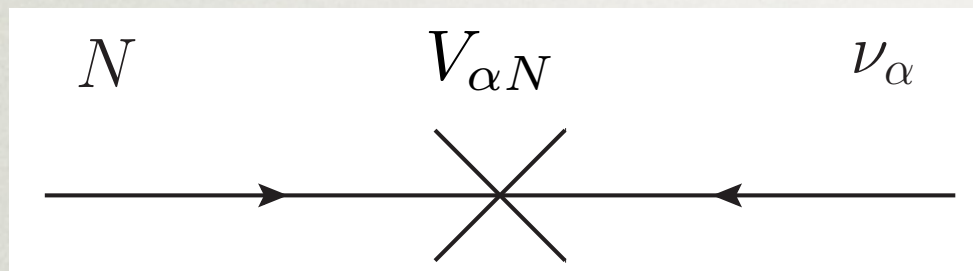
- N can be directly discovered at accelerators!



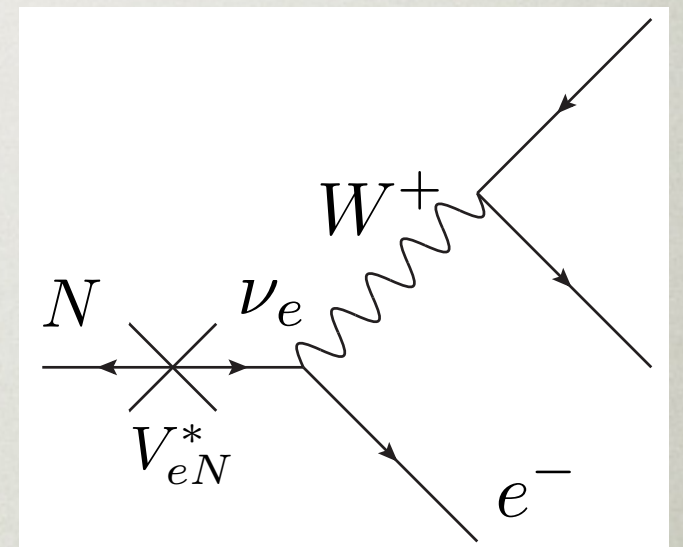
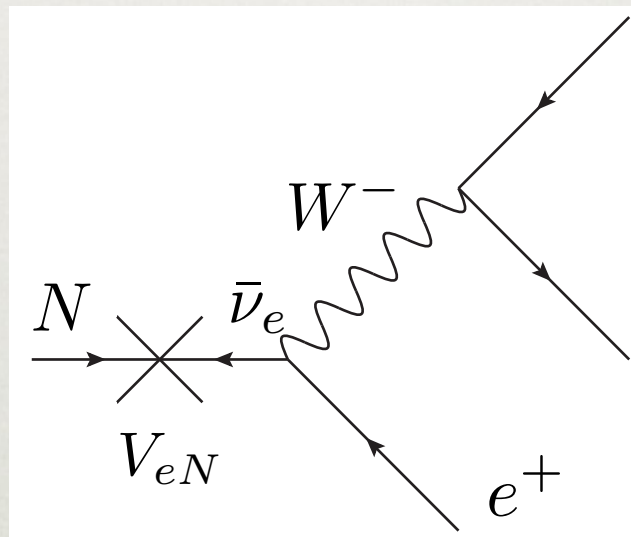
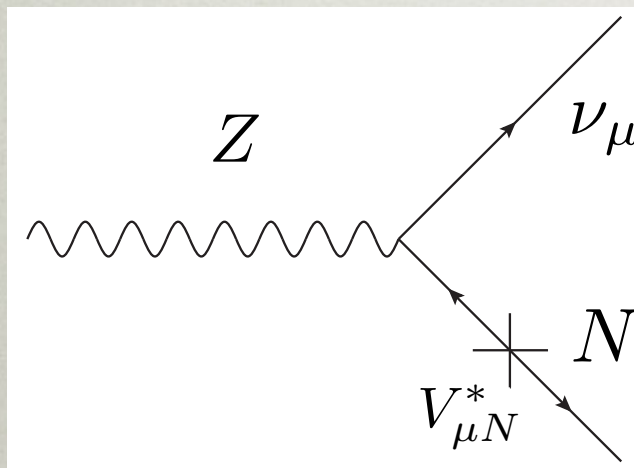
- Also indirect probes (BS, Yavin, 1401.2459)

Neutrinos at Colliders

- After EWSB, the LH and RH neutrinos mix



$$V_{\alpha N} \sim \frac{F_{\alpha} \langle \Phi \rangle}{M_N}$$

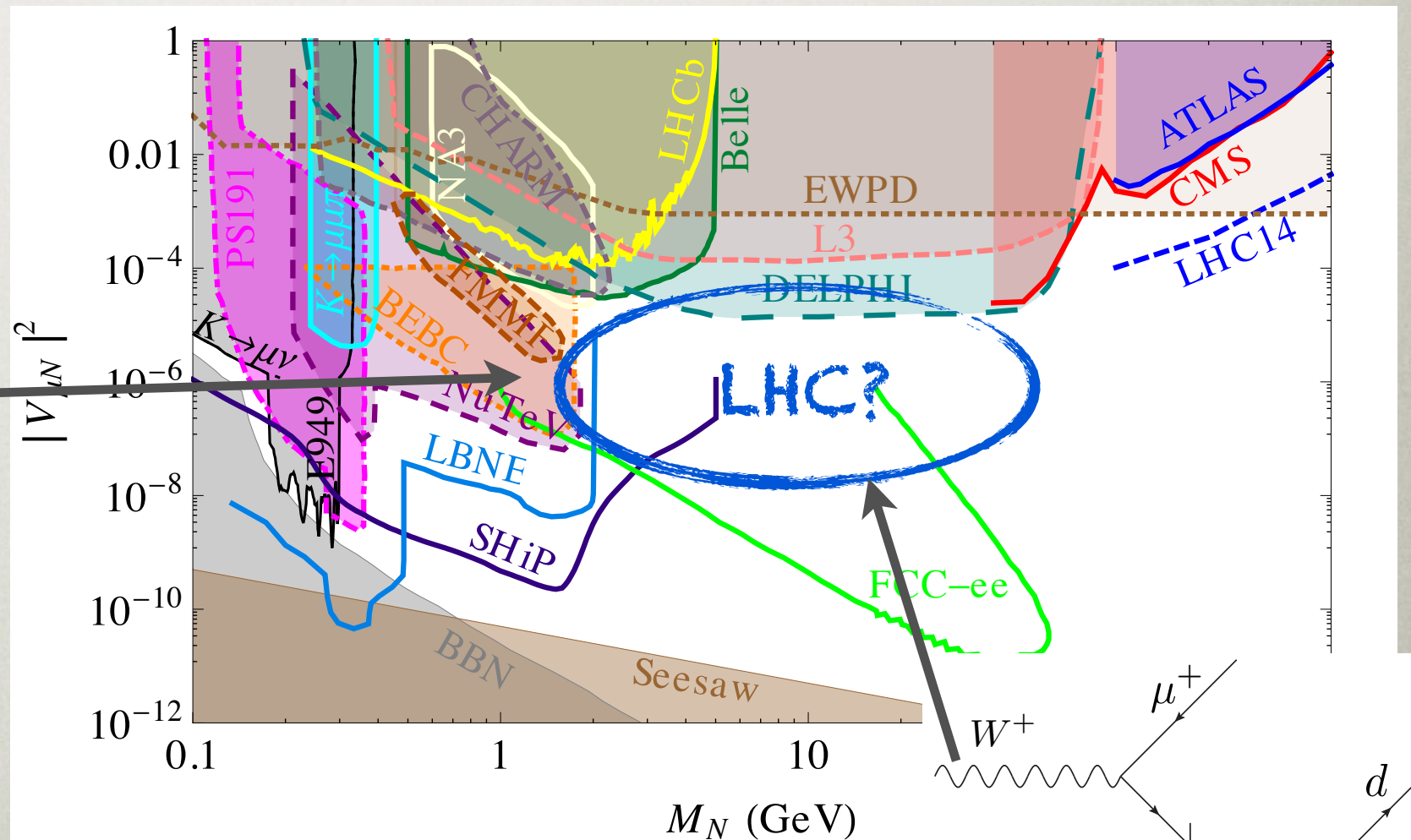
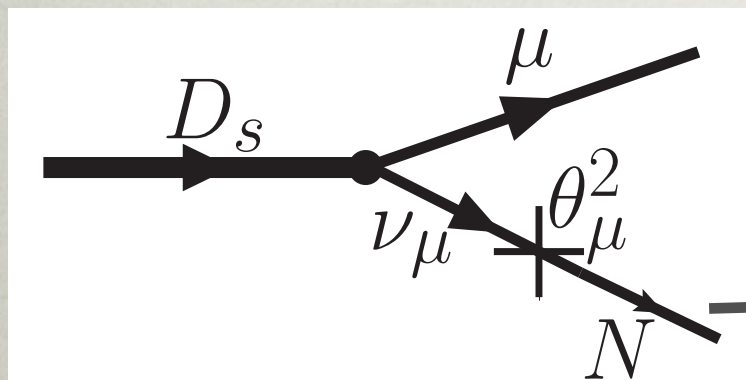


$$\Gamma_N \sim |V_{\alpha N}|^2 G_F^2 M_N^5$$

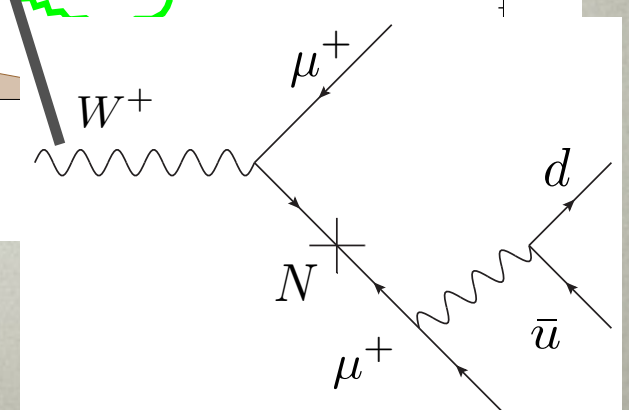
Neutrinos at Colliders: Status

- Simplified Model

- M_N
- $|V_{\mu N}|$ (single-flavour mixing)



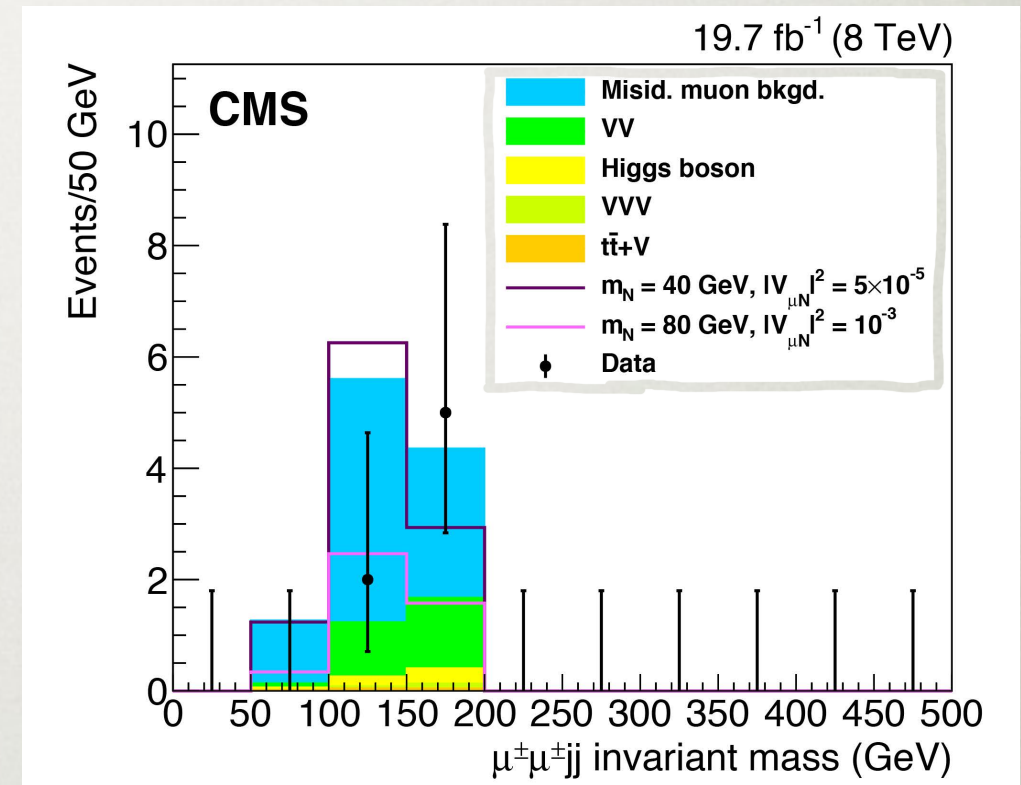
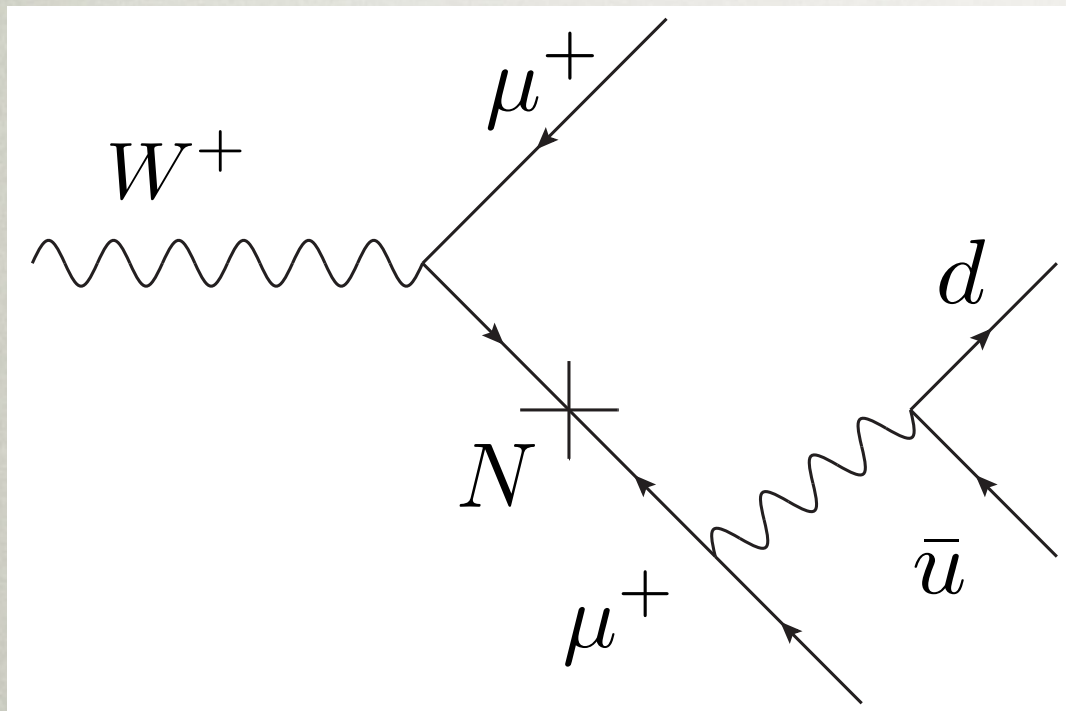
- SHiP: proposed SPS fixed target expt.



Deppisch, Dev, Pilaftsis, 1502.06541

RH Neutrinos at Colliders: Status

- Current strategy: fully reconstructible decay with same-sign leptons
(Keung, Senjanovic, 1983)



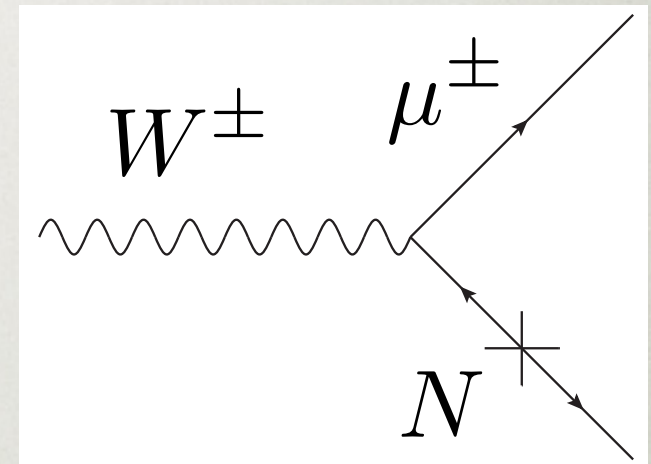
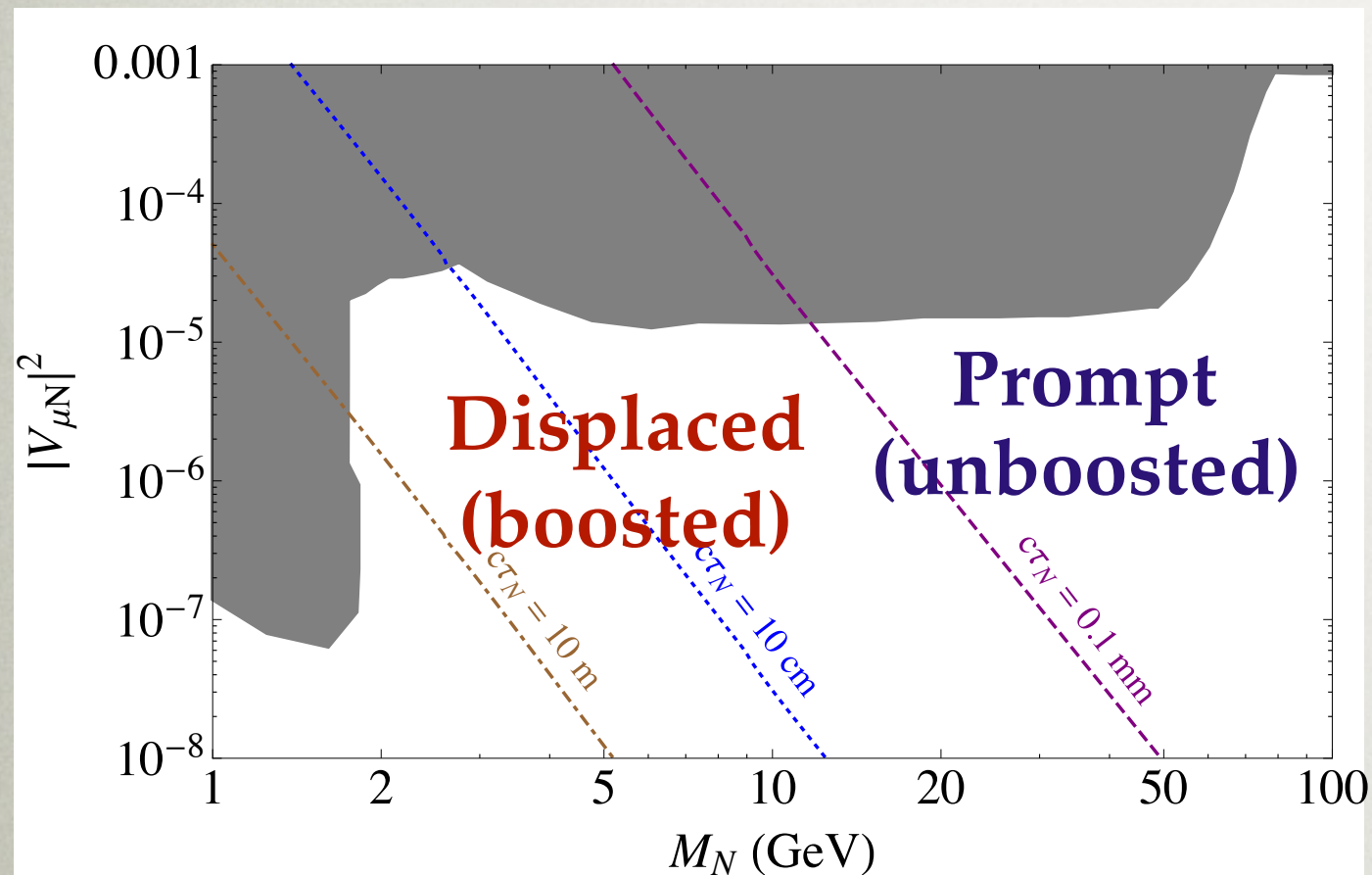
(arXiv:1501.05566)

- We propose cleaner **fully leptonic signatures**

Izaguirre, BS, 1504.02470

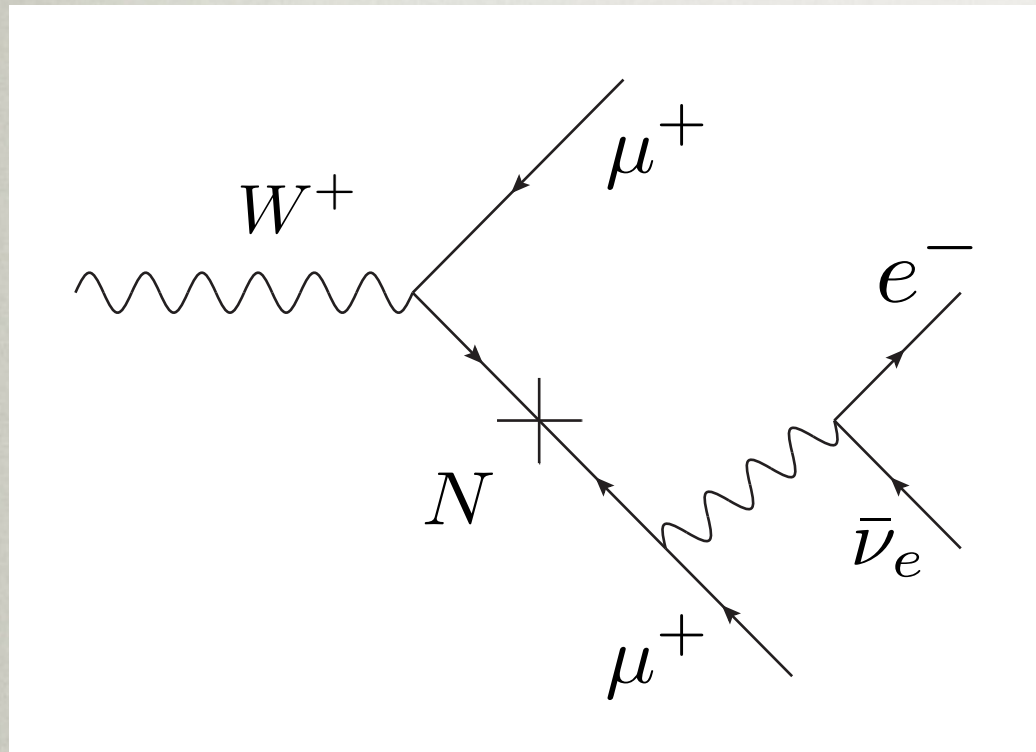
(Proposed for heavy, Dirac N : del Aguila, Aguilar-Saavedra 2008, +de Blas 2009)

Neutrinos at the LHC



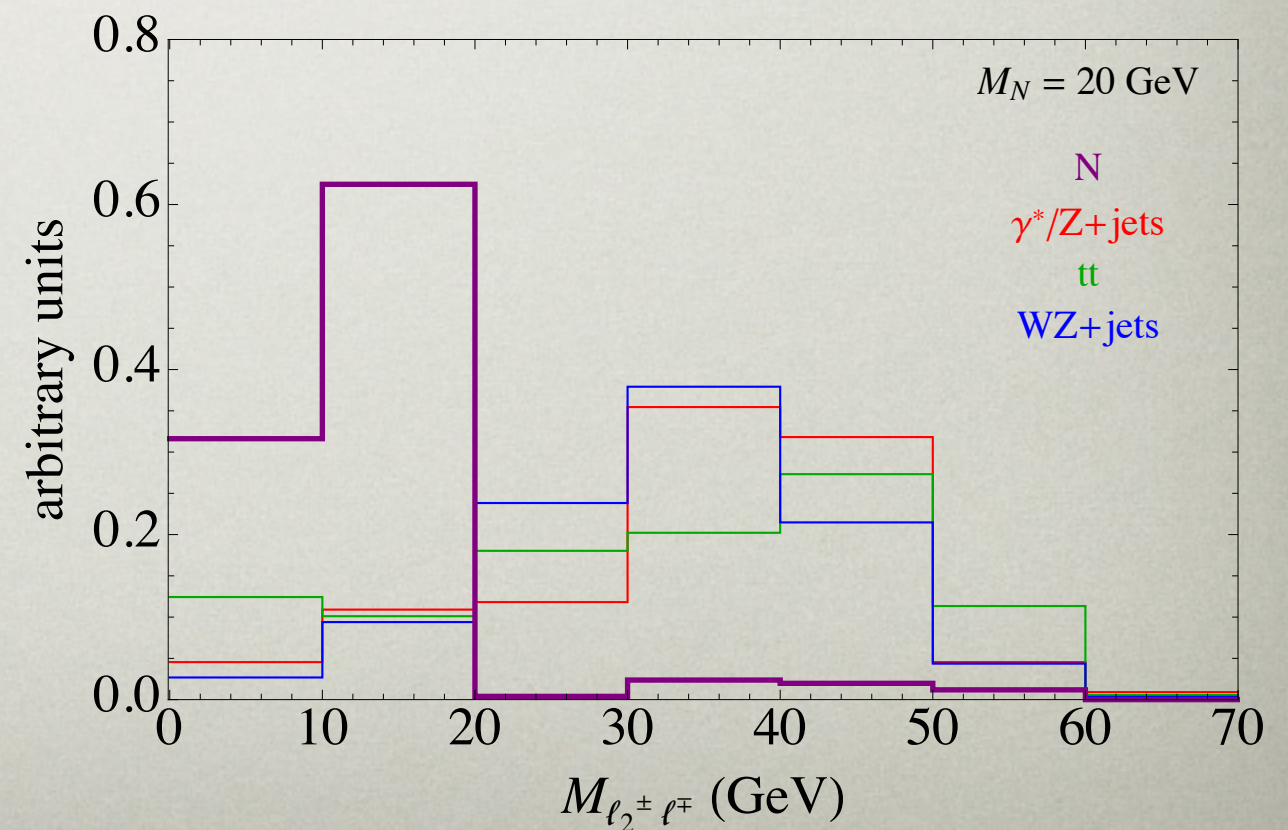
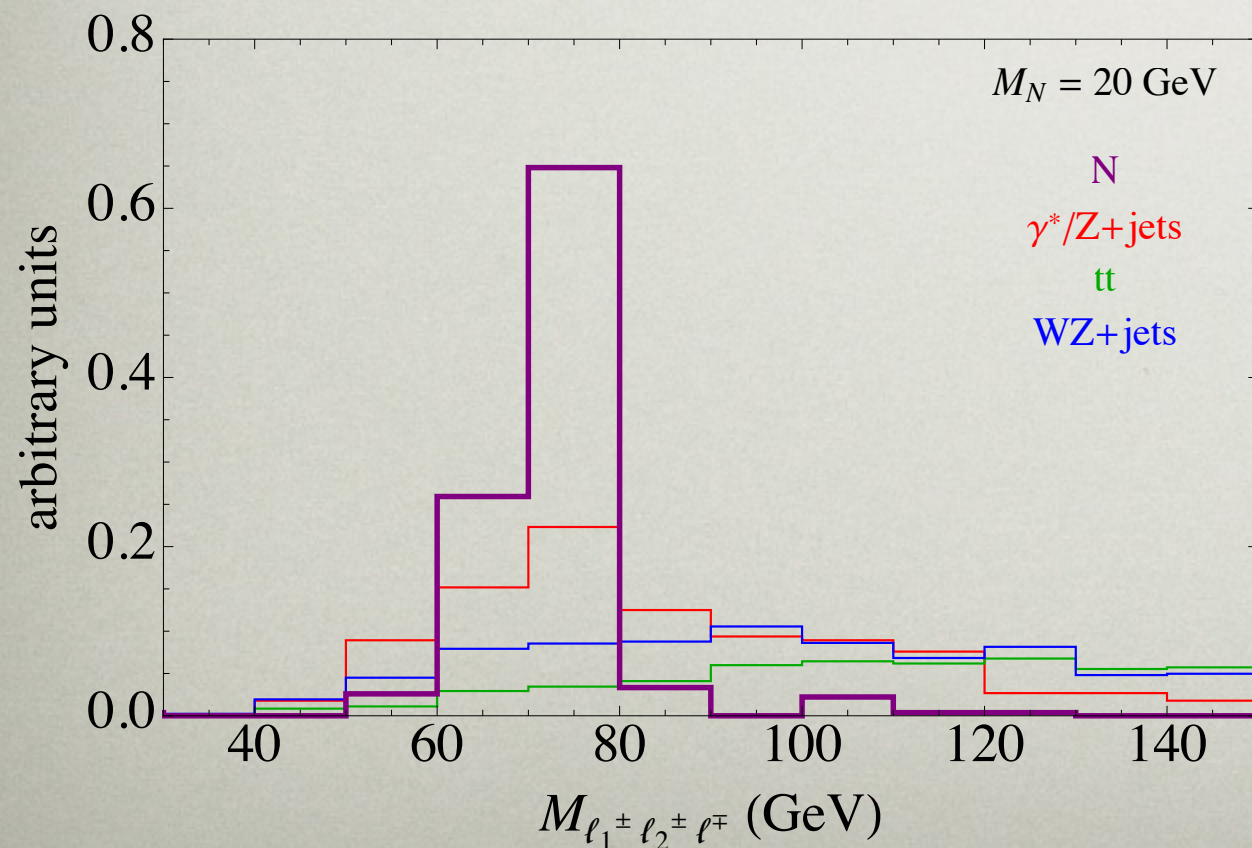
$$\Gamma_N \sim |V_{\alpha N}|^2 G_F^2 M_N^5$$

Prompt/Unboosted Signatures



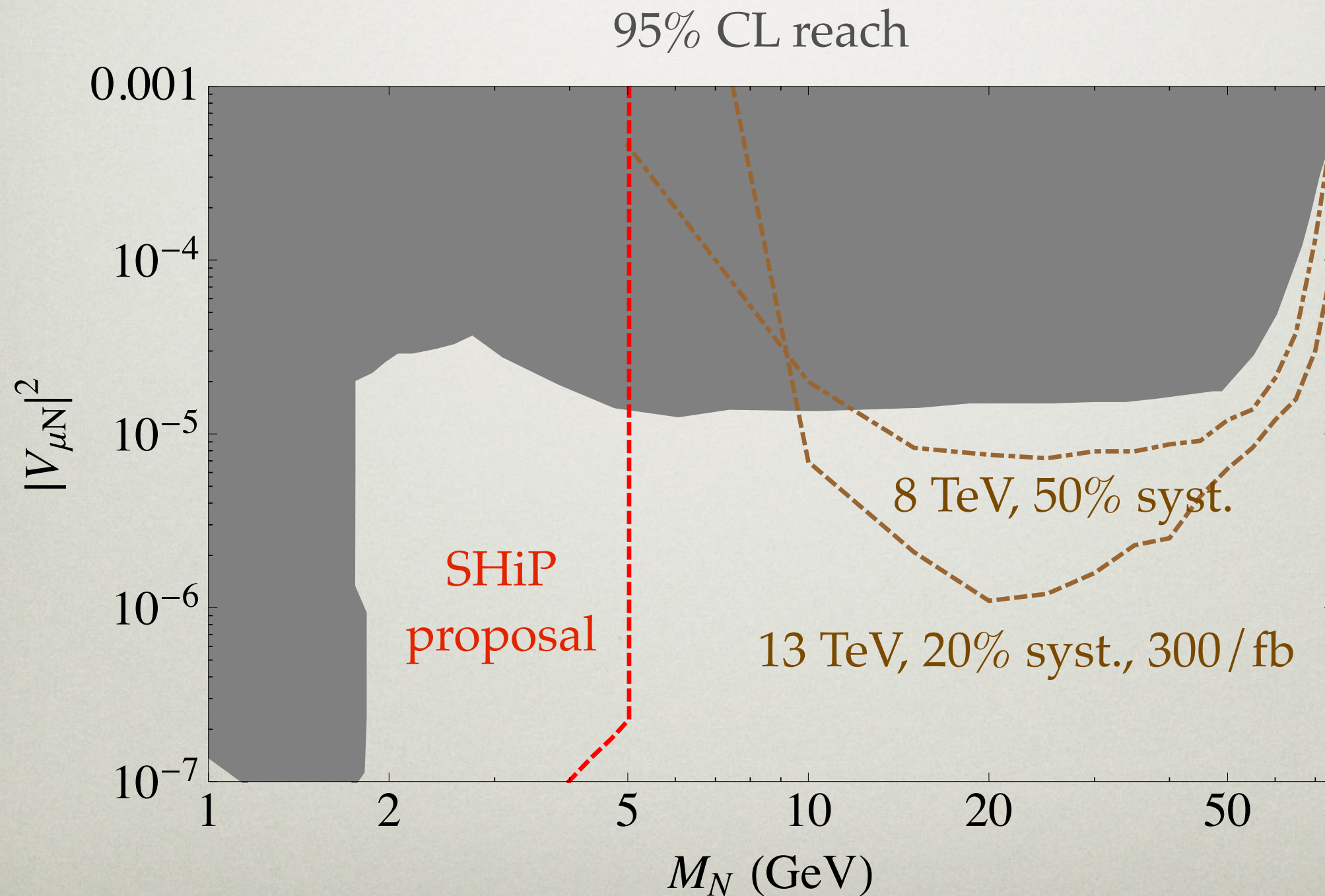
Generic trilepton: large backgrounds from Z , top, diboson

Majorana gives striking OSSF-0 signatures!

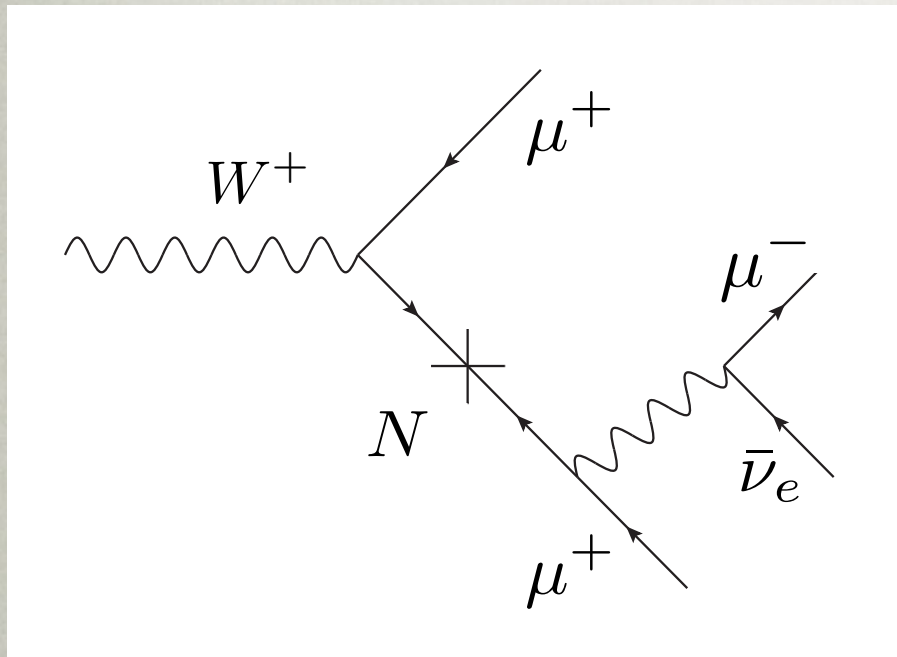


Prompt Neutrino Signatures

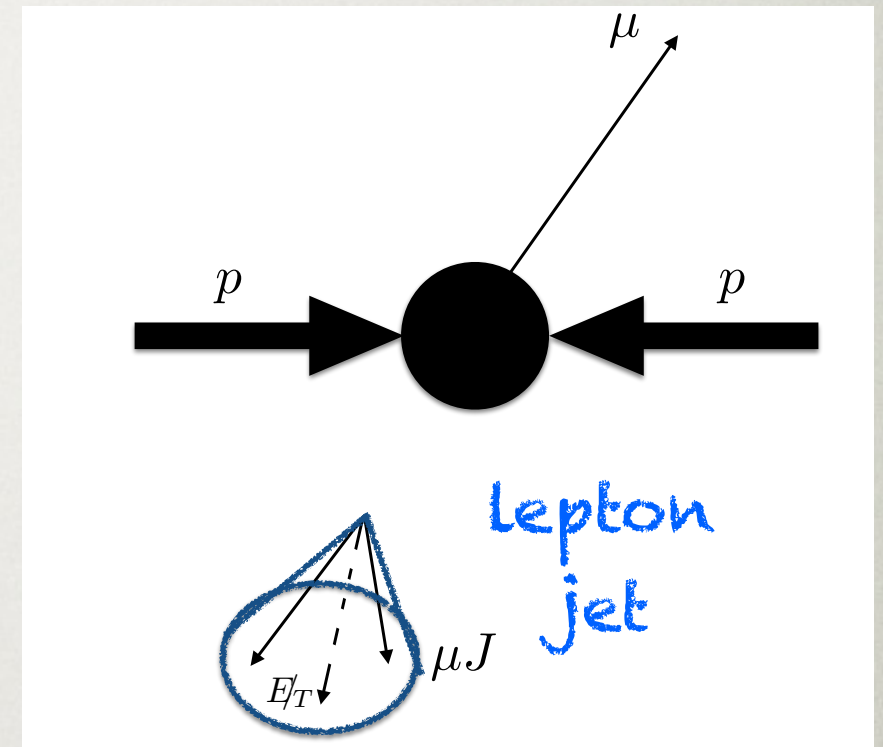
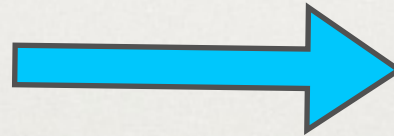
- Require 2 same-sign muons, 1 electron (all isolated, $p_T > 10$ GeV), trigger requirement, low hadronic/missing energy, kinematic cuts



Displaced/Boosted Signatures



$$M_N \ll M_W$$



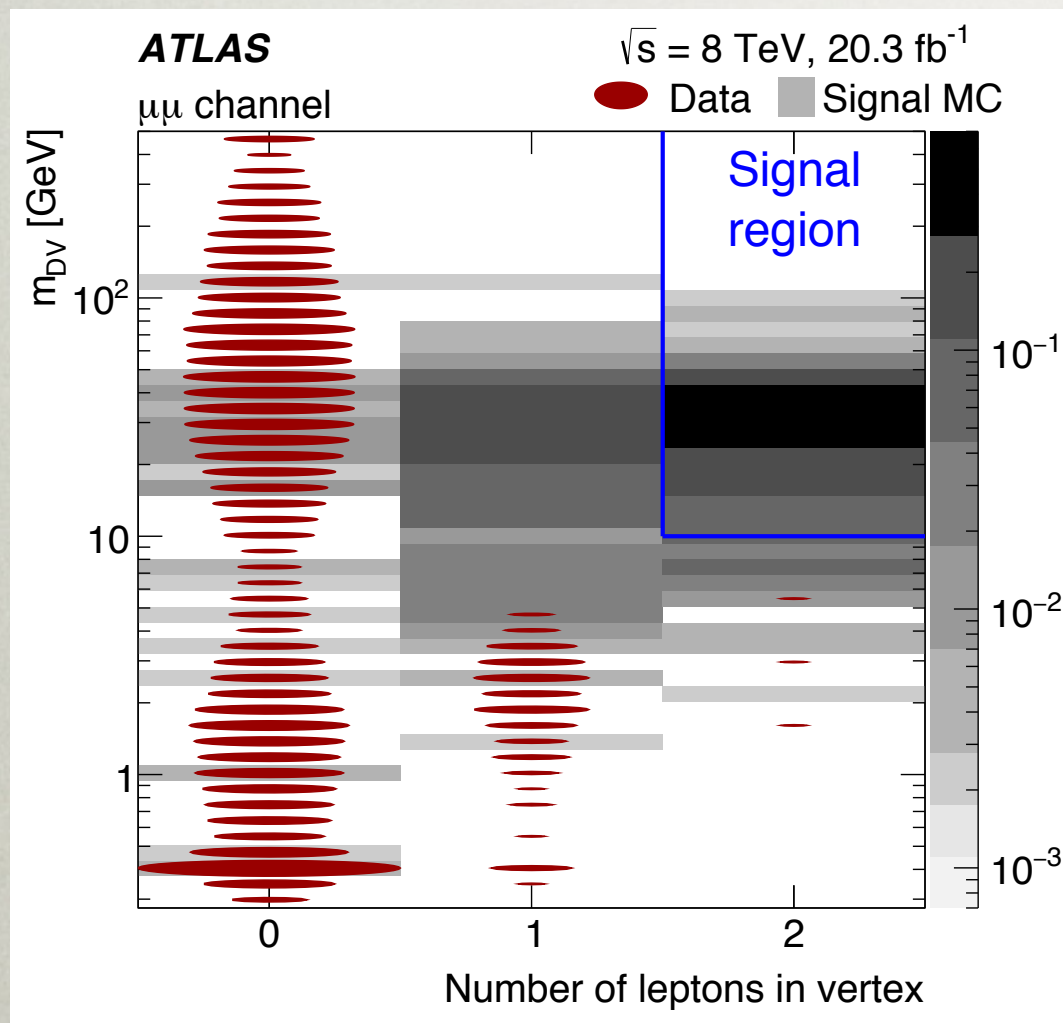
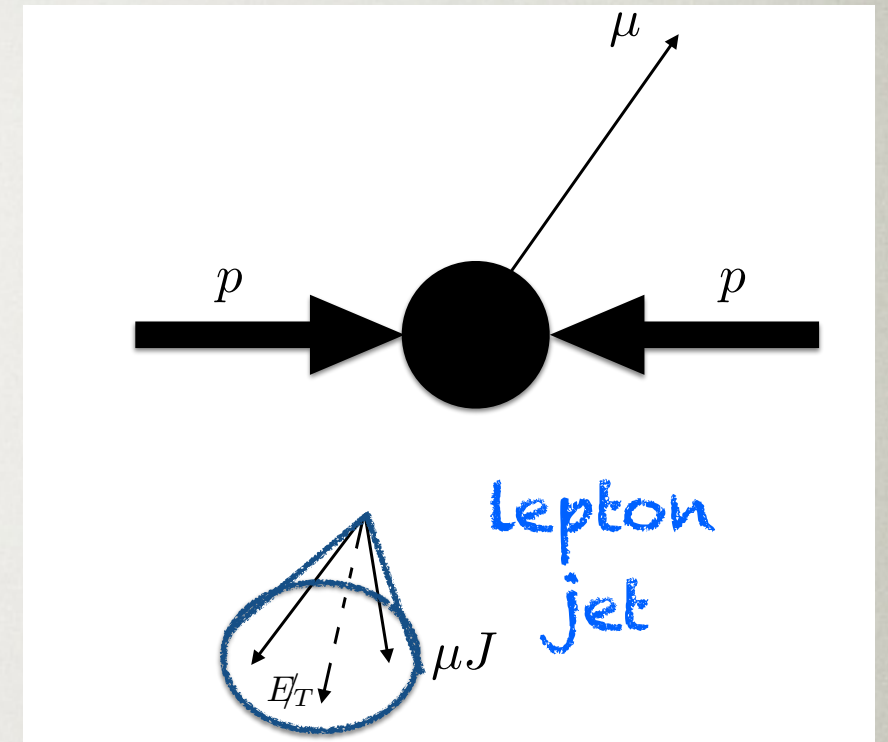
Arkani-Hamed, Weiner, 0810.0714

- Hadronic displaced vertices also possible, but backgrounds could be very large

(Helo, Hirsch, Kovalenko, 1312.2900)

Displaced/Boosted Signatures

- By contrast, leptonic backgrounds expected to be negligible

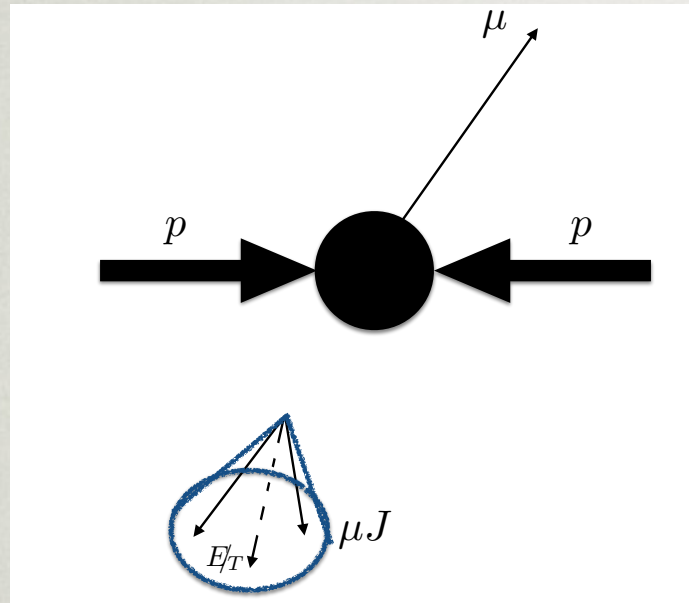


(a)

- Also by extrapolation from existing 2-LJ searches

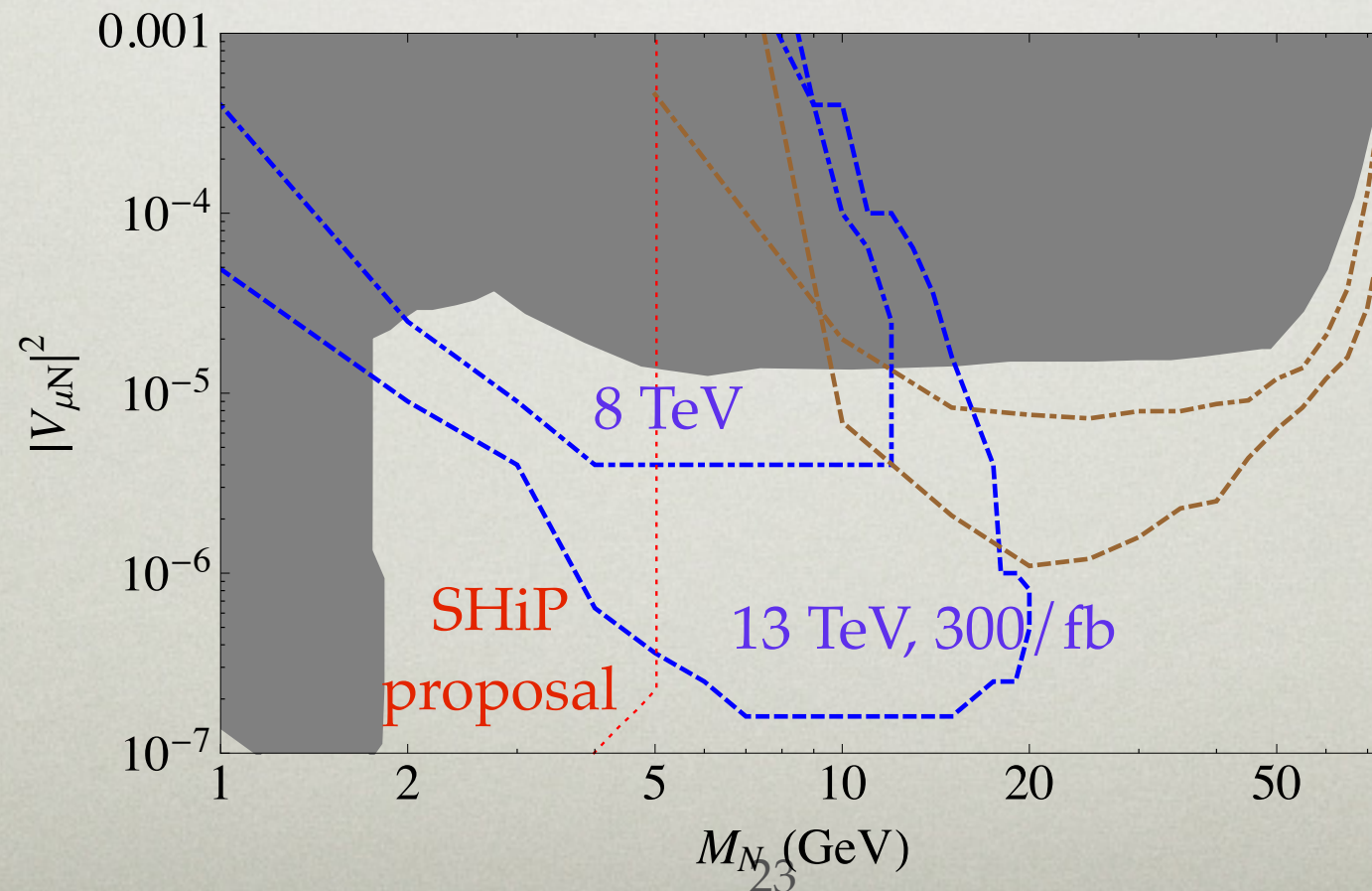
(ATLAS: 1409.0746)

Displaced trilepton signatures



- LJ selections:
 - Hard lepton for trigger, two soft muons in MS
 - Expect zero backgrounds when require a displacement of > 1 mm
 - Veto back-to-back muons

95% CL reach (signal yield ≈ 3)

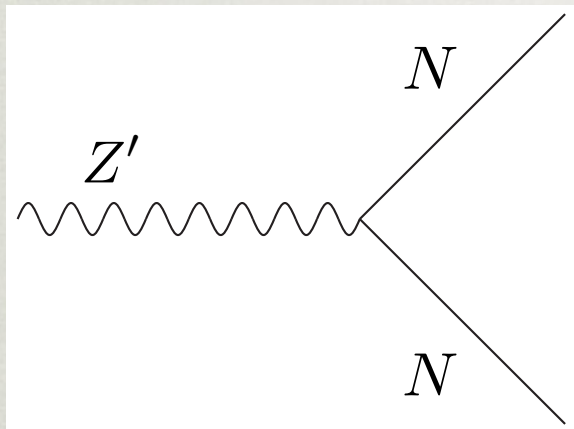


The Z' Portal: Discovery of new B/L forces

B. Batell, M. Pospelov, BS, arXiv:1603.xyzab

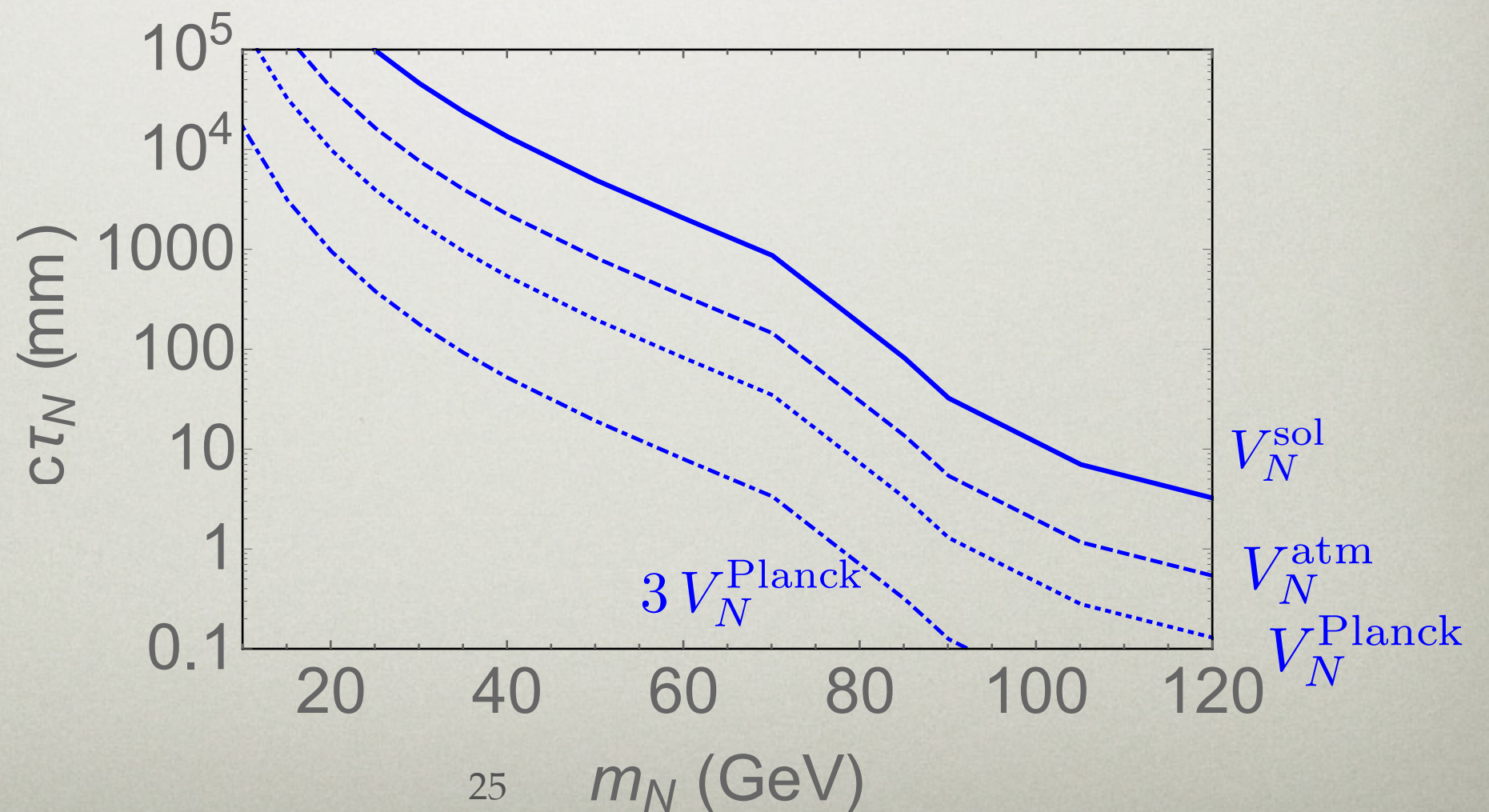
B/L gauge forces at colliders

- Simplest anomaly-free extensions of the SM ($B-L$, $L_i - L_j$, etc.)



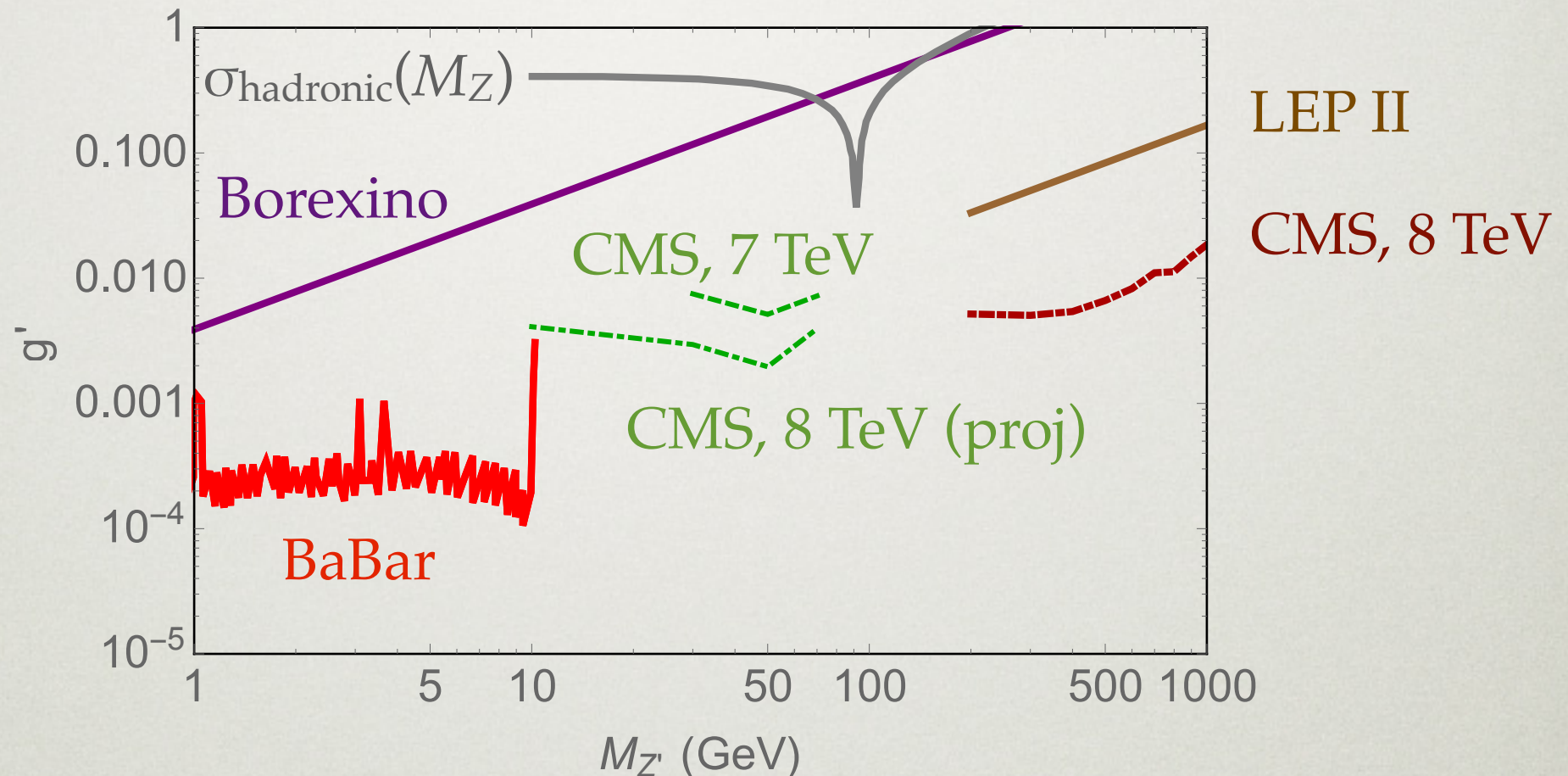
$$V_N \sim \sqrt{\frac{m_\nu}{M_N}}$$

- For $m_\nu \sim 0.1$ eV, $M_N \sim$ GeV, $V_N \sim 10^{-5}$



$B-L$ gauge forces at colliders

- Current constraints on $B-L$:



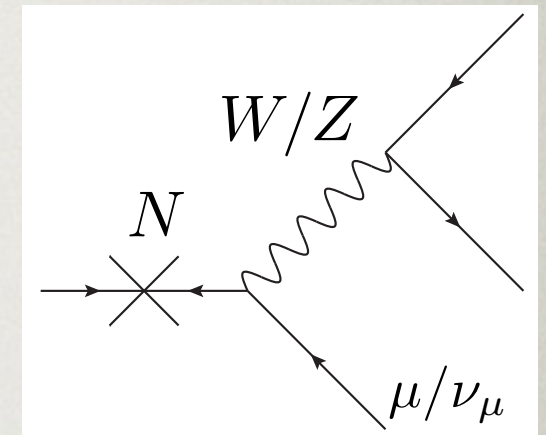
- Drell-Yan rate enormous ($\sim 10^5$ / GeV bin), sensitivity $\sim \sqrt{\mathcal{L}}$

Drell-Yan constraints recast from Hoenig, Samach, Tucker-Smith, 1408.1075

Borexino constraints from Harnik, Kopp, Machado, 1202.6073

B-L gauge forces at colliders

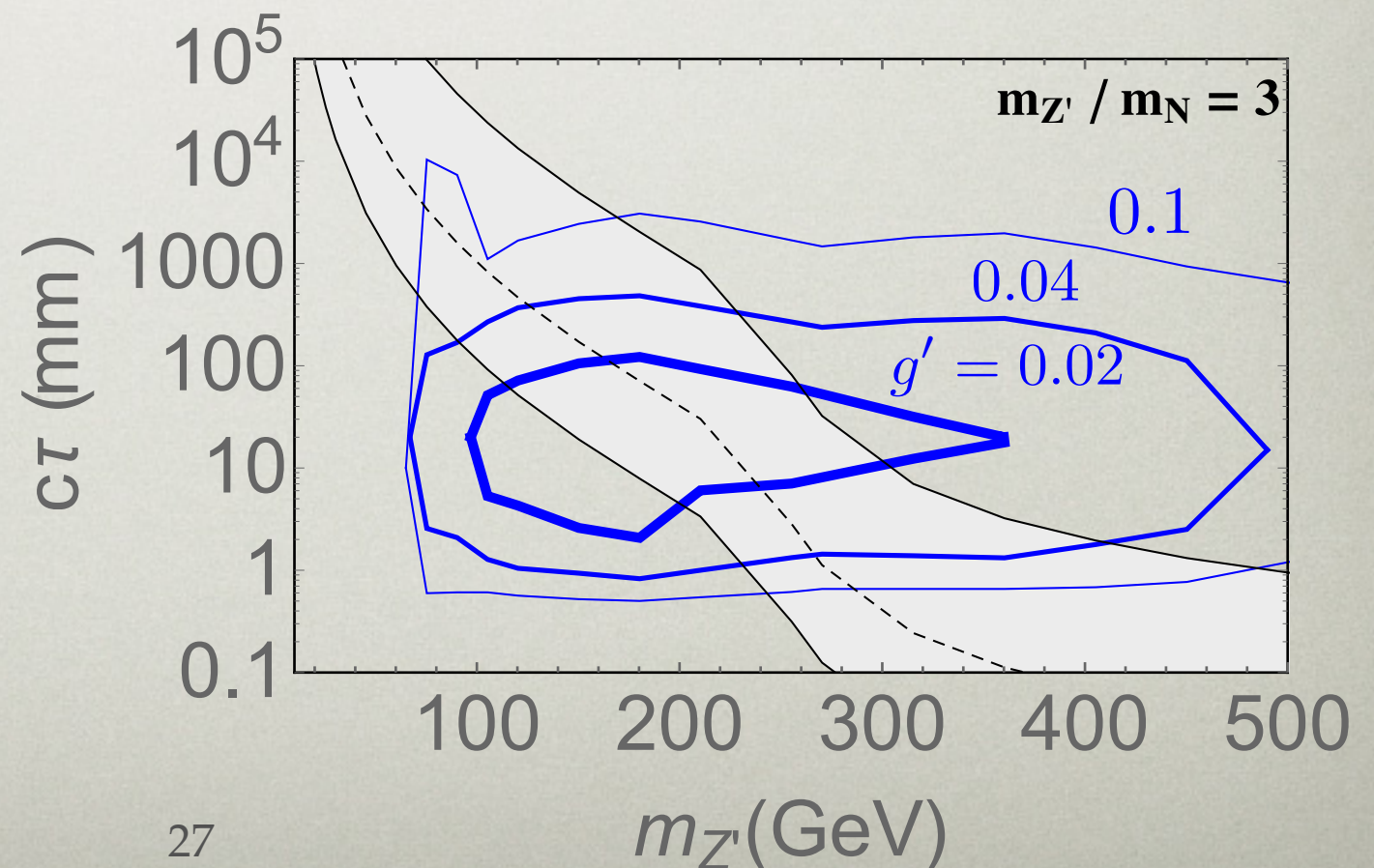
- N decays via (off-shell) W/Z
 - Get displaced muon in $> 50\%$ of decays
 - Use lepton triggers



- Current searches are background-free but have high thresholds or unnecessary restrictions

- CMS: displaced $e + \mu$ on

(1409.4789)

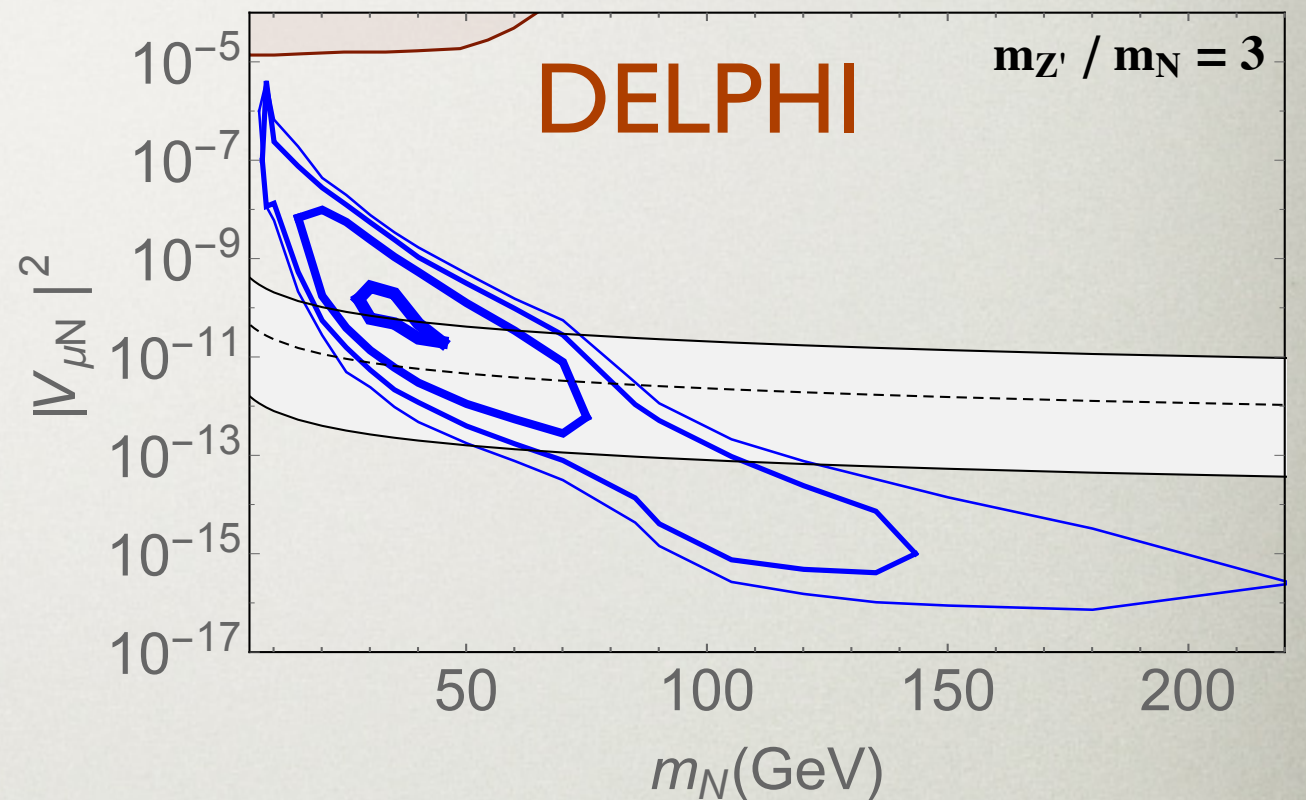
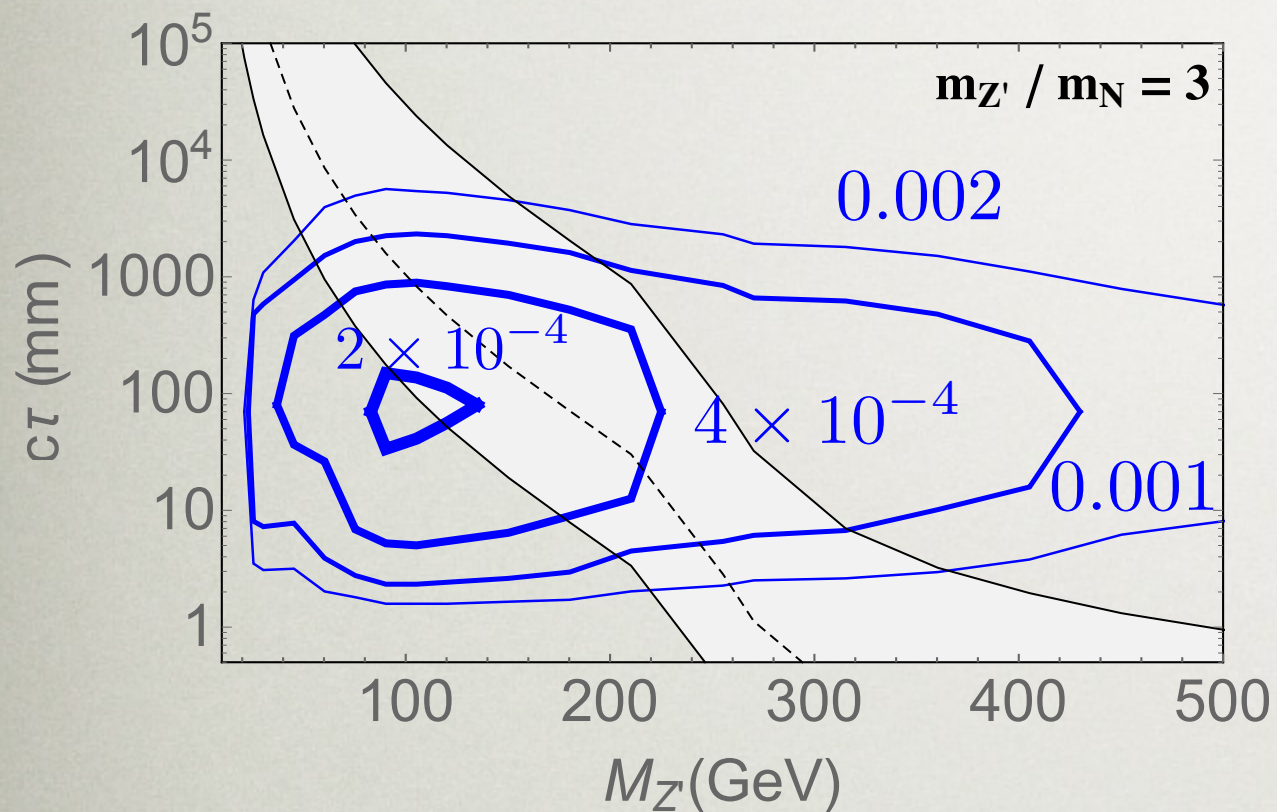


***B-L* gauge forces at colliders**

Our approach: since current searches with 2 displaced leptons or 1 displaced lepton at displaced vertex are bkd-free, we require **both** of these in same event at HL-LHC

- Overly conservative, can relax some kinematic / selection requirements

B-L gauge forces at colliders



- At best, Drell-Yan limits would be:
 - **Below Z:** $g' \lesssim 5 \times 10^{-4}$ ($M_{Z'} \sim 50-60$ GeV)
 - **Above Z:** $g' \lesssim 10^{-3}$ ($M_{Z'} \sim 150-400$ GeV)
- SHiP can reach $g' \sim 10^{-5}$ ($M_{Z'} \sim 1-10$ GeV)!!

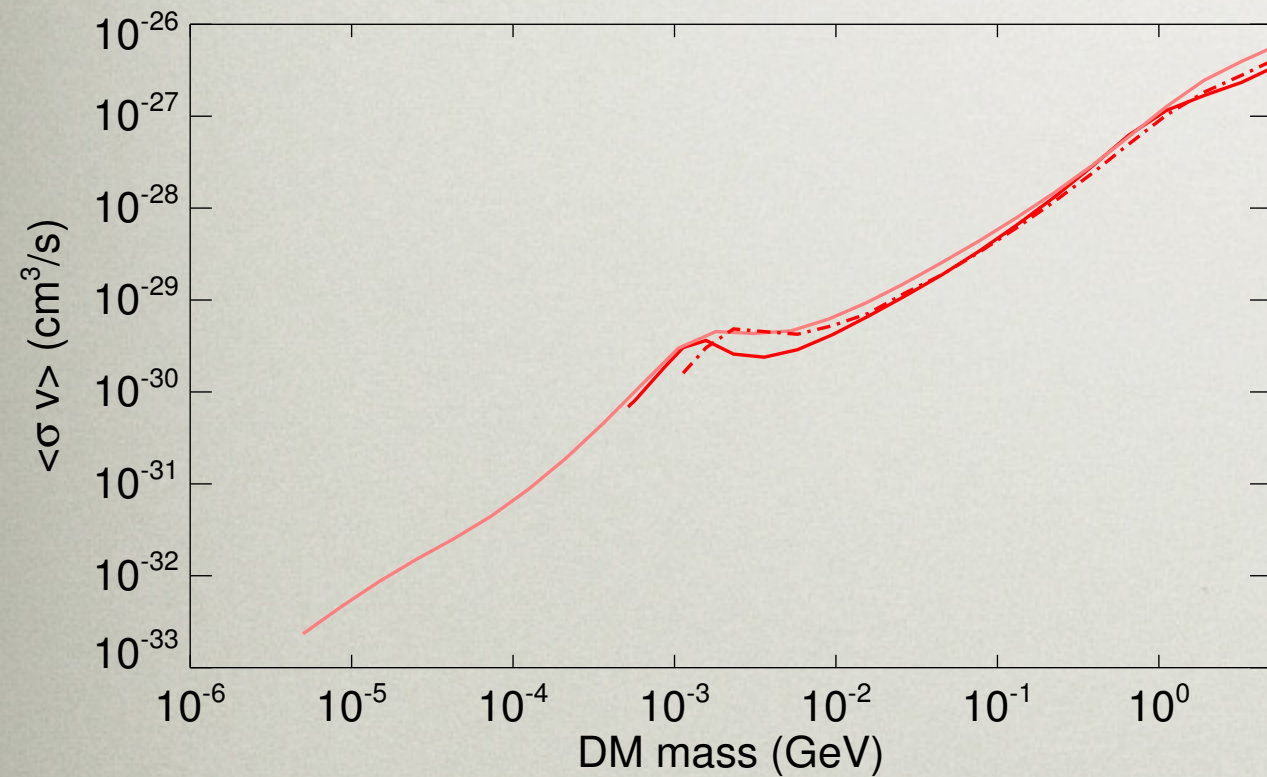
The Vector Portal & Inelastic Dark Matter

E. Izaguirre, G. Krnjaic, BS, arXiv:1508.03050

Discovering Dark Matter

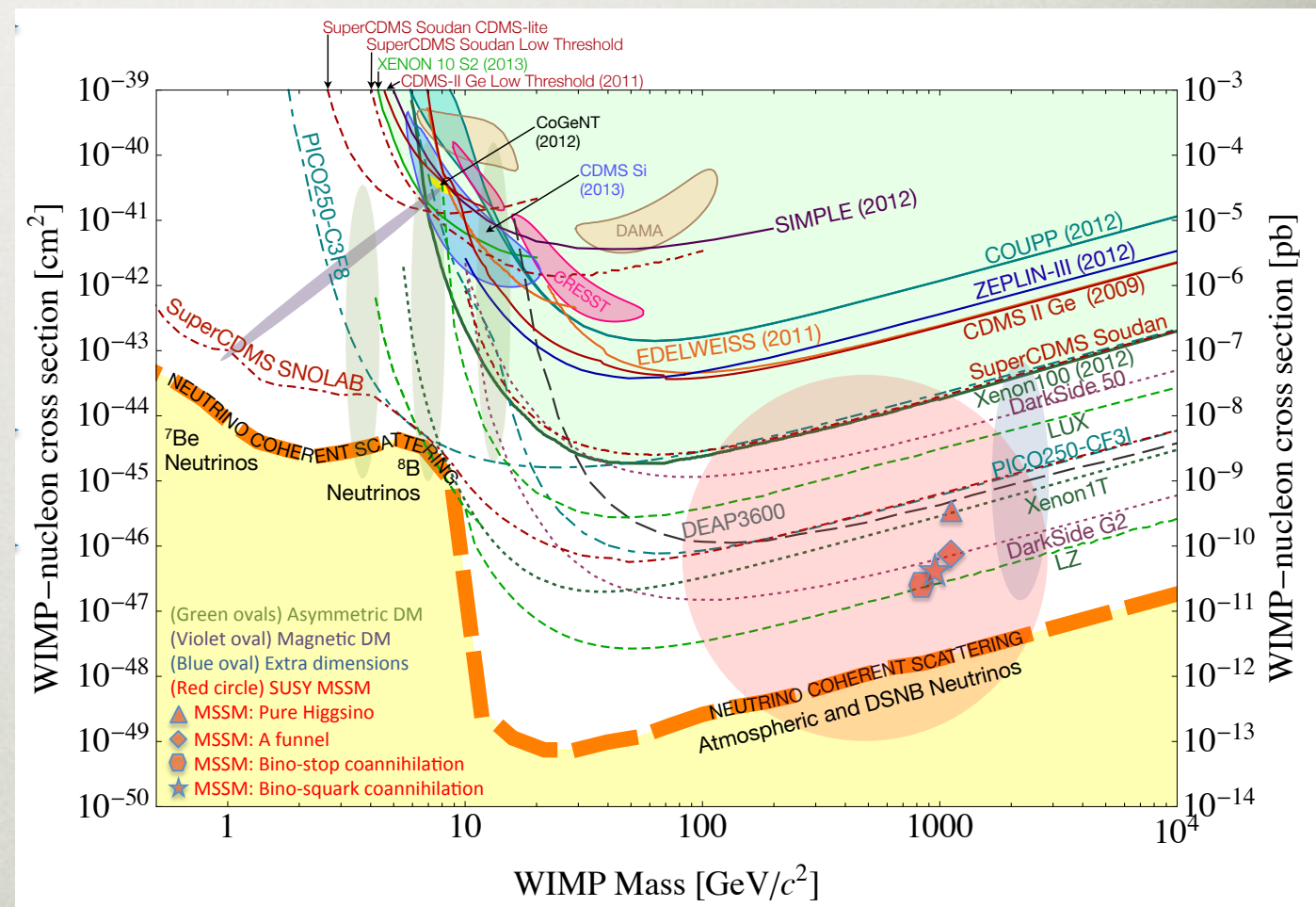
- Generically strong constraints on low-mass dark matter

CMB/indirect detection



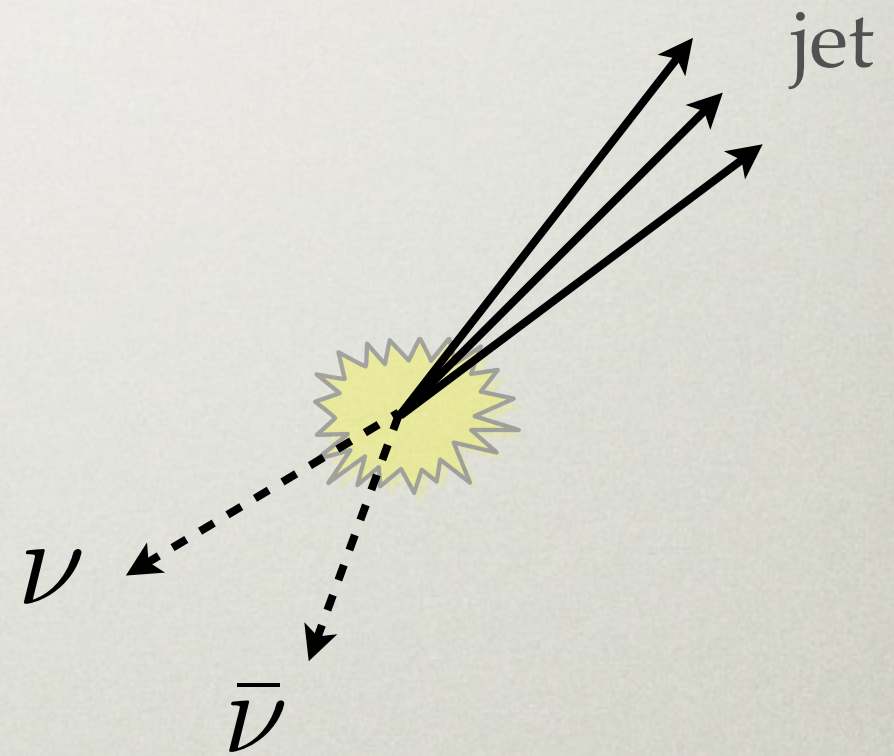
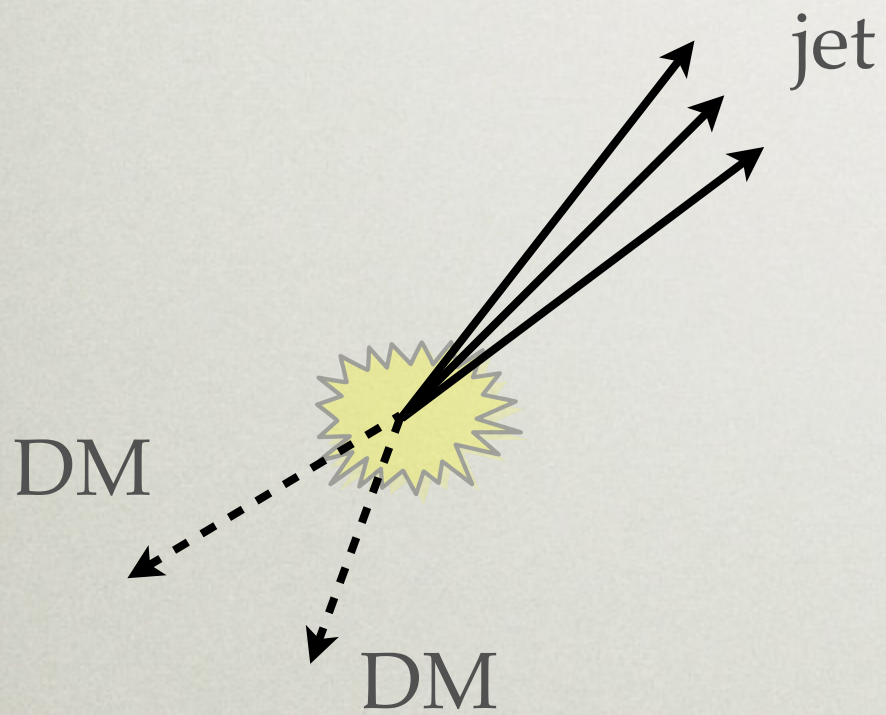
from Slatyer, 1506.03811

direct detection



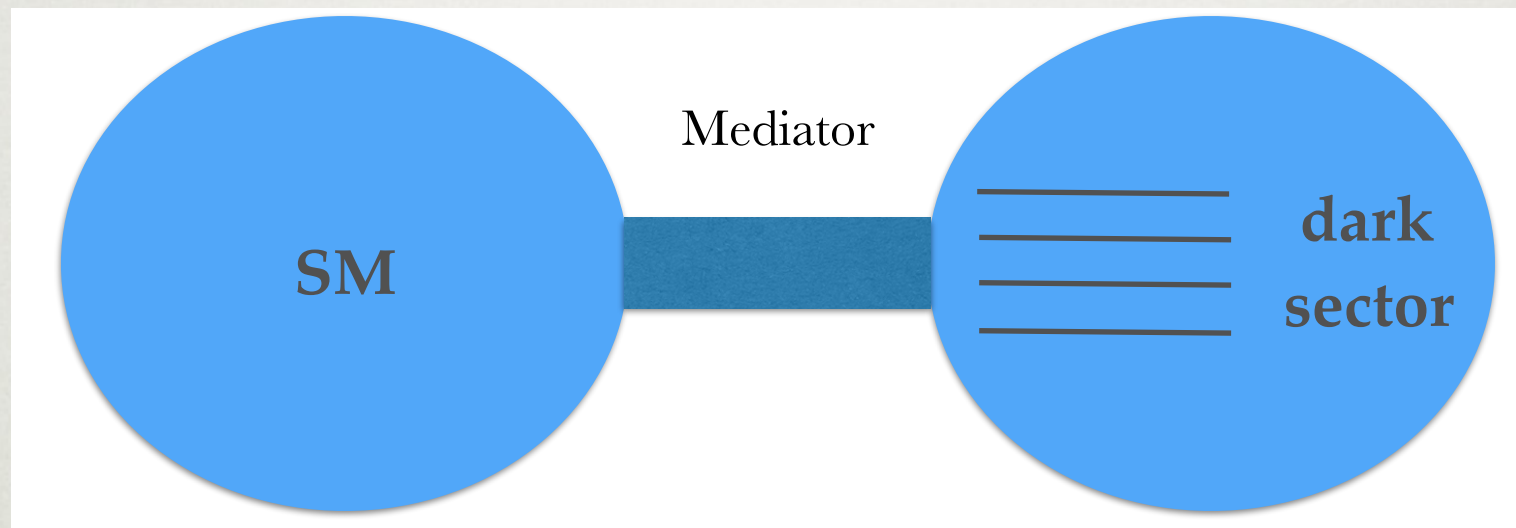
Discovering Dark Matter

- By contrast, high-energy colliders can be tricky...



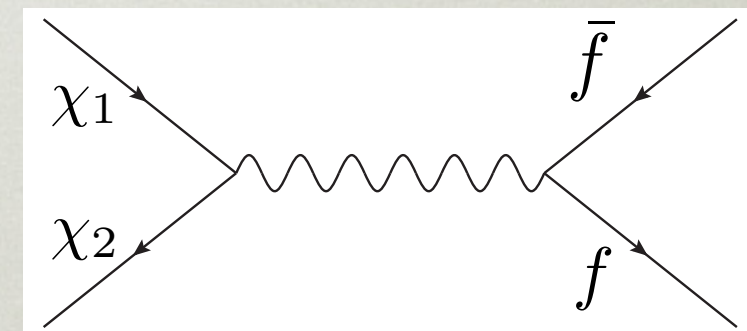
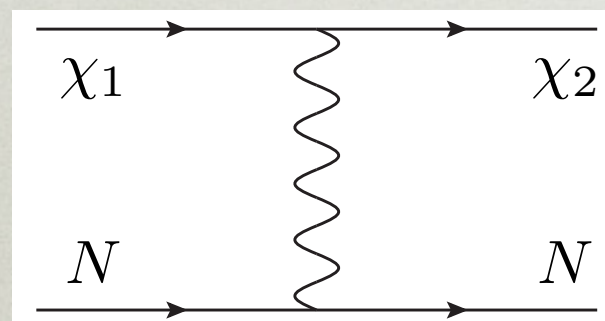
Dark Matter in a Dark Sector

- In a dark sector, the situation can be flipped



- Ex: inelastic dark matter

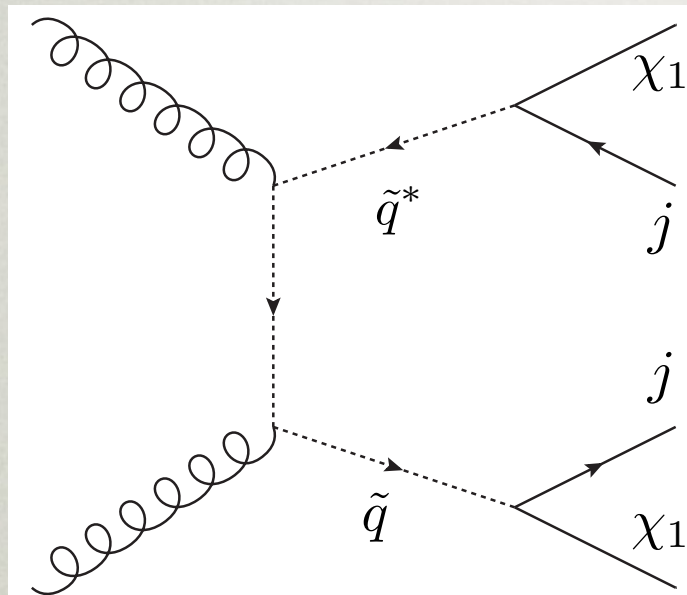
(Tucker-Smith, Weiner, hep-ph/0101138)



$$E_{\text{collider}} \gg \Delta M \gg E_{\text{DM}}^{\text{kin}}$$

Dark Matter in a Dark Sector

- For large splittings, decays of heavier DS states give **hard** objects



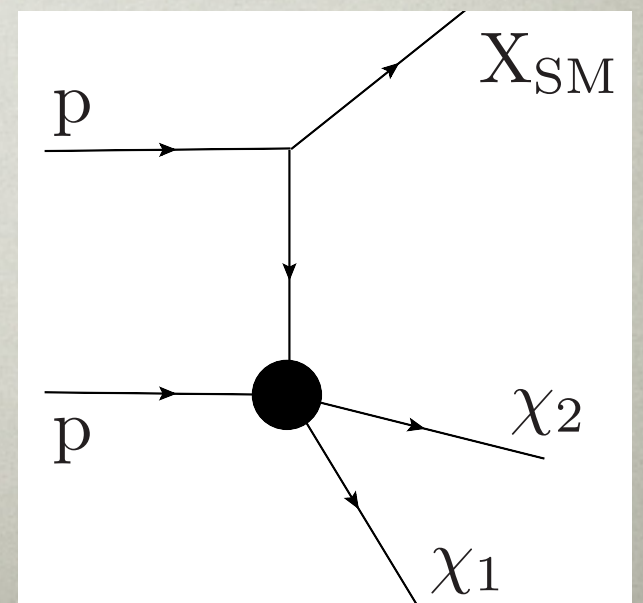
(Chang *et al.*, 1307.8120; An *et al.*, 1308.0592; Bai, Berger, 1308.0612; Papucci, Vicchi, Zurek, 1402.2285; Primulando, Salvioni, Tsai, 1503.04204)

- For small splittings, get conventional monojet **plus additional soft radiation** from heavier DS decay

(Bai, Tait, 1109.4144; Izaguirre, Krnjaic, BS, 1508.03050)

- Similar to compressed SUSY searches

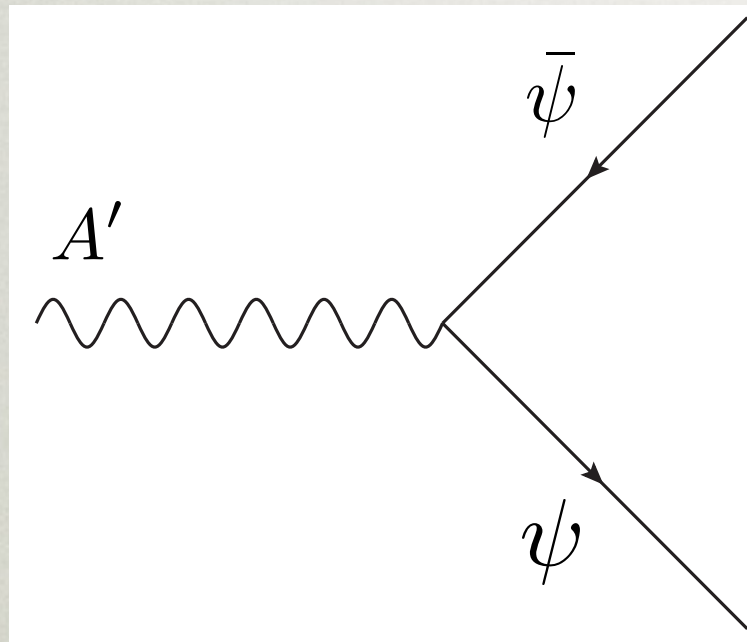
Recent examples: Giudice *et al.*, 1004.4902; Gori *et al.*, 1307.5952; Buckley *et al.*, 1310.4827, and many others...



An Inelastic Benchmark

(Izaguirre, Krnjaic, BS, 1508.03050)

- Dark matter with Higgsed dark QED



$$\mathcal{L}_\psi \supset -M_\psi \bar{\psi} \psi + \frac{y \langle \Phi \rangle}{2} \bar{\psi}^c \psi + \text{h.c.}$$

$$M_\psi \gg y \langle \Phi \rangle$$

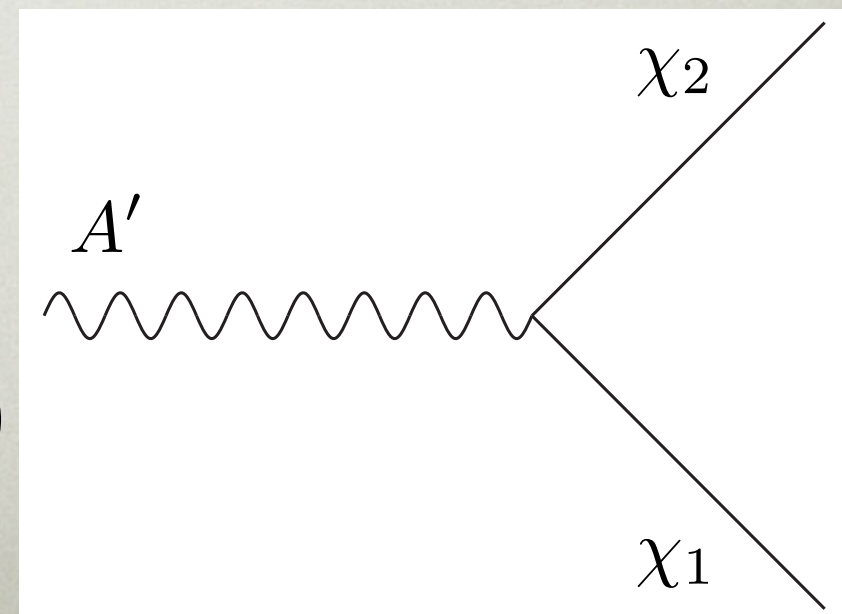


$$M_{1,2} \approx M_\psi \pm \Delta/2$$

$$\Delta \equiv 2y \langle \Phi \rangle$$

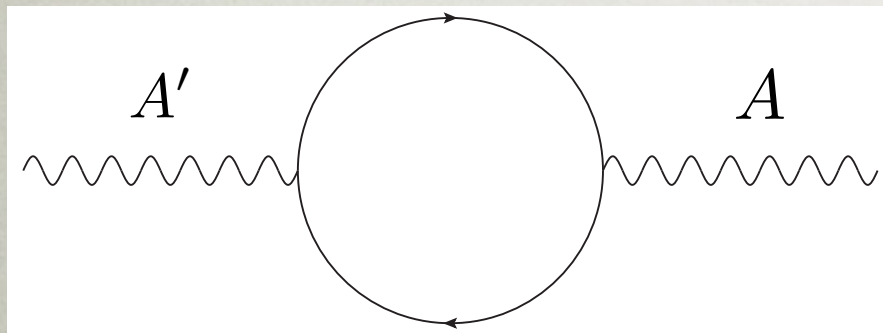
- Parity conservation = off-diagonal

$$\mathcal{J}_\chi^\mu = i \left(\chi_2^\dagger \bar{\sigma}^\mu \chi_1 - \chi_1^\dagger \bar{\sigma}^\mu \chi_2 \right) + \mathcal{O}(\Delta/M)$$



An Inelastic Benchmark

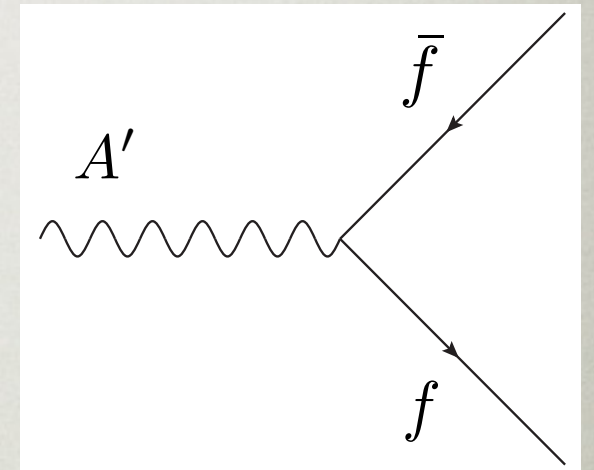
- Coupling to SM comes from kinetic mixing



$$\mathcal{L}_{\text{mix}} = -\frac{\epsilon}{2} F_{\mu\nu} F'^{\mu\nu}$$



$$\mathcal{L}_{A'-\text{SM}} \approx e\epsilon J_{\text{SM}}^{\mu} A'_{\mu}$$

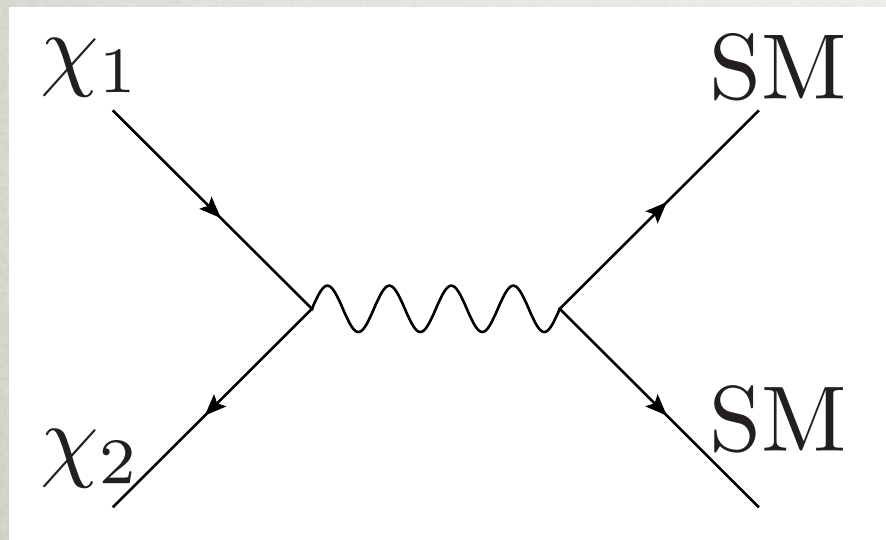


- We consider the spectrum hierarchy $\Delta \ll M_{1,2} \lesssim M_{A'}$

See also: “Secluded DM”, Pospelov, Ritz, Voloshin, 0711.4866;
 Autran *et al.*, 1504.01386; Bai *et al.*, 1504.01395; Buschmann *et al.*, 1505.07549

Inelastic Freeze-out

- Many parameters -- choose a parameterization that connects freeze-out to lab probes



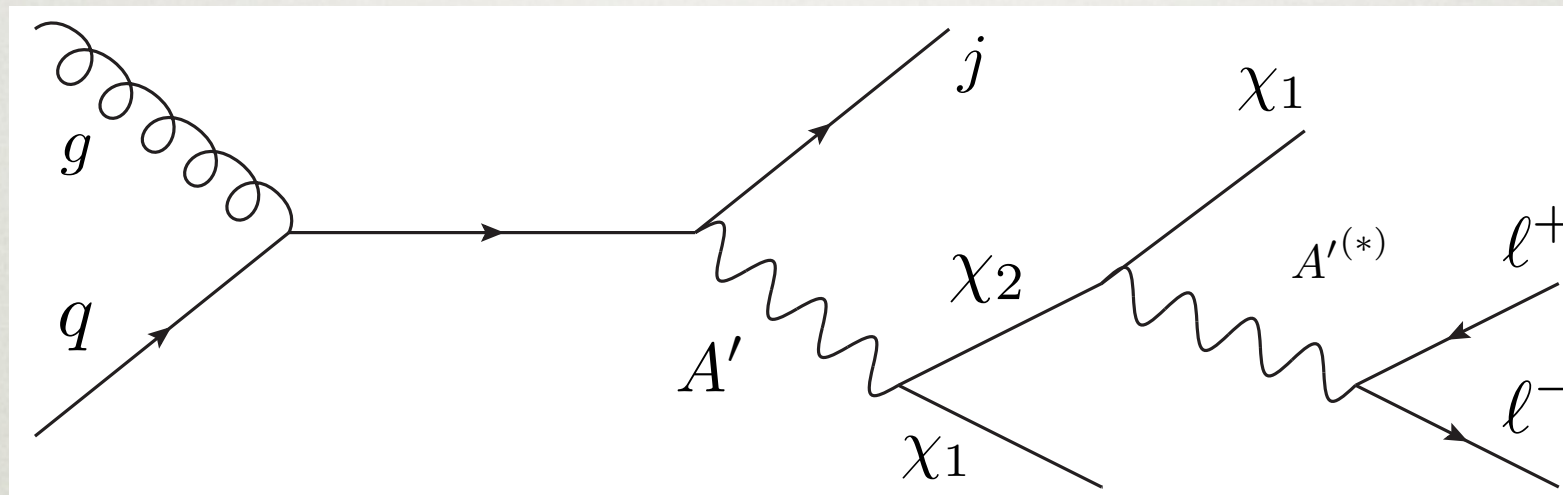
$$\langle \sigma v \rangle \propto \frac{\epsilon^2 \alpha_D M_1^2}{M_{A'}^4} = \frac{y}{M_1^2} \quad (M_{A'} \gg M_1)$$

$$y \equiv \epsilon^2 \alpha_D \left(\frac{M_1}{M_{A'}} \right)^4$$

- Abundance depends on y, M_1 , while lab constraints depend on $\epsilon, M_{A'}$
 - Choose large value of α_D to avoid over-stating bounds
(Izaguirre *et al.*, 1505.00011)

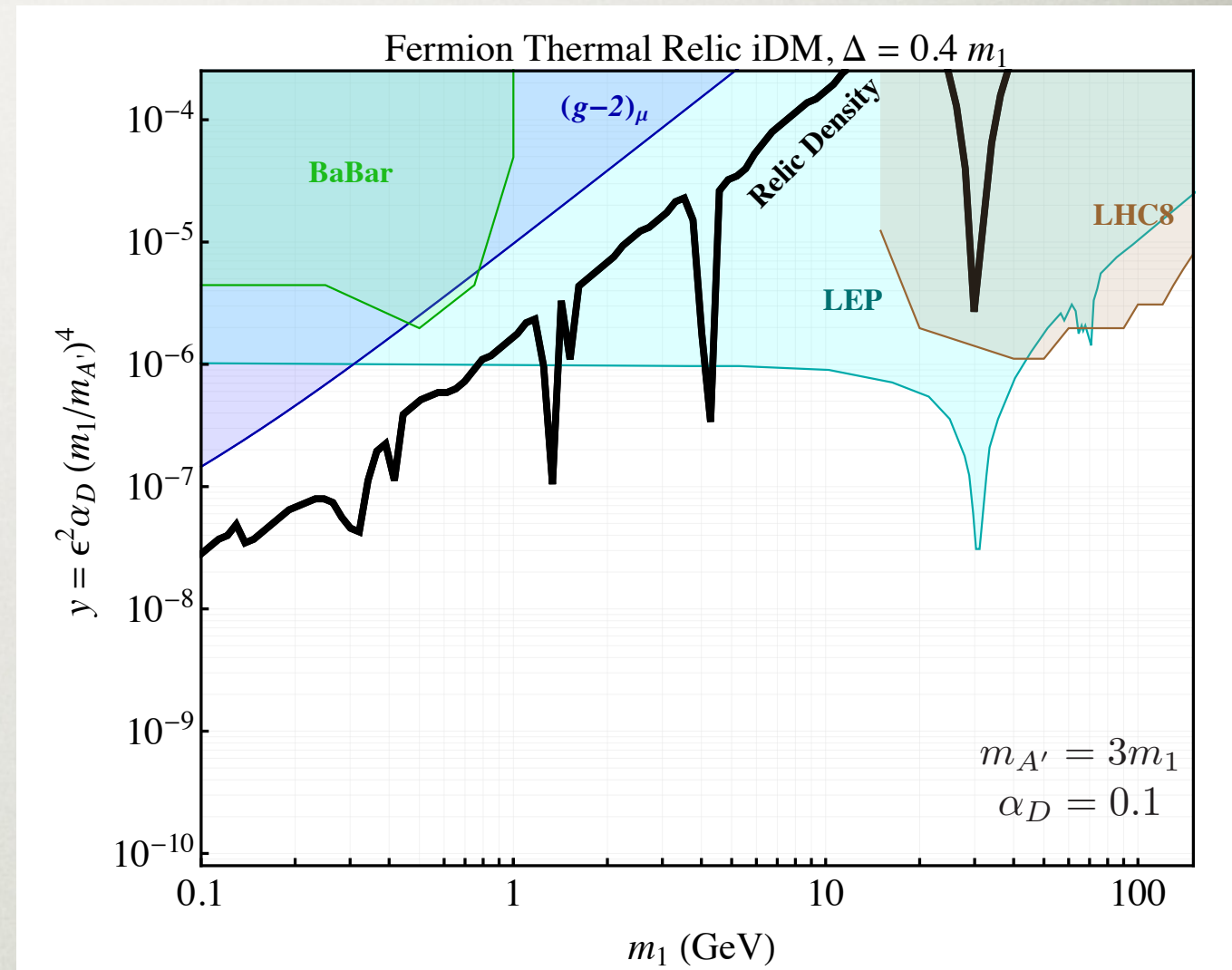
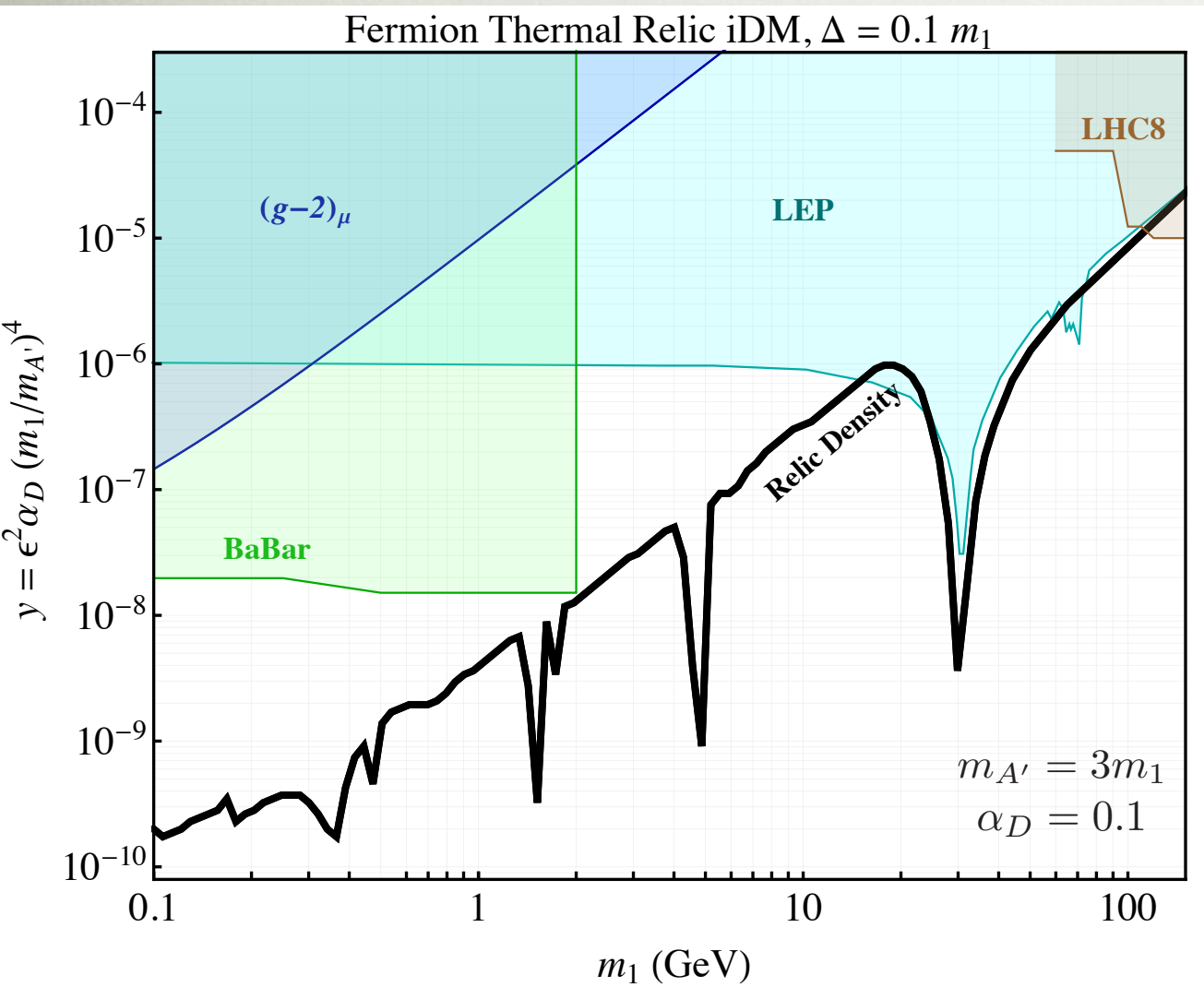
Inelastic Freeze-out

- Collider production:

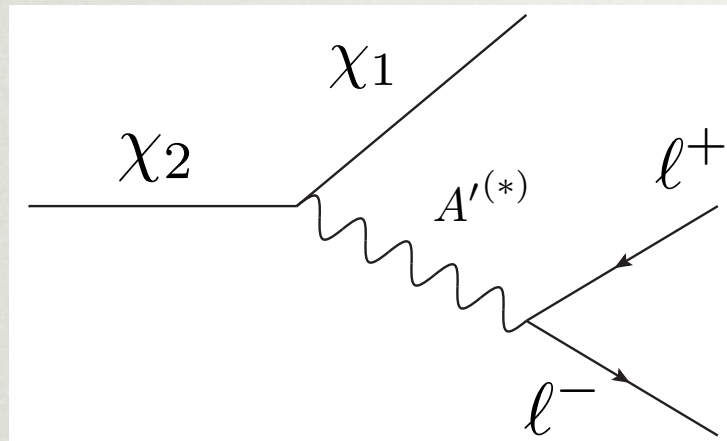


- Also have bounds from EWPT, monophoton, compressed SUSY

Current Status



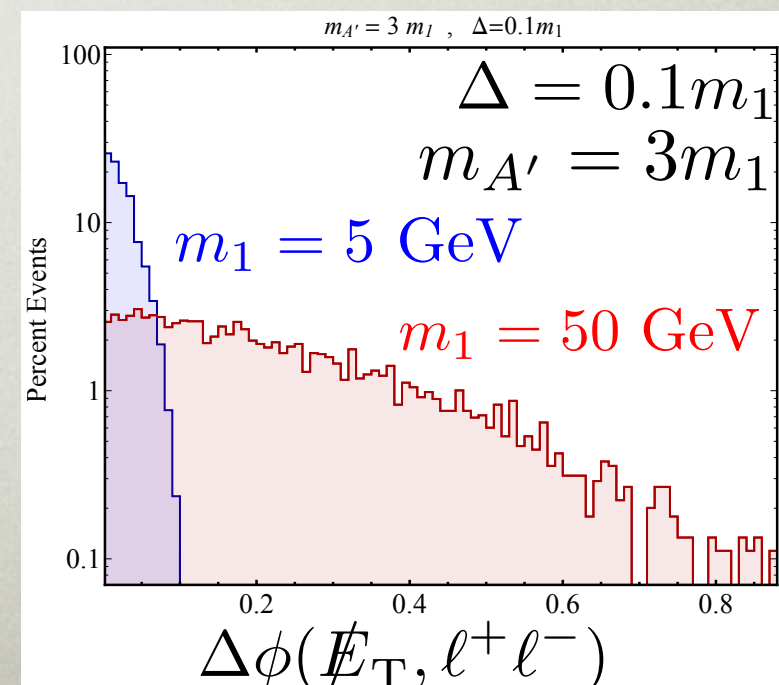
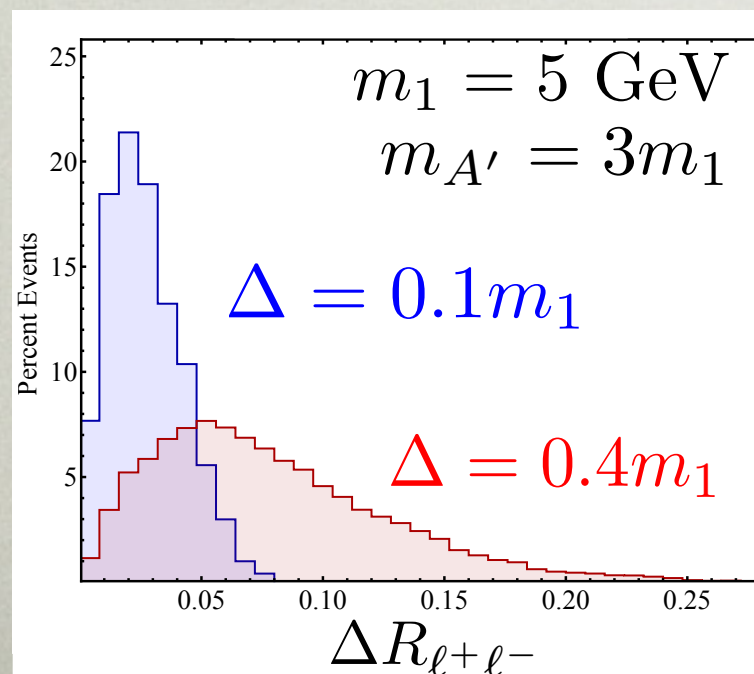
Improving the Searches



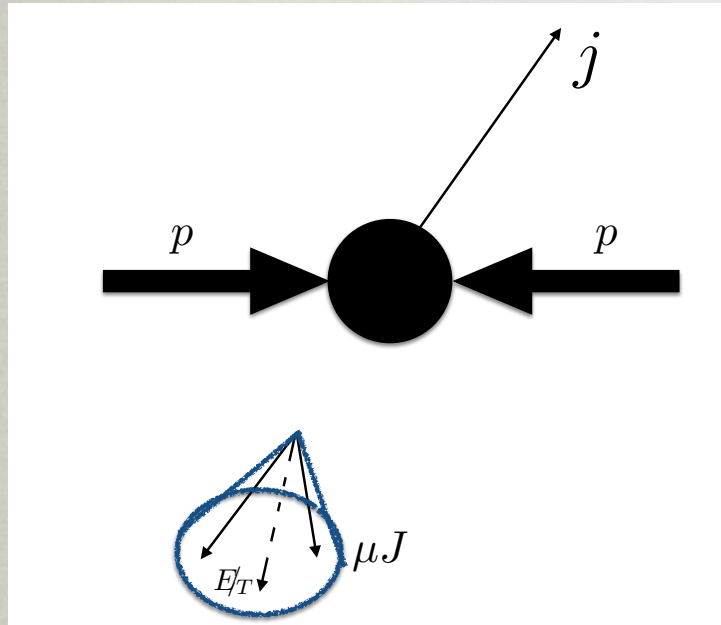
$$\Gamma_{\chi_2} \sim \frac{\alpha \alpha_D \epsilon^2 \Delta^5}{M_{A'}^4}$$

- Get displaced decay!

- The leptons are typically soft, so trigger on monojet + MET
- The DM produced through on-shell A' , so typically **boosted**



Improving the Searches

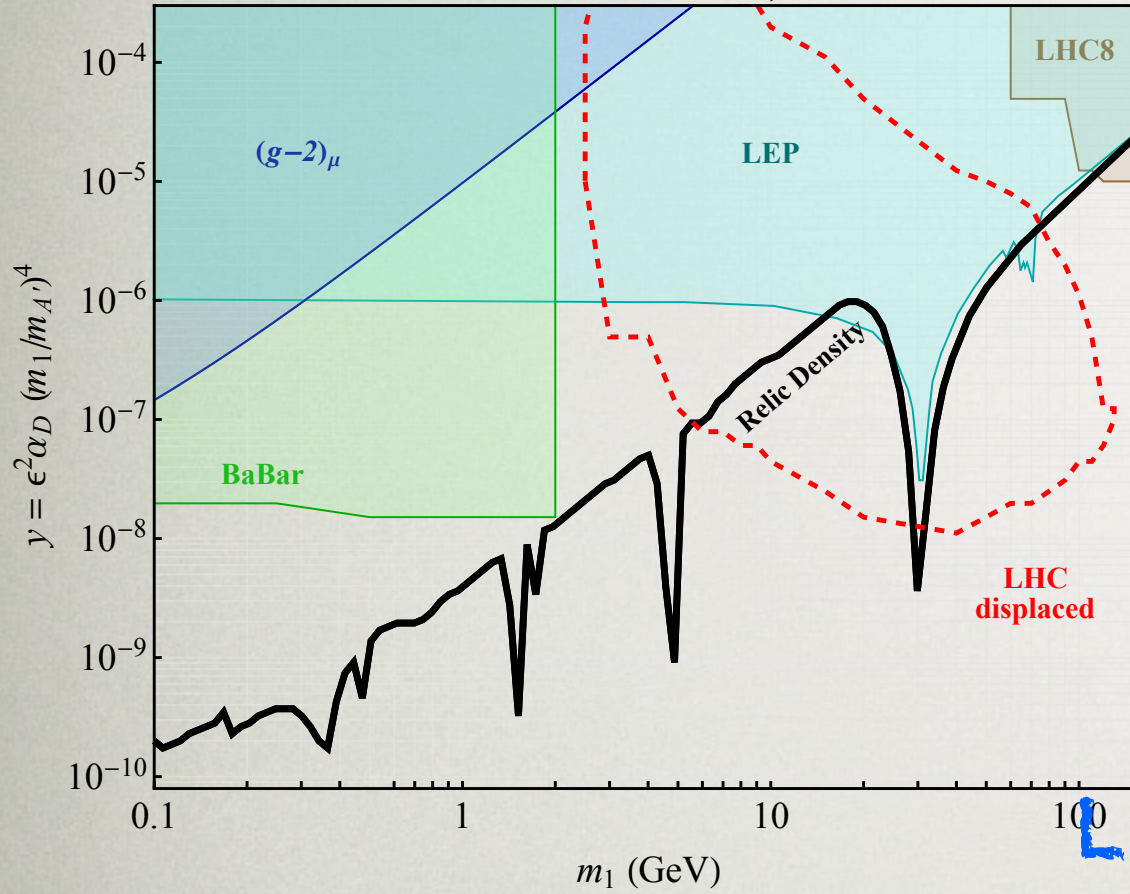


- Monojet + soft displaced lepton jet + MET
- Could be background free:
plot sensitivity for **ten** signal events

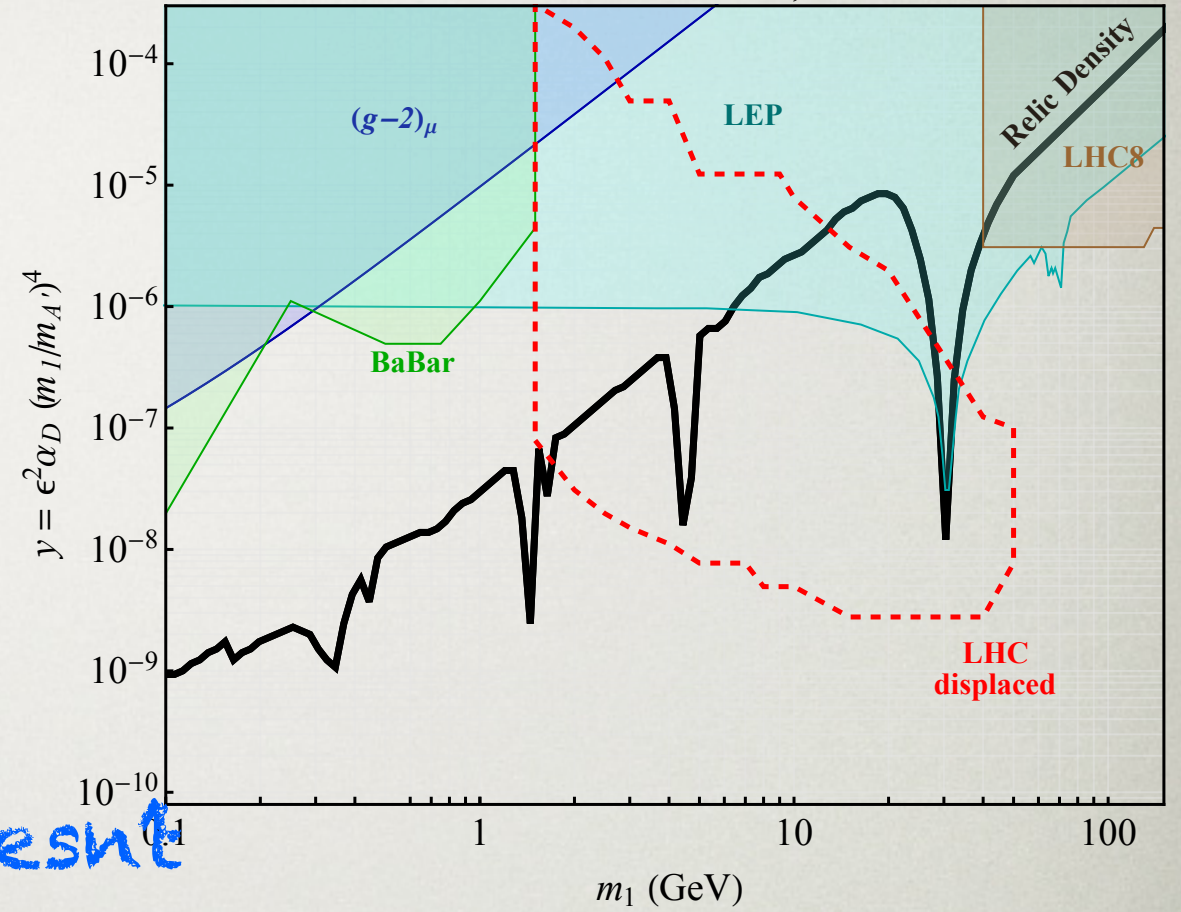
- Require $\cancel{H}_T > 120$ GeV
- Leading jet $p_T > 120$ GeV, veto 3rd jet $p_T > 30$ GeV
- Two displaced muon tracks, $p_T > 5$ GeV, crossing within 1 mm of one another
- $\Delta R < 0.4$ between muons
- $|\Delta\phi| < 0.4$ between lepton jet and MET

LHC Results

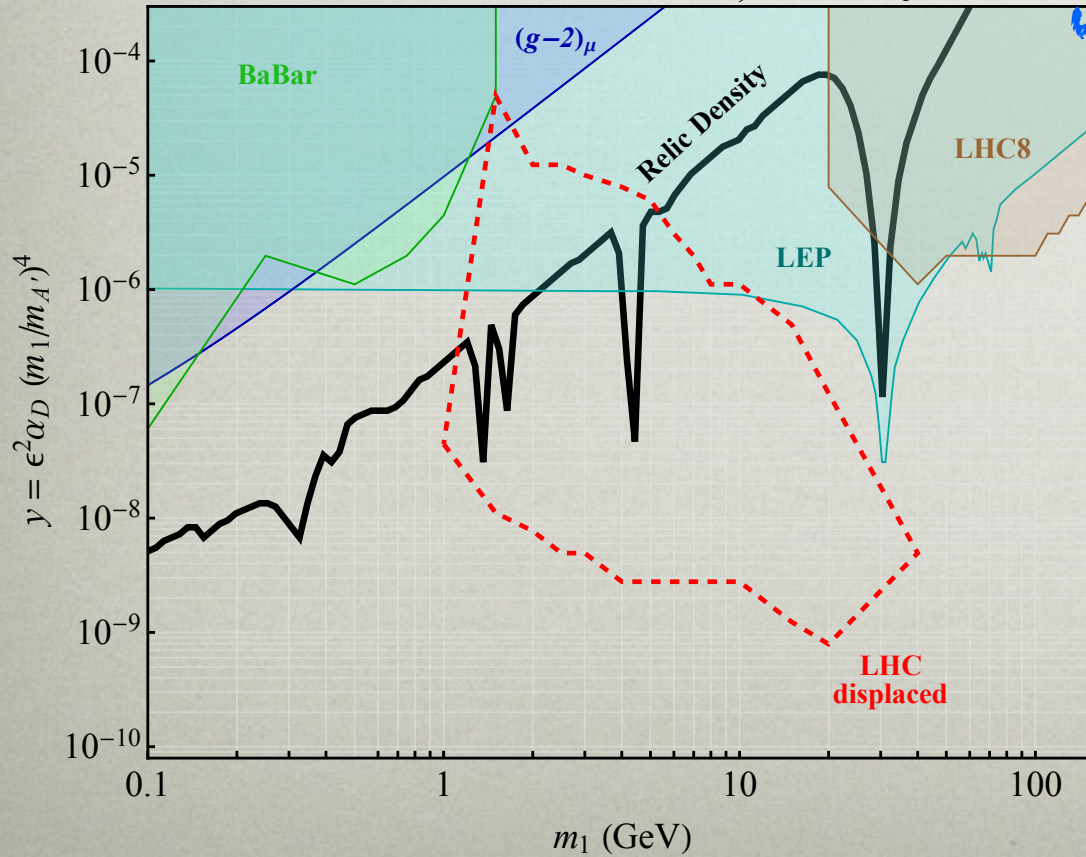
Fermion Thermal Relic iDM, $\Delta = 0.1 m_1$



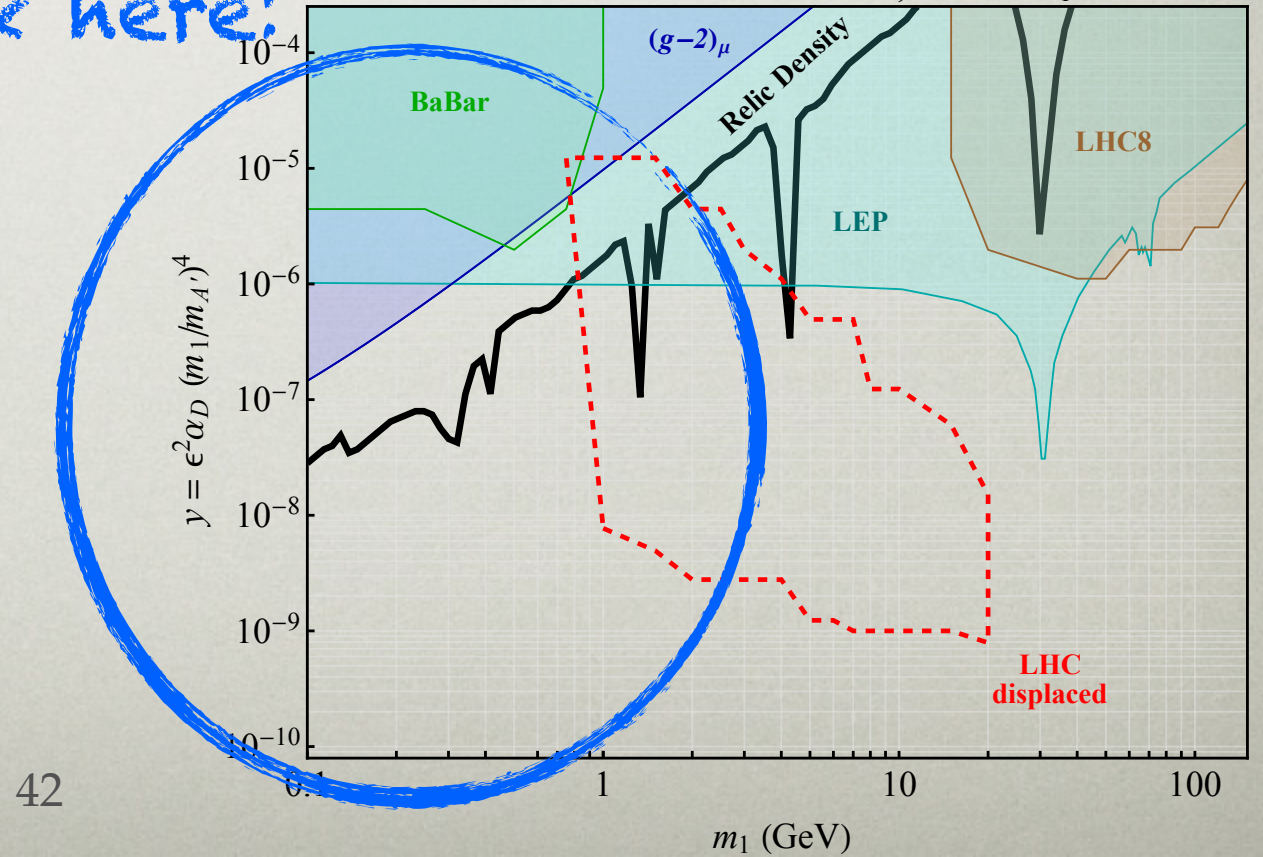
Fermion Thermal Relic iDM, $\Delta = 0.2 m_1$



Fermion Thermal Relic iDM, $\Delta = 0.3 m_1$



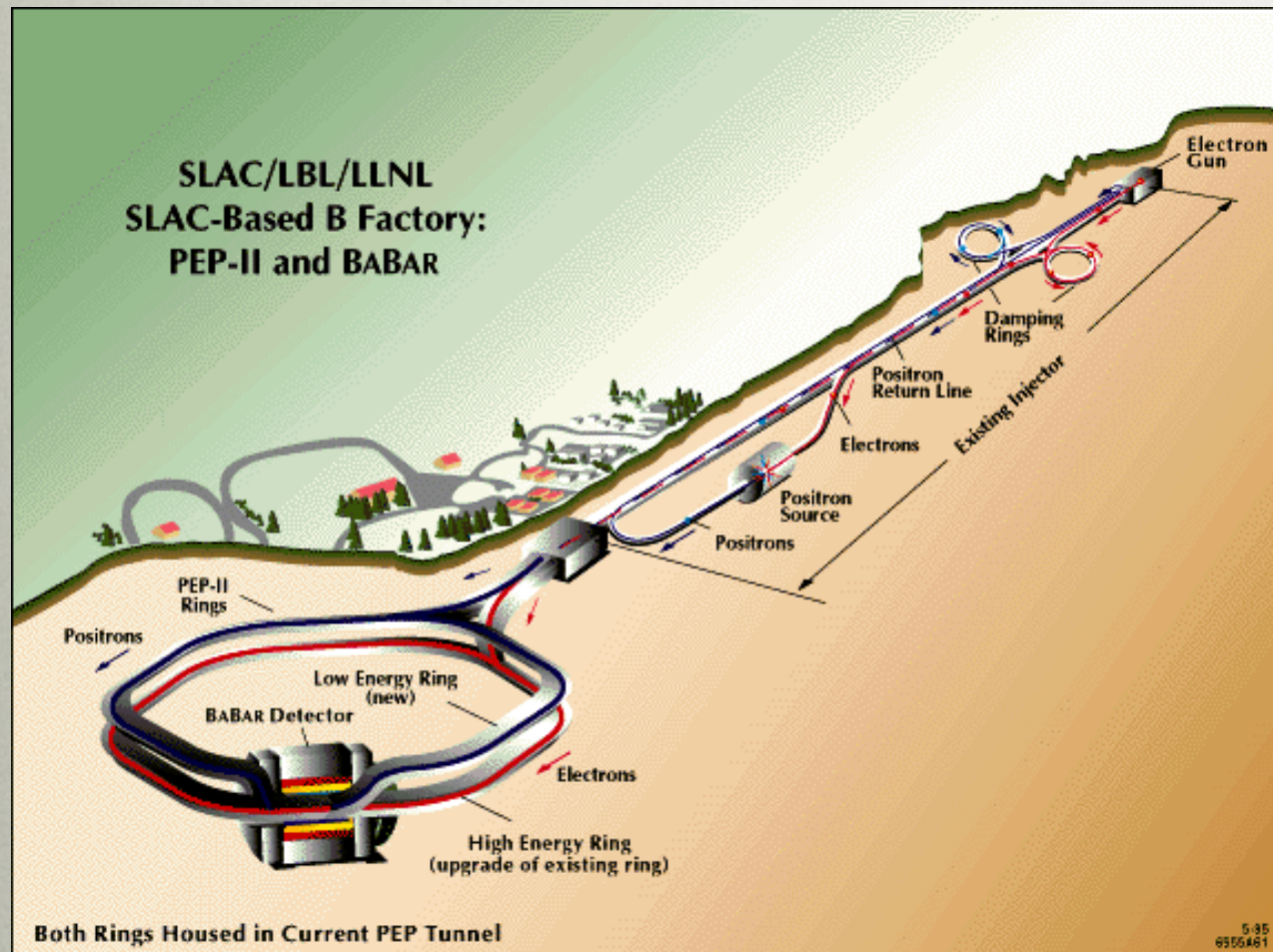
Fermion Thermal Relic iDM, $\Delta = 0.4 m_1$



LHC doesn't work here!

What About Low-Mass DM?

- Can use a **lower energy** collider to study these scenarios!



LHC: 13 TeV

BaBar / Belle II: 10.6 GeV

Other Collider Probes

1. BaBar:

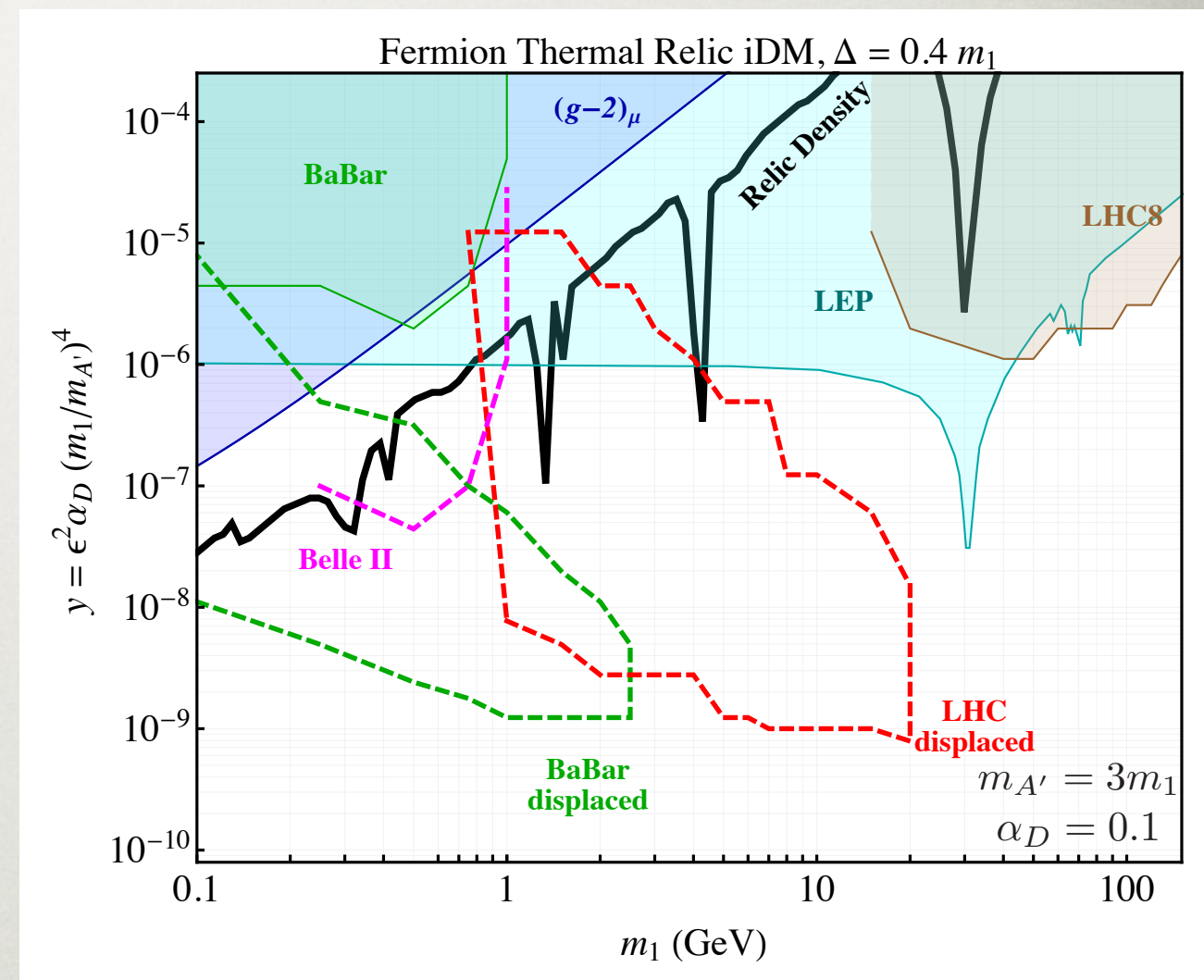
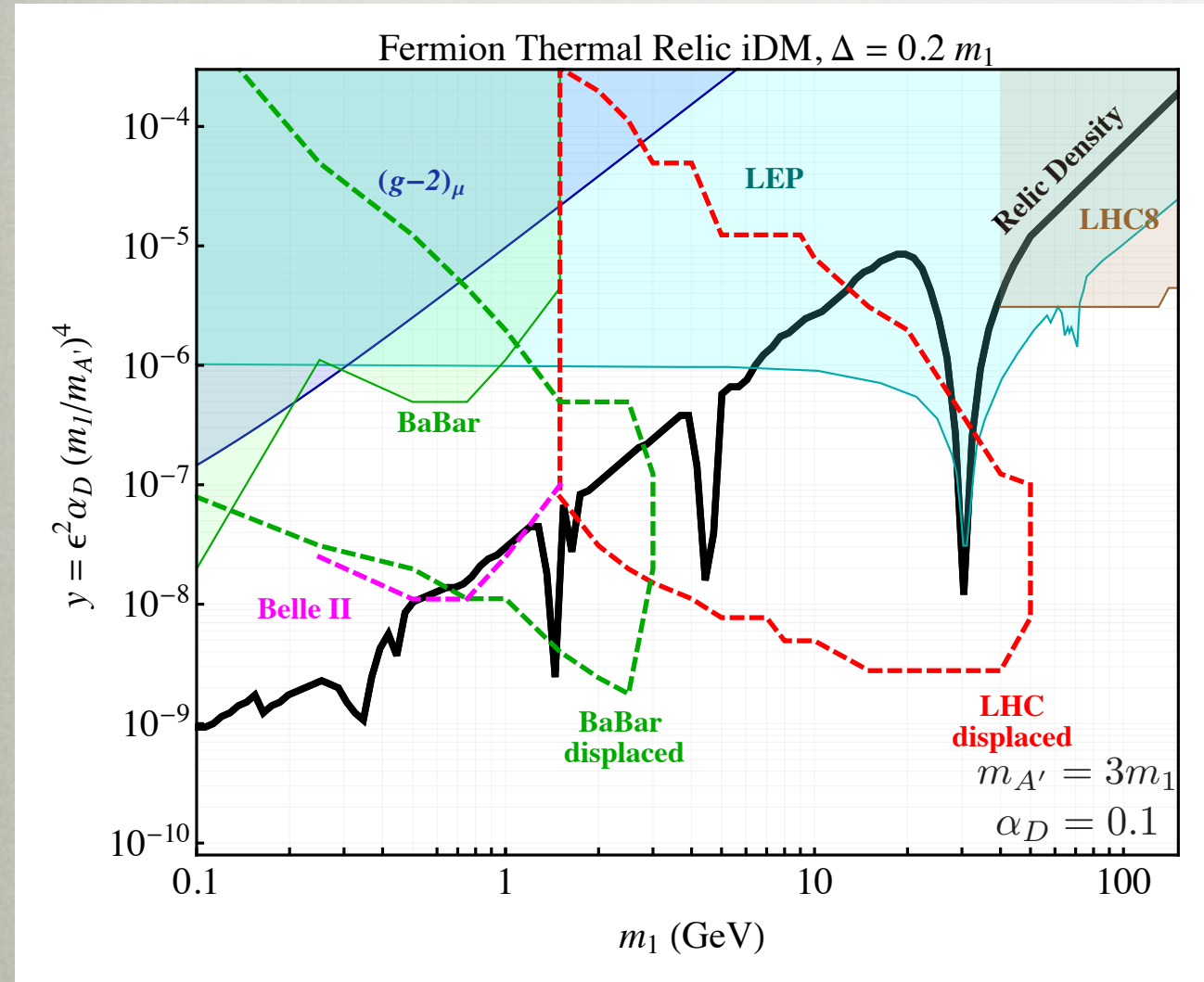
- Monophoton trigger for $\sim 60 / \text{fb}$
- Look for **monophoton** + displaced vertex (between 1-50 cm)
- Need DV discrimination of 100-1000 for background-free search

2. Belle II:

- Looks like it will be instrumented with a monophoton trigger
- Bounds from monophoton + missing energy search when tracks are below threshold / outside detector

(Essig *et al.*, 1309.5084)

LHC Results



Other Possibilities

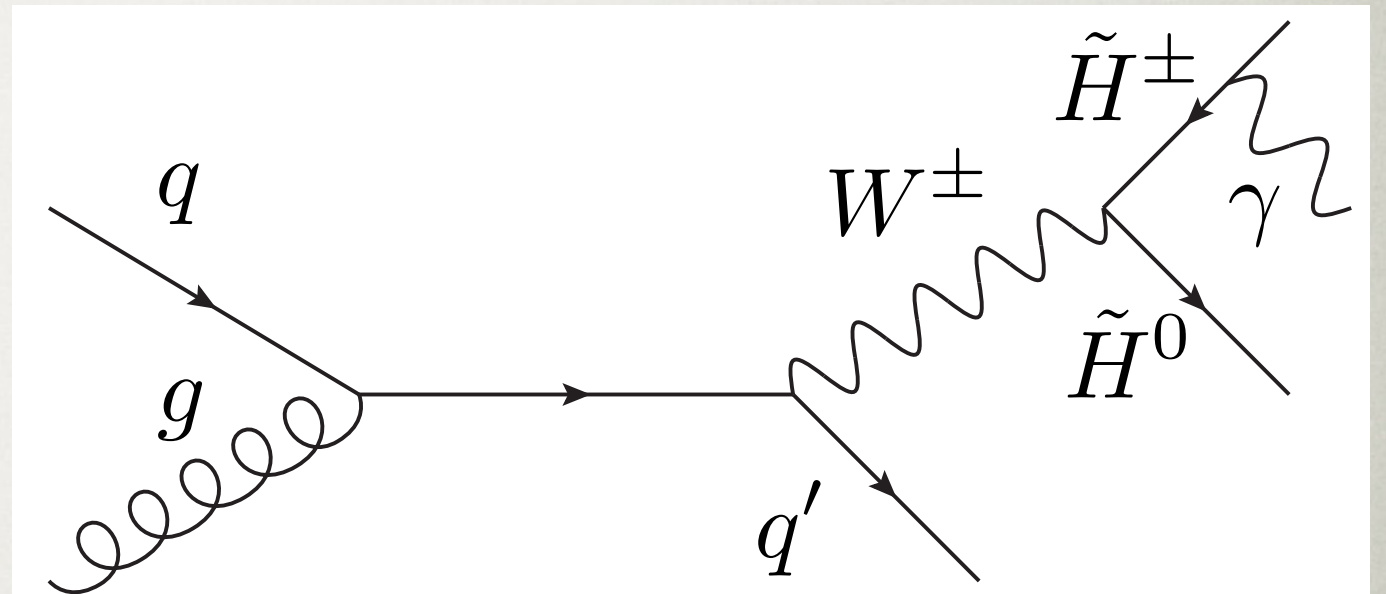
- Can use monojet + soft object tagging for other DM scenarios
- *E.g.*, doublet/“pure Higgsino”

$$\frac{\tilde{H}^\pm}{\underline{\underline{\tilde{H}^0}}} \quad c\tau(\tilde{H}^\pm \rightarrow \tilde{H}^0 \pi^\pm) \sim 5 \text{ mm}$$

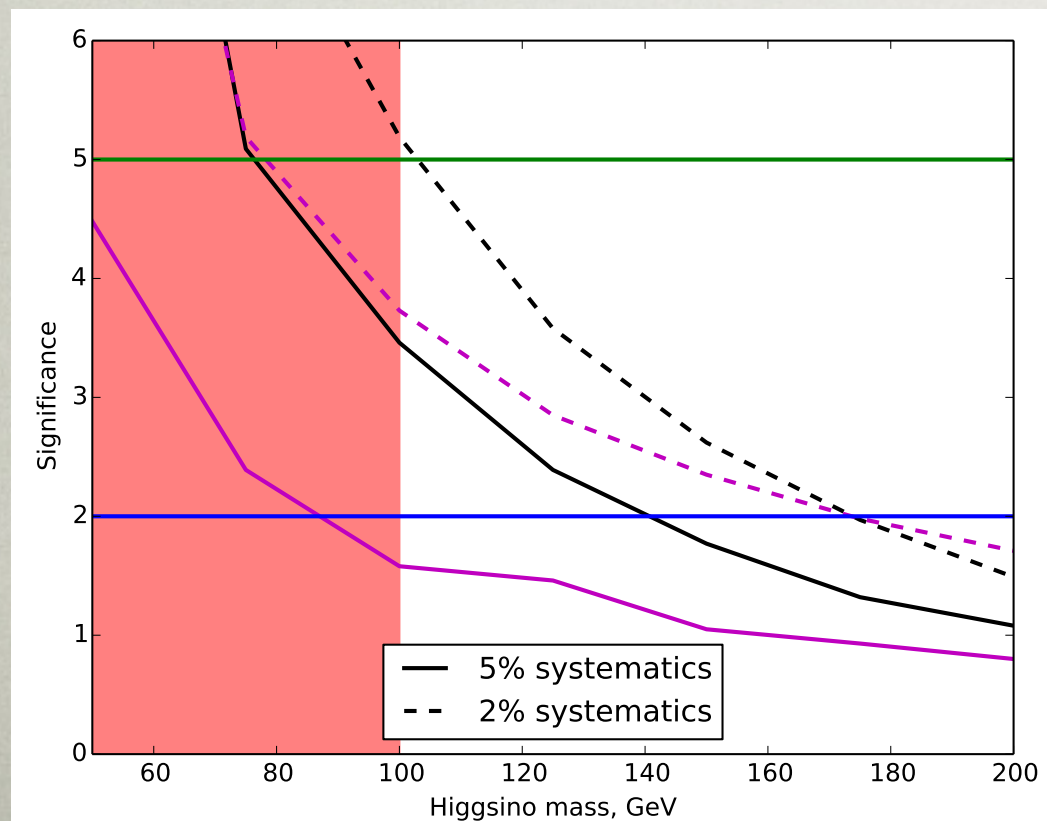
- Disappearing charged track, but **too short!**

Other Possibilities

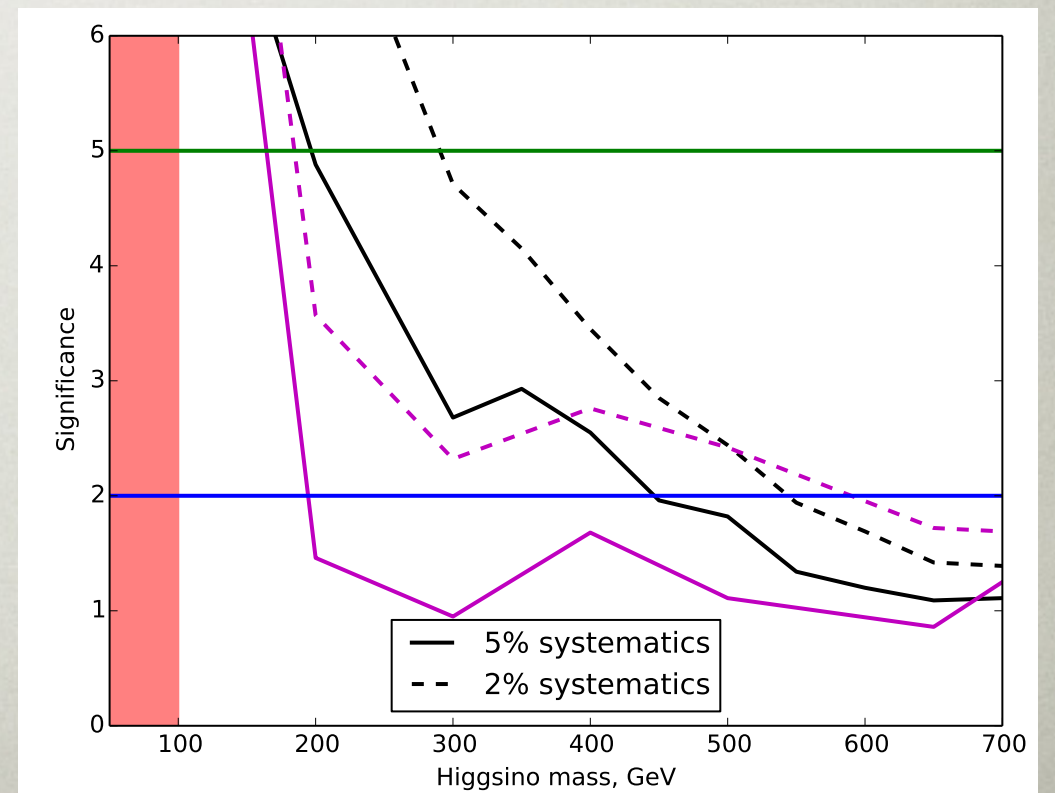
- Instead, use photon FSR
- Helps when systematics dominated



HL-LHC



100 TeV, 3/ab



Summary

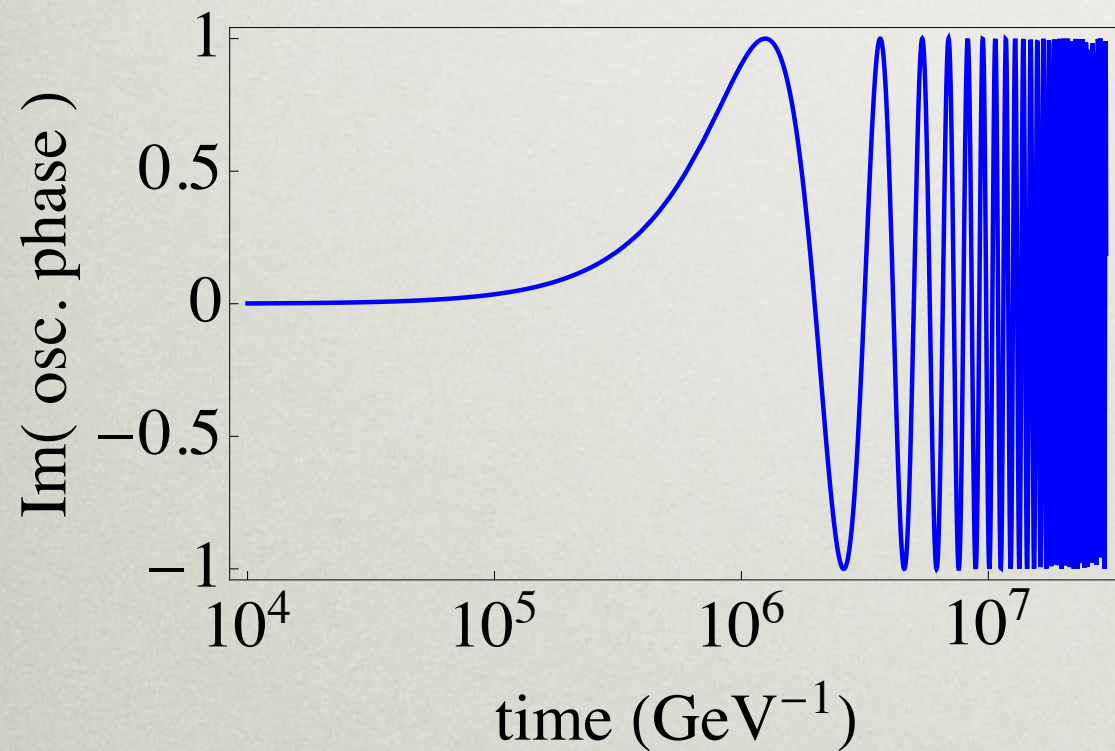
- New physics at or below the weak scale is motivated by naturalness, dark matter, baryogenesis, and neutrino masses
- Many diverse models and frameworks predict **similar signatures**
- In many cases backgrounds are so low that a discovery is possible even in very soft final states
- Let's hope for discovery in both high- and low-mass new physics!

Back-up slides

Asymmetry Generation

- The CP-violating rate comes from the interference of the diagrams

$$\Gamma(L_\alpha \rightarrow L_\beta) - \Gamma(\bar{L}_\alpha \rightarrow \bar{L}_\beta) \propto \text{Im} \left[\exp \left(-i \int_0^t dt' \frac{M_3^2 - M_2^2}{2T(t')} \right) \right] \text{Im} [F_{\alpha 3} F_{\beta 3}^* F_{\alpha 2}^* F_{\beta 2}]$$

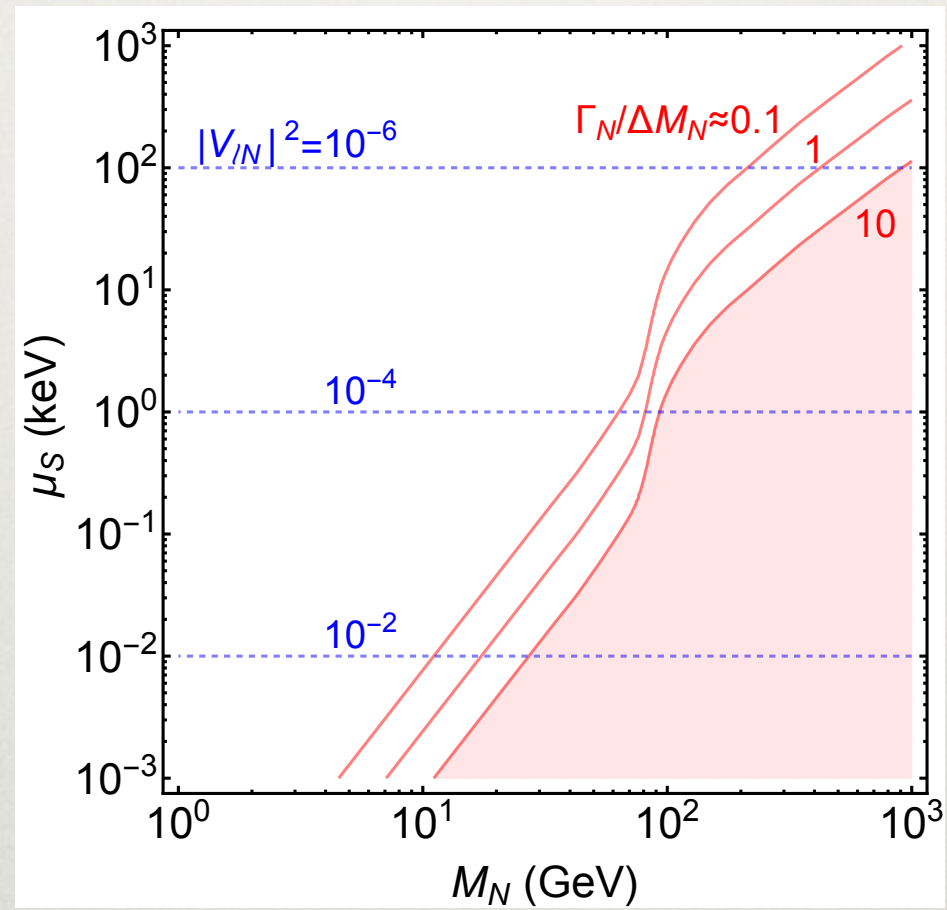


Asymmetry generation effectively **stops** when

$$(M_3^2 - M_2^2)/T \sim H$$

$$(T \gg M_N)$$

- Asymmetry is **larger** for oscillation at **later** times because integration time is longer



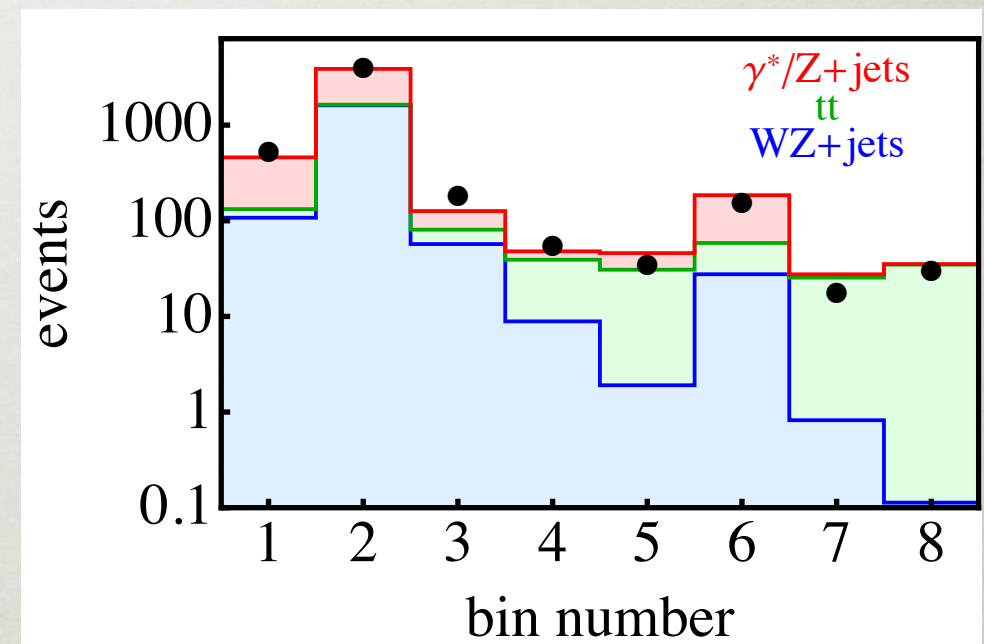
Resolved prompt decays

- Problem: these backgrounds are dominated by jets faking lepton
- A “fake simulator” for theorists has been proposed (Curtin, Galloway, Wacker 2013)

$$\epsilon_{j \rightarrow \ell} = \text{p}_T\text{-dependent probability of jet faking lepton} \times \text{map from jet kinematics to lepton kinematics}$$

CMS trilepton search (low HT, low MET)

MadGraph 5 + Pythia 6 (matched)



OSSF1, < Z
OSSF1, = Z
OSSF1, > Z
OSSF0
OSSF1, < Z
OSSF1, = Z
OSSF1, > Z
OSSF0

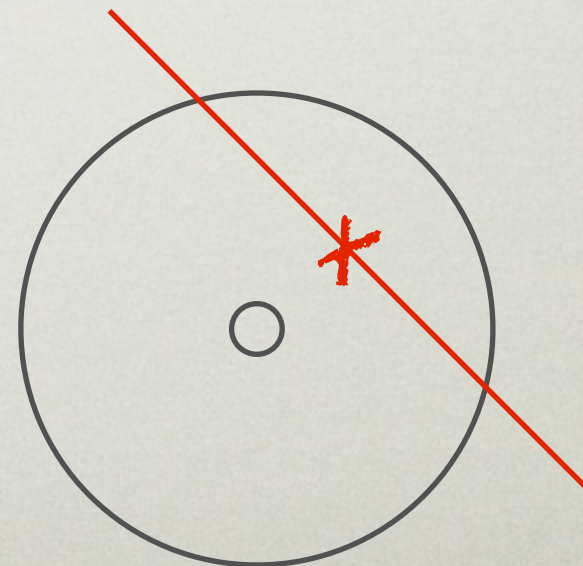
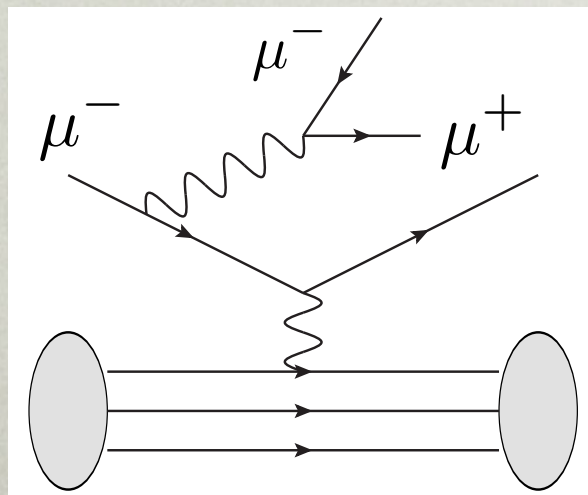
0b

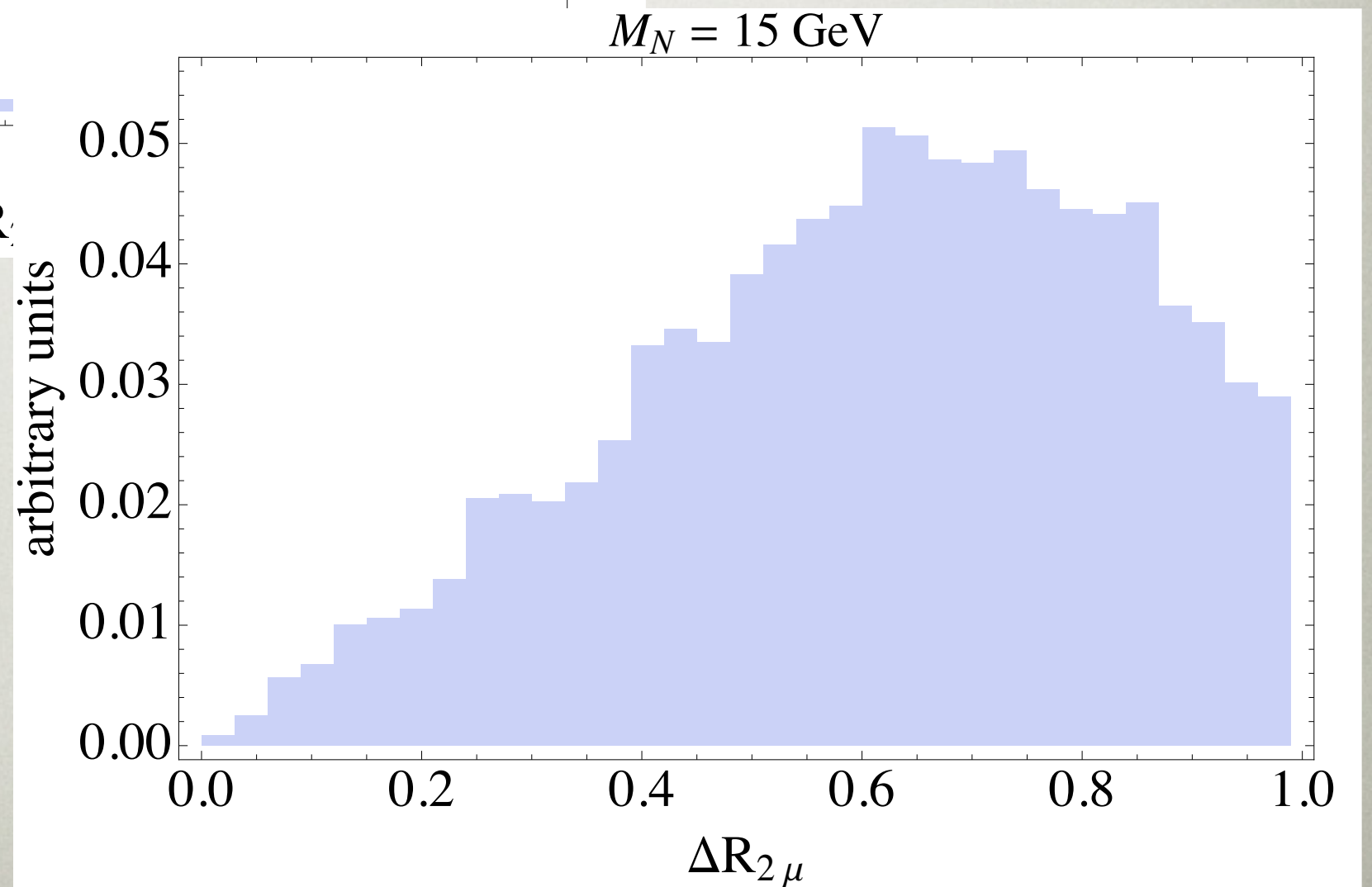
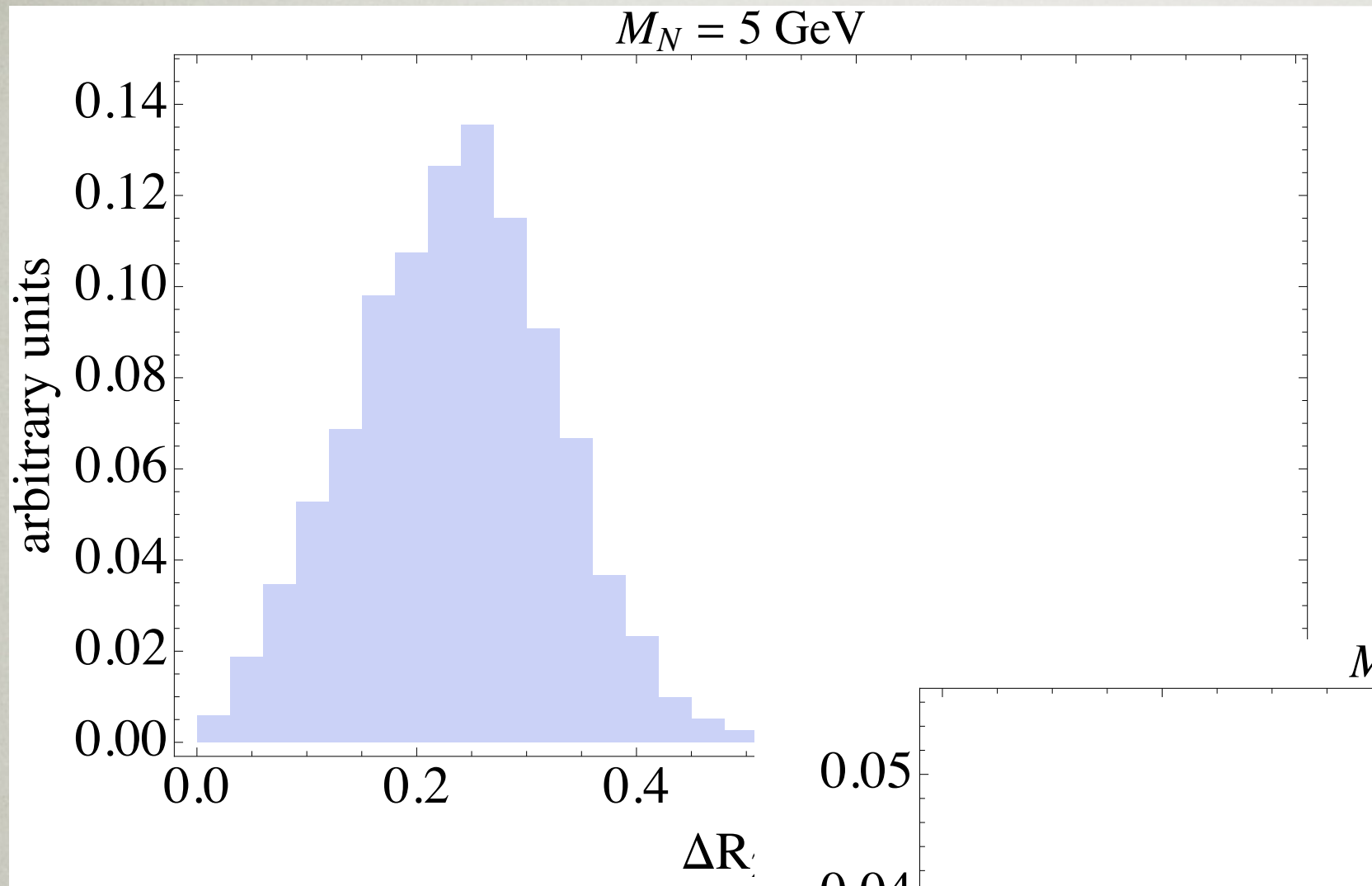
1b

Also checked method with ATLAS same-sign muon & CMS same-sign muon + jets analyses

Lepton jets from RH neutrinos

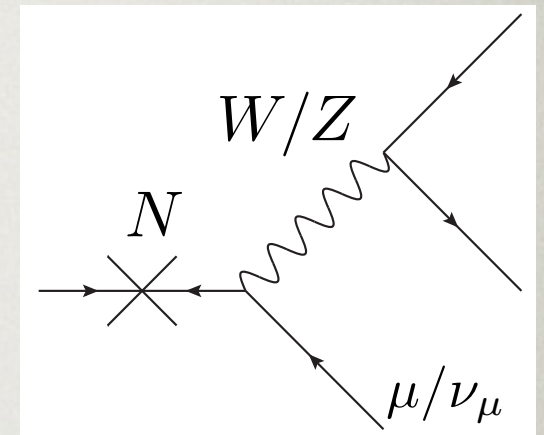
- ATLAS has a lepton jet search, but it is for a **pair** of muon jets
 - Veto reconstructed tracks in inner detector (no displaced vertex reconstruction), but require track impact parameter to be in inner detector
 - They see ~ 10 background events
- There are some unknowns in extrapolating results to single lepton jet
 - Extrapolation based on single cosmic muon flux gives vanishing result!
 - Can get “muon bundles” from muon trident production
 - Single cosmic muon can fake back-to-back muons from inner detector



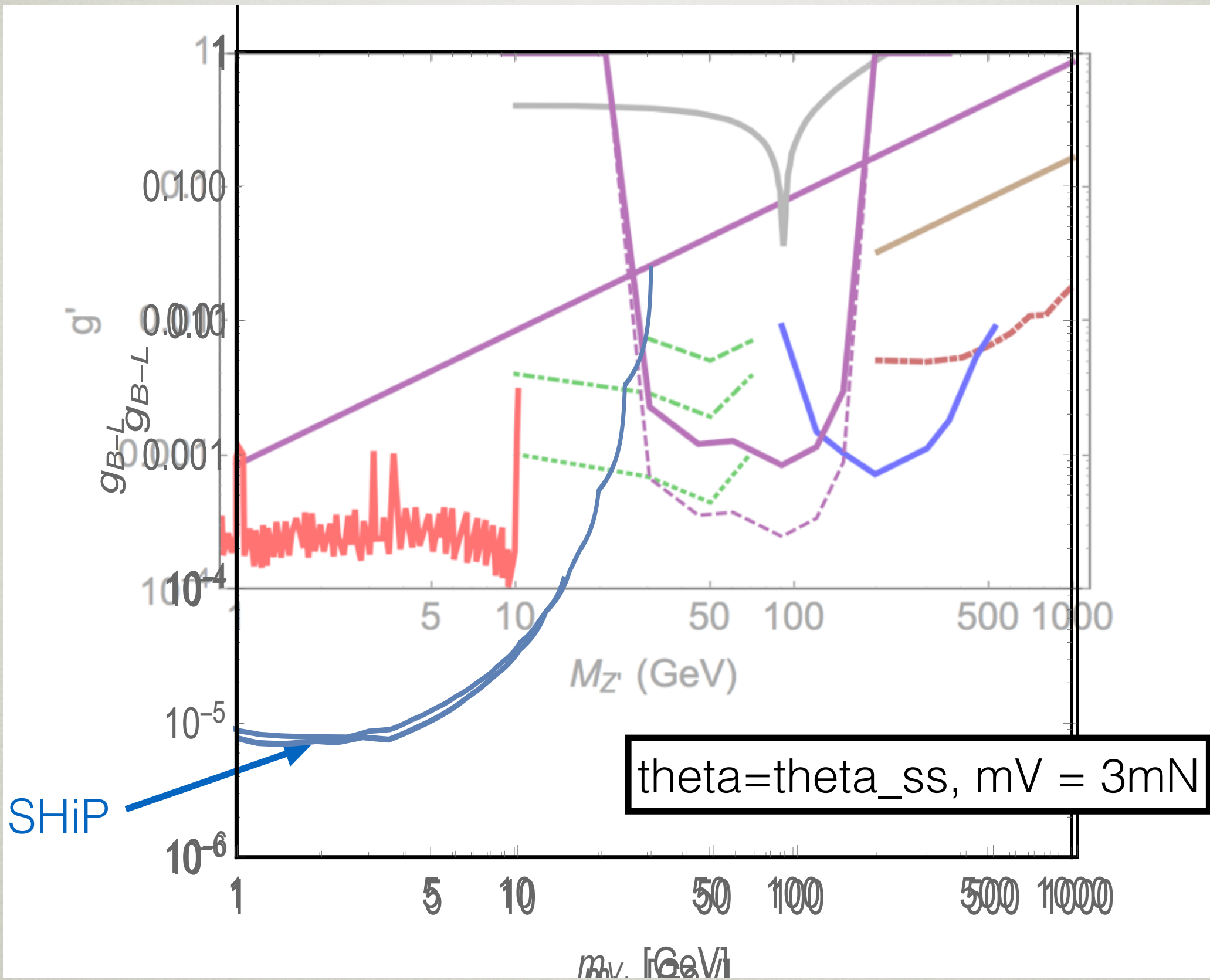


B/L gauge forces at colliders

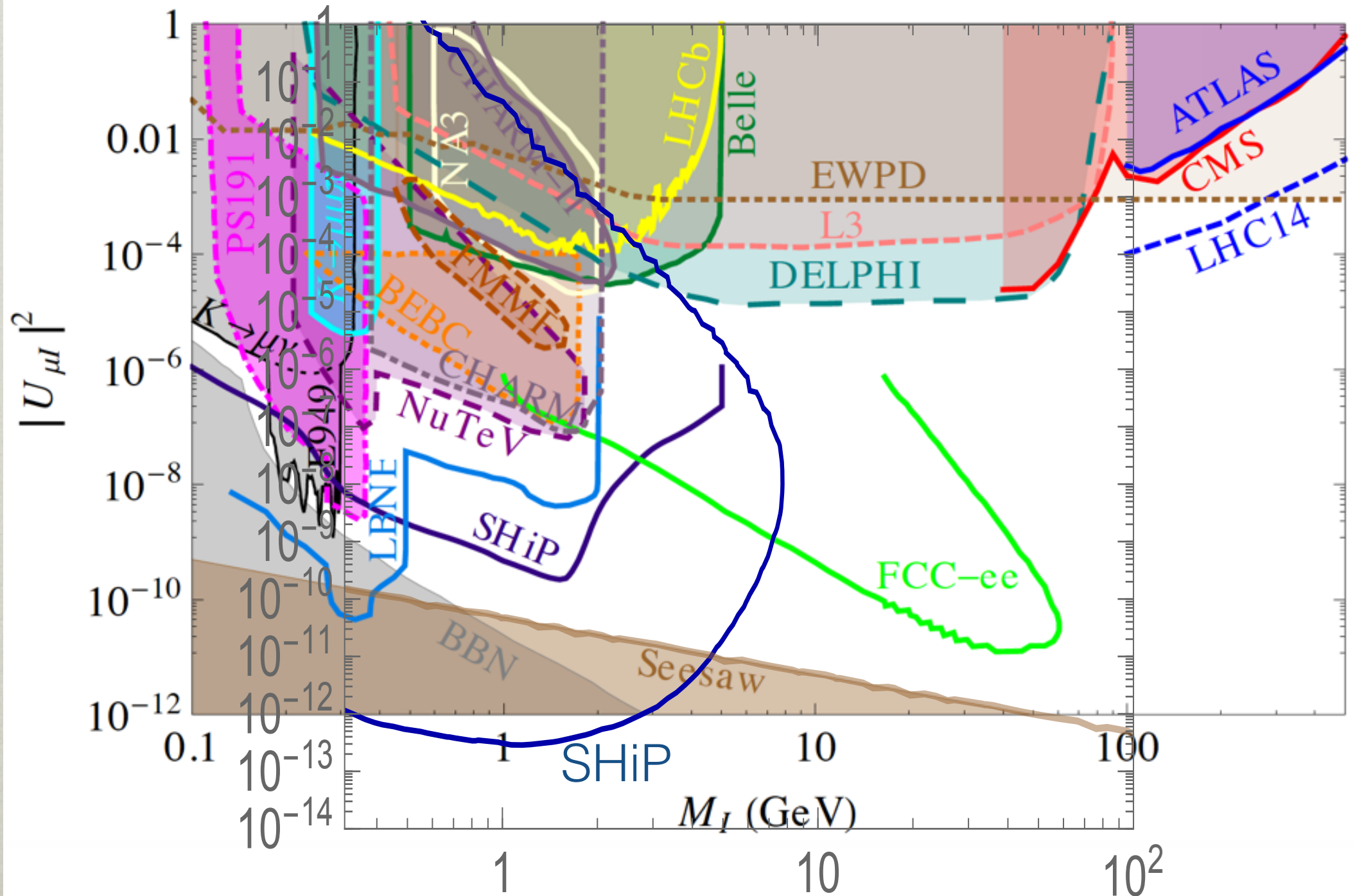
- N decays via (off-shell) W/Z
 - Get displaced muon in $> 50\%$ of decays
 - Use lepton triggers



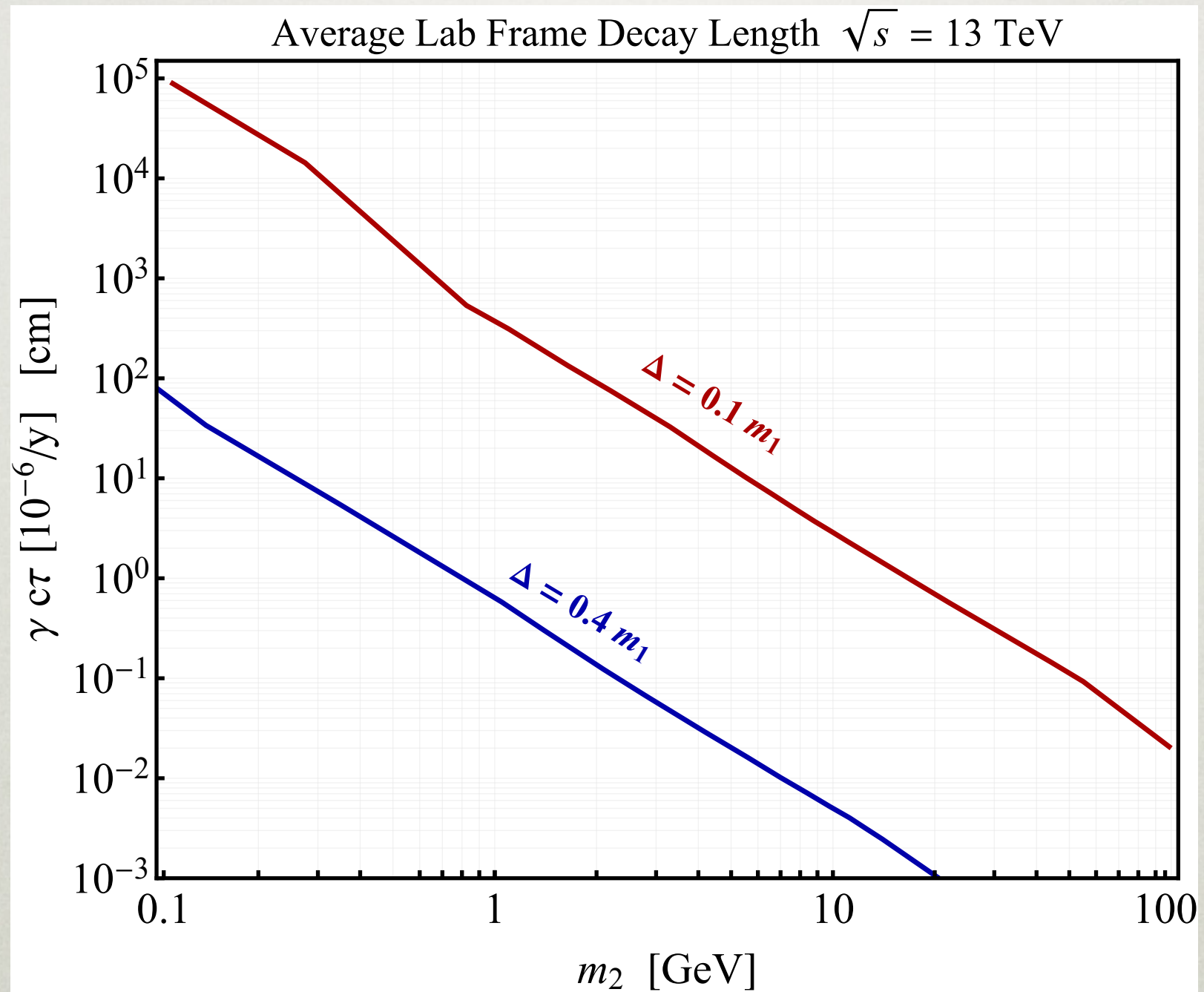
- Current searches have high thresholds, unnecessary restrictions
 - 1 electron and 1 muon, not required to reconstruct vertex, > 1 mm displacement, 0.05 bkd events (CMS, 1409.4789)
 - Pair of electrons or muons at vertex, $m_{\ell\ell} > 15$ GeV, > 0.2 mm displacement, 0 bkd event prediction with 0 bkd in control region (CMS, 1411.6977)
 - Pair of leptons at vertex, leading lepton > 50 -100 GeV, $m_{\ell\ell} > 6$ GeV, ≈ 1 mm displacement, 10^{-3} bkd events (ATLAS, 1504.05162)
 - Muon + at least 4 tracks at vertex, muon > 50 GeV, $m_{\ell\ell} > 6$ GeV, ≈ 1 mm displacement, 10^{-3} bkd events (ATLAS, 1504.05162)
- We require one of the above + some other displaced object \rightarrow expect bkd-free



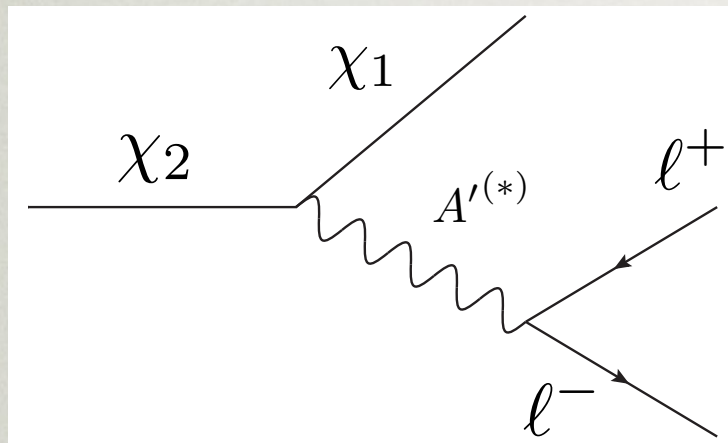
$g' = 0.001, mV = 3mN$



iDM Lifetimes

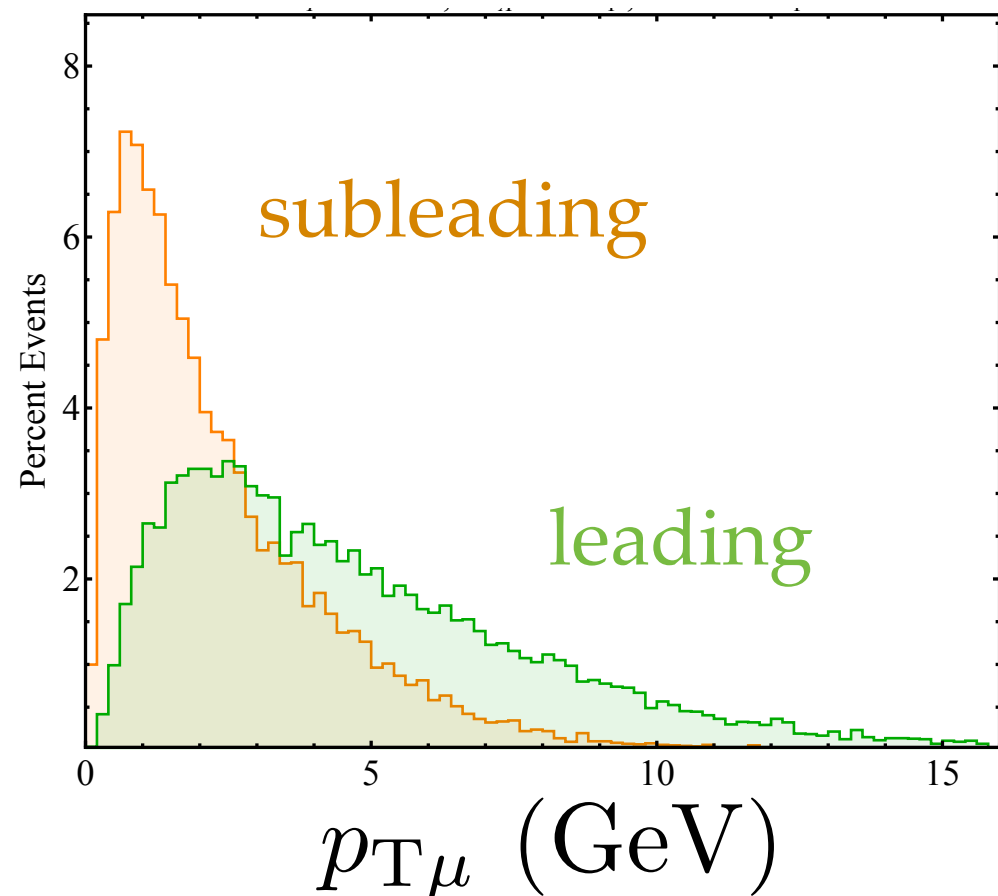


Improving the Searches

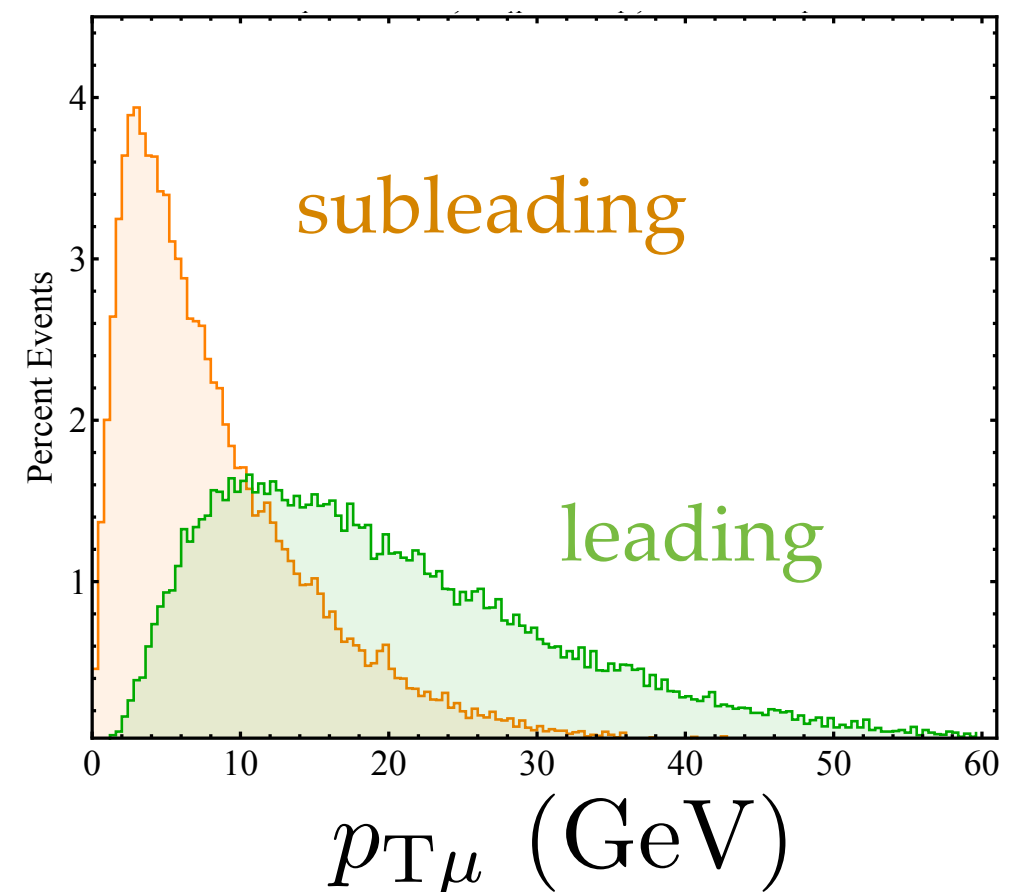


- For small splittings, leptons are soft, so trigger on monojet + MET

$m_1 = 5 \text{ GeV}$ $\Delta = 0.1 m_1$



$m_1 = 50 \text{ GeV}$ $\Delta = 0.4 m_1$



Backgrounds

- Random track crossings
 - Can't do first principles estimate
 - We look at QCD events ($p_{Tj} > 120$ GeV, no MET cut) and find the efficiency for two isolated muon tracks satisfying the signal requirements
 - We find no events, bounds QCD contribution < 100 fb
 - Adding requirement for additional invisible Z/W, kinematic requirements leads to expectation of \approx few events

2. Photon conversion to muons

- Cross section for Z + jet + gamma is ~ 100 fb after jet p_T , photon E_T cut
- Even though the probability for conversion to leptons is $O(1)$, the ratio of e/mu is

$$\frac{\sigma(\gamma \rightarrow \mu\mu)}{\sigma(\gamma \rightarrow ee)} \sim \frac{m_e^2}{m_\mu^2}$$

3. Pile-up crossings

- Since LJ is collinear with χ_2 , require that muons point back to same vertex as jet

