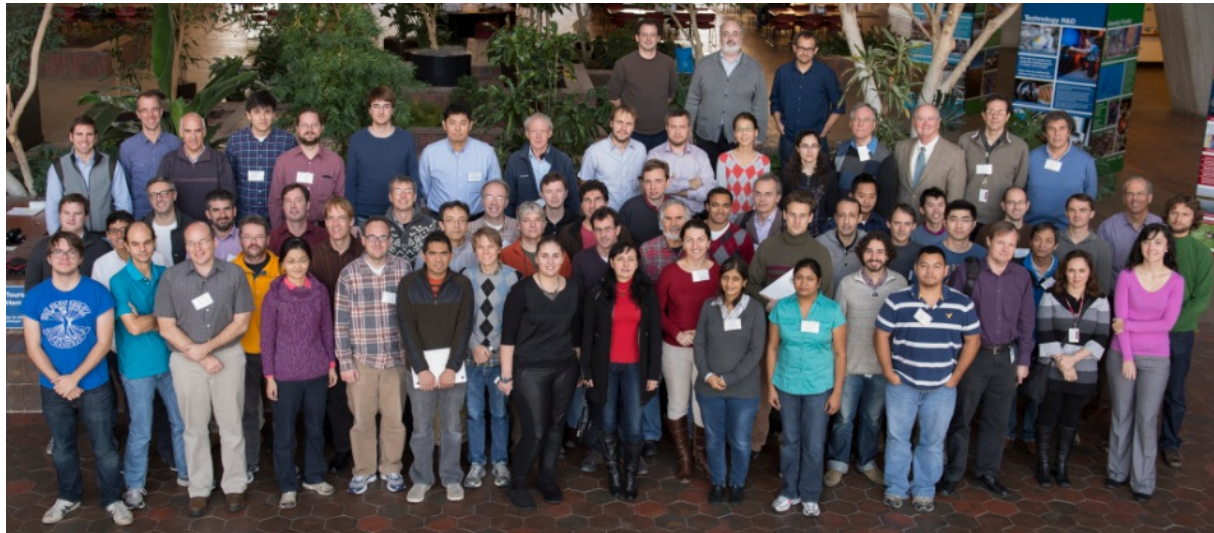


SUSY scenarios we're
not looking for* but
should be

Markus Luty
UC Davis

SUSY at the Near Energy Frontier



11-13 November 2013
Fermi National Accelerator Laboratory

Disclaimer

Many important “loopholes” are covered in existing or planned searches

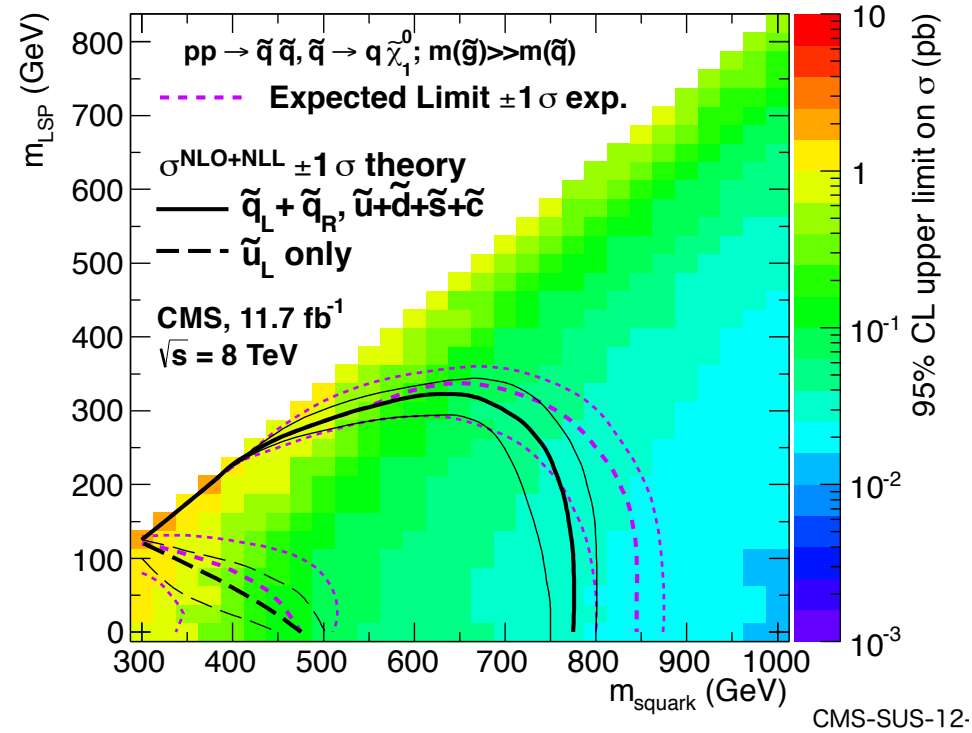
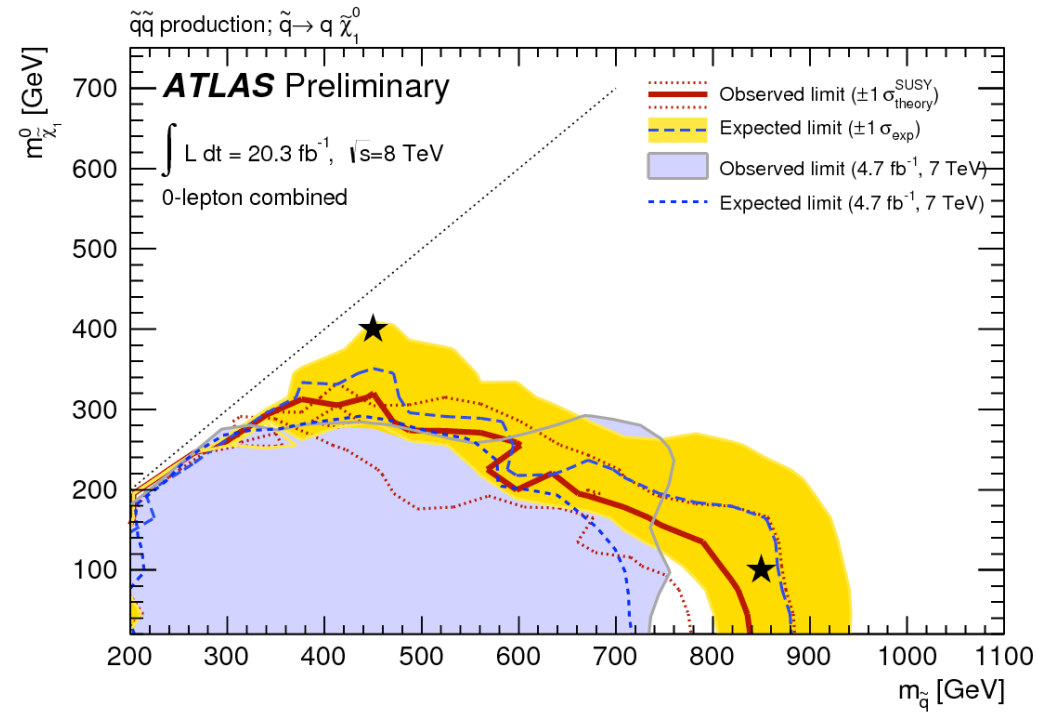
- RPV
- Reduced MET
compressed spectra, “stealth” SUSY,...
- Long decay chains
- NLSP \rightarrow gravitino + X
- \vdots

This talk will discuss a few areas that may benefit from more attention

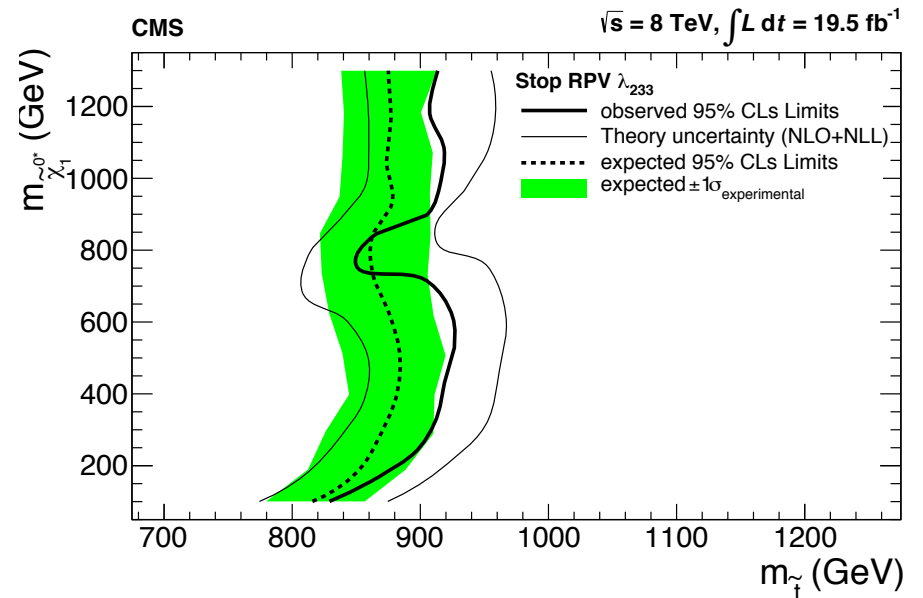
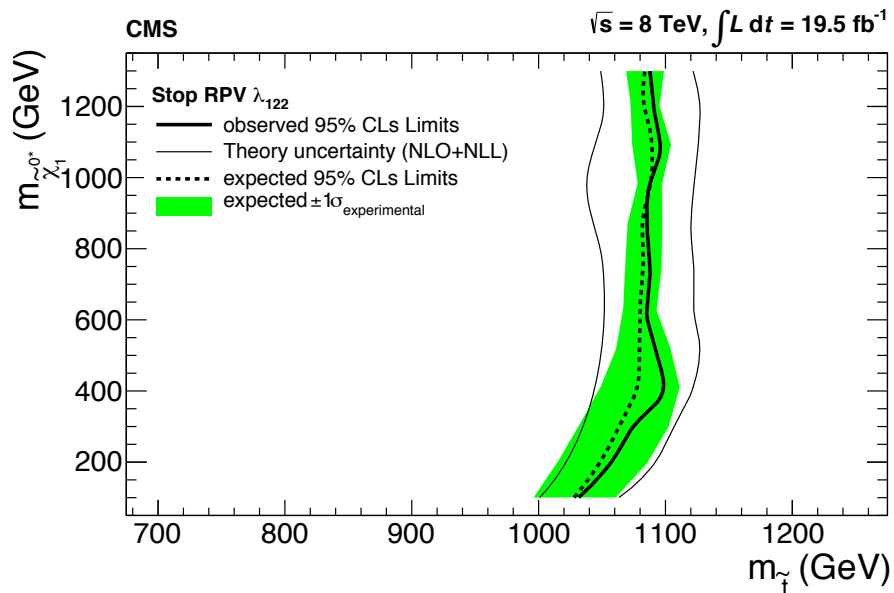
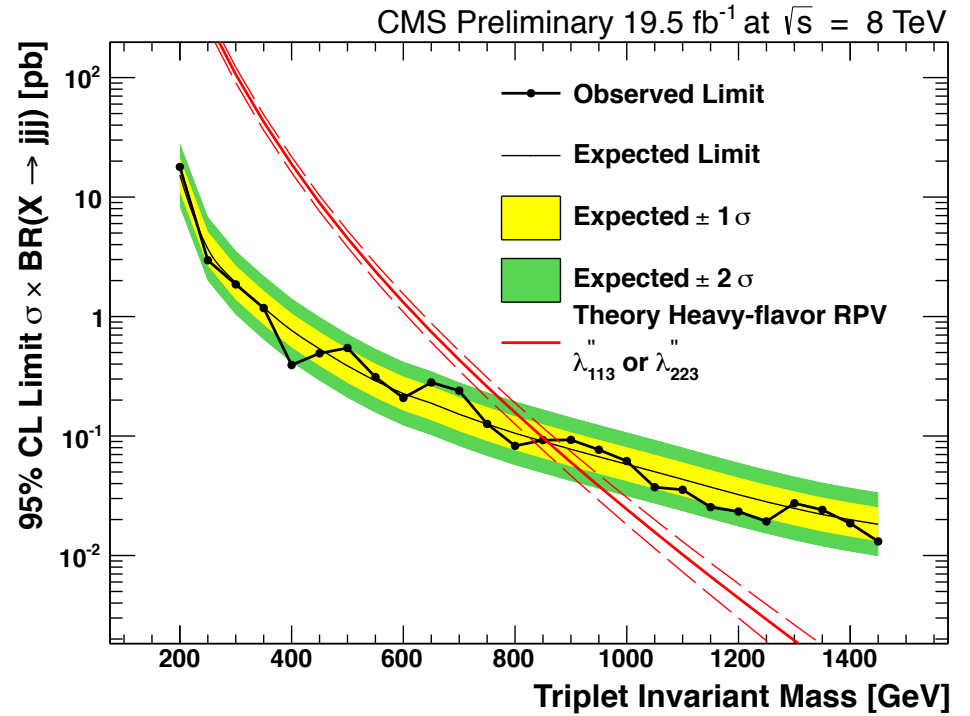
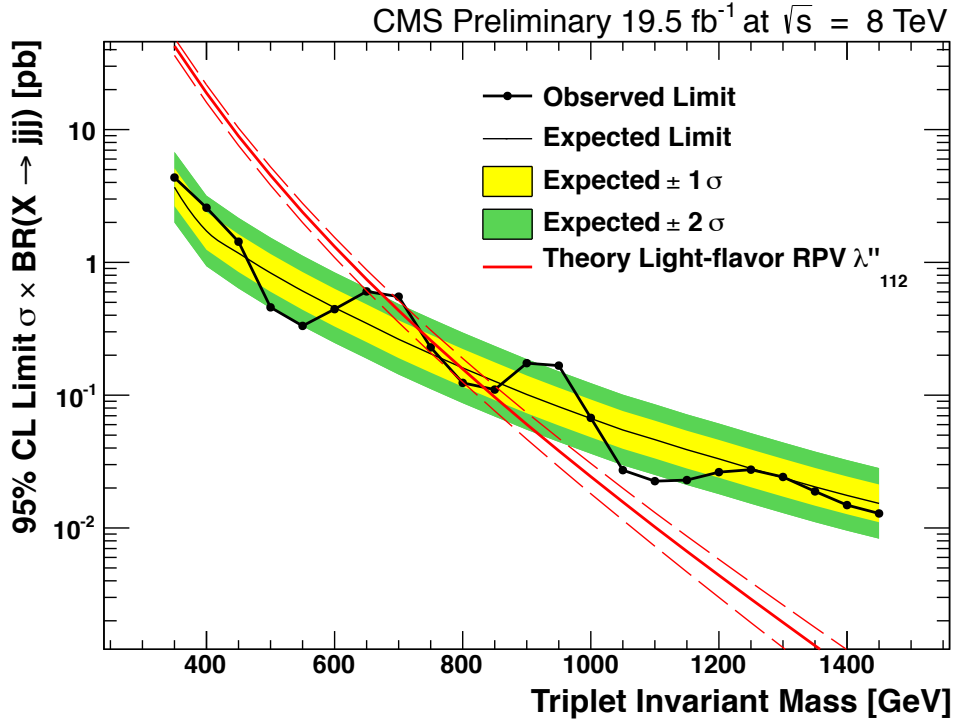
A personal selection, no attempt at completeness

Existing SUSY Searches

Compressed Spectra

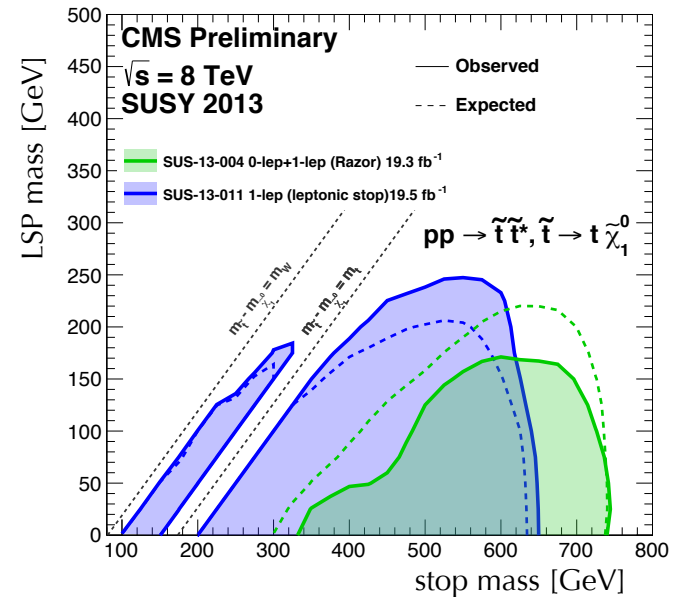
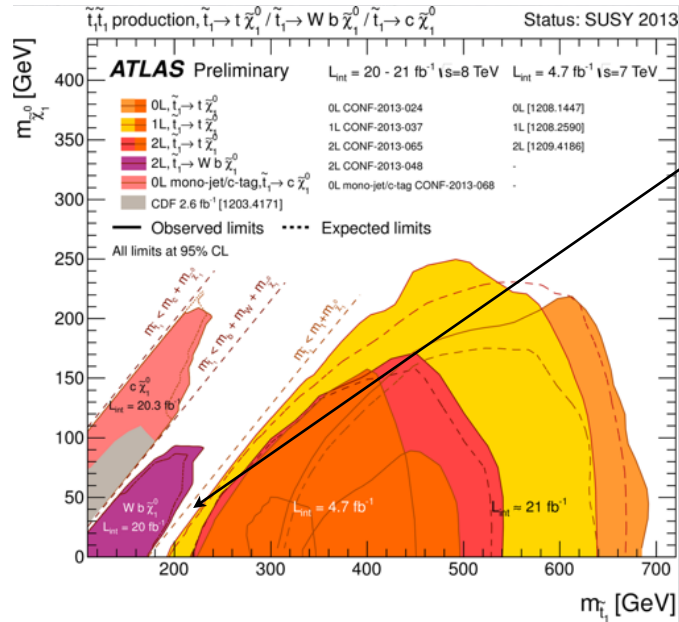


R-Parity Violation

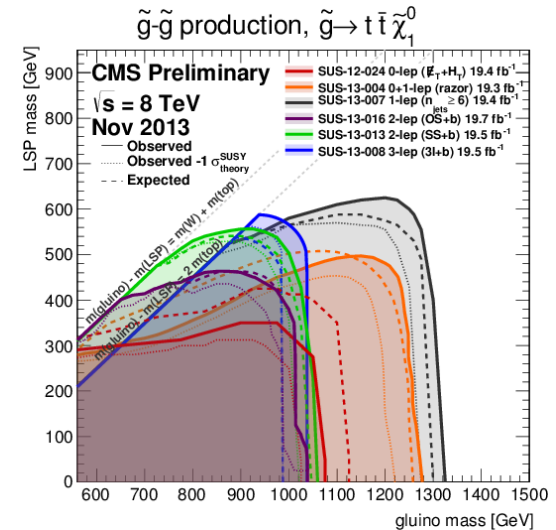
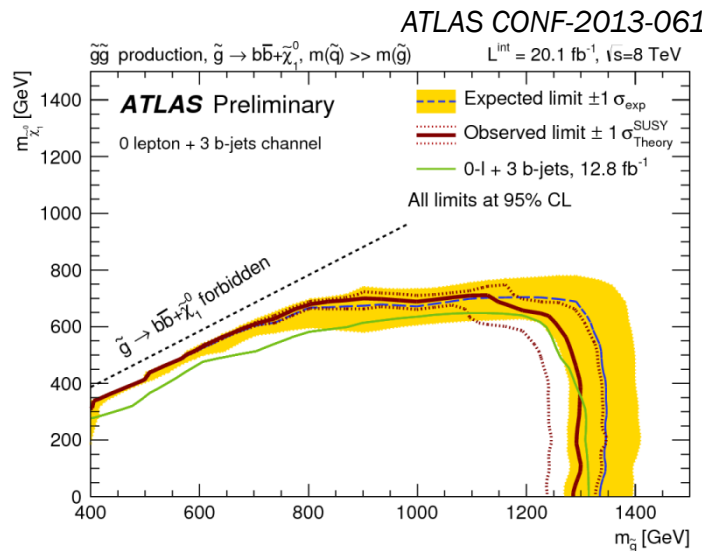


Natural SUSY

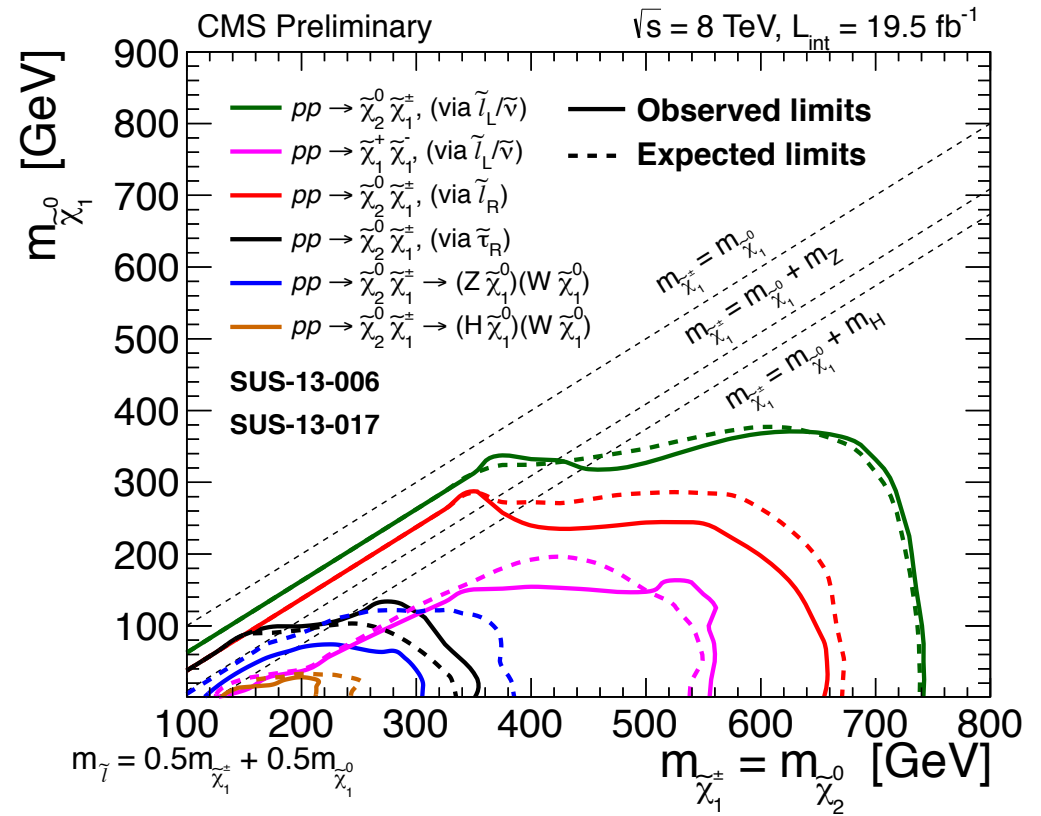
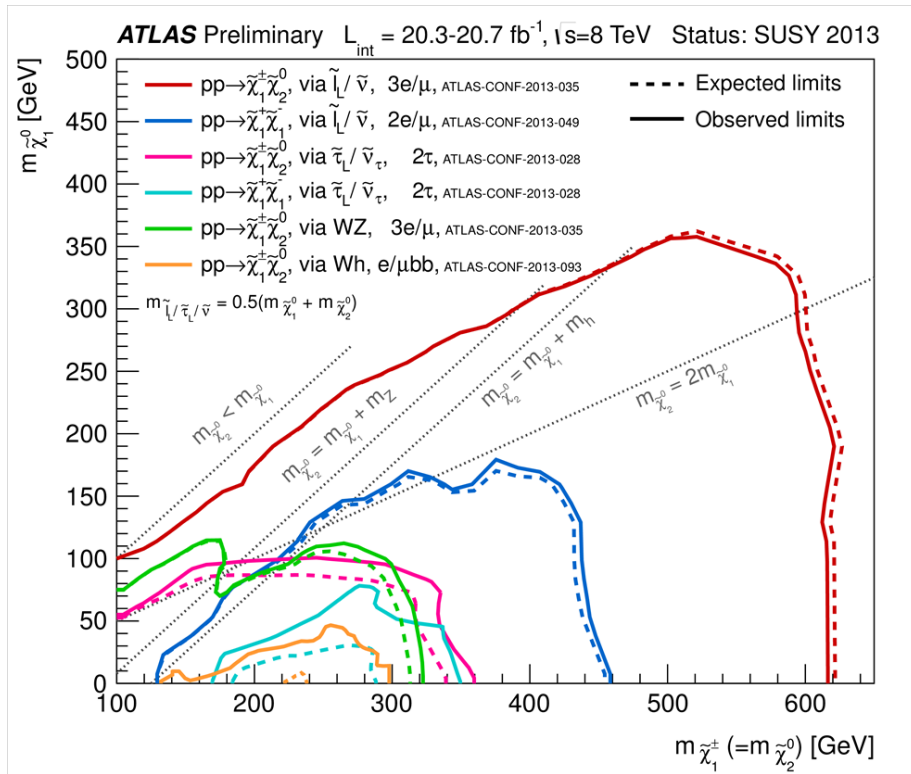
Stops



Gluininos:



Electroweakinos



The Big Picture



125 GeV Higgs
SM-like couplings

SUSY

Composite
Higgs

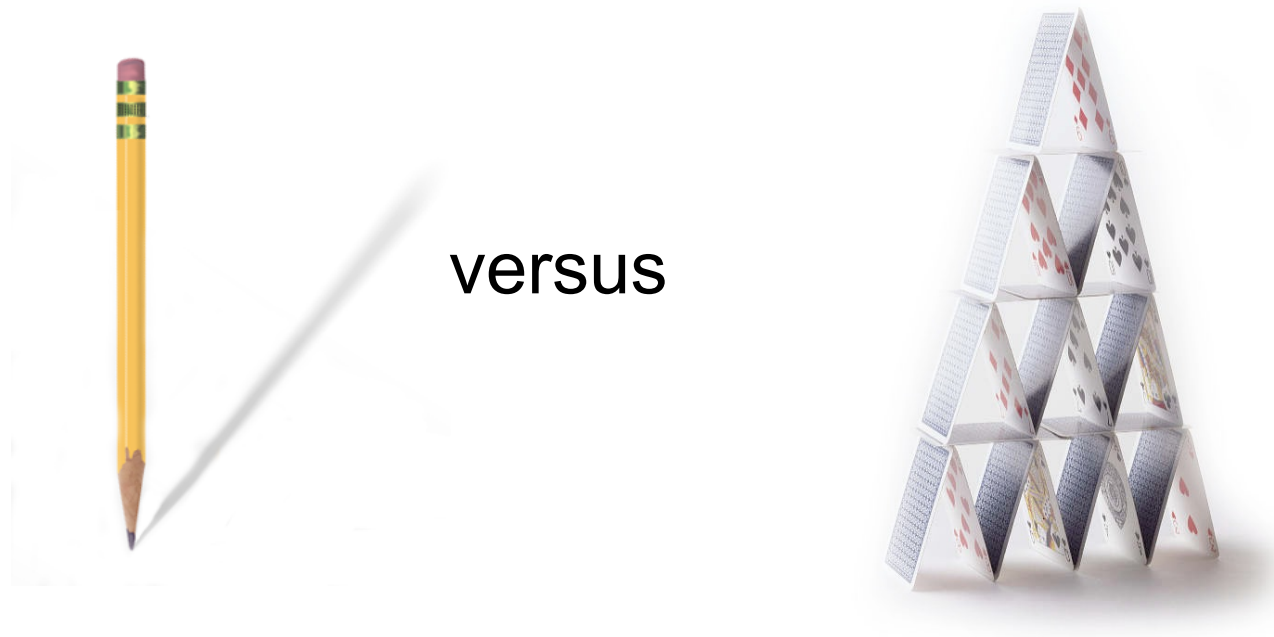
We don't know
what we're doing

Simplicity vs. Naturalness

SM is the perfect effective theory...

...if one parameter is tuned

$m_{H,\text{eff}}^2 =$ function of more fundamental parameters



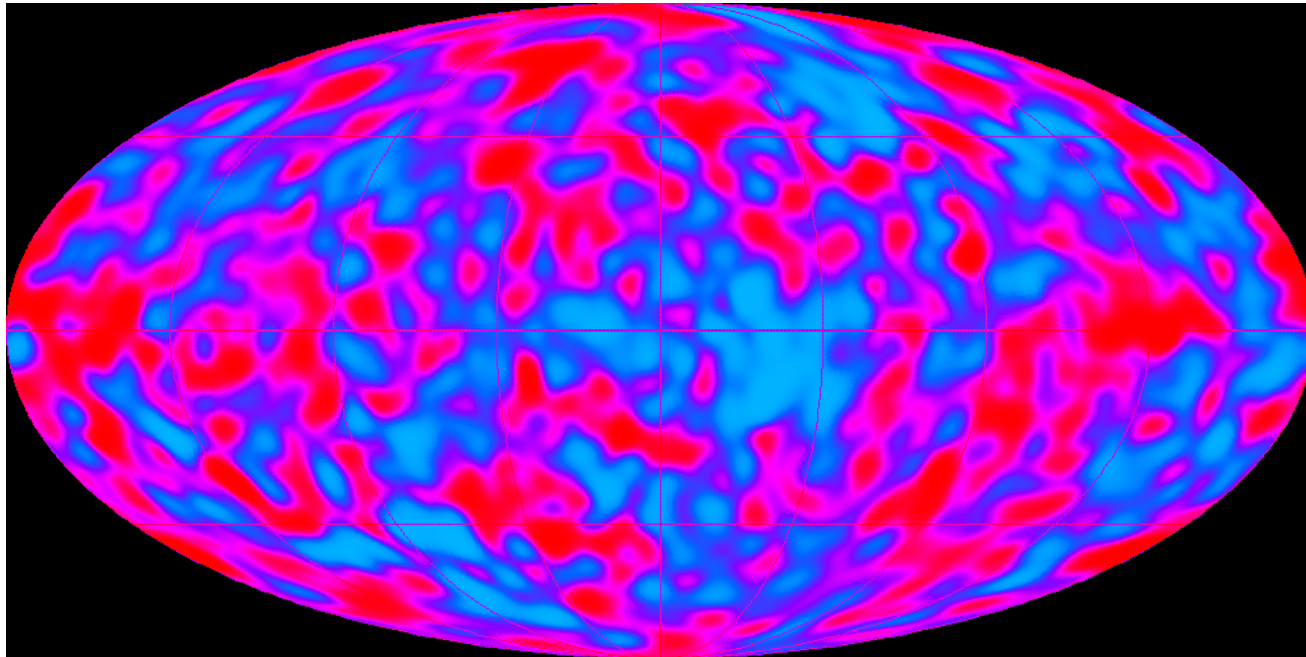
Naturalness = no tuning of parameters

Simplicity = minimality of structure & assumptions

An Unnatural Story

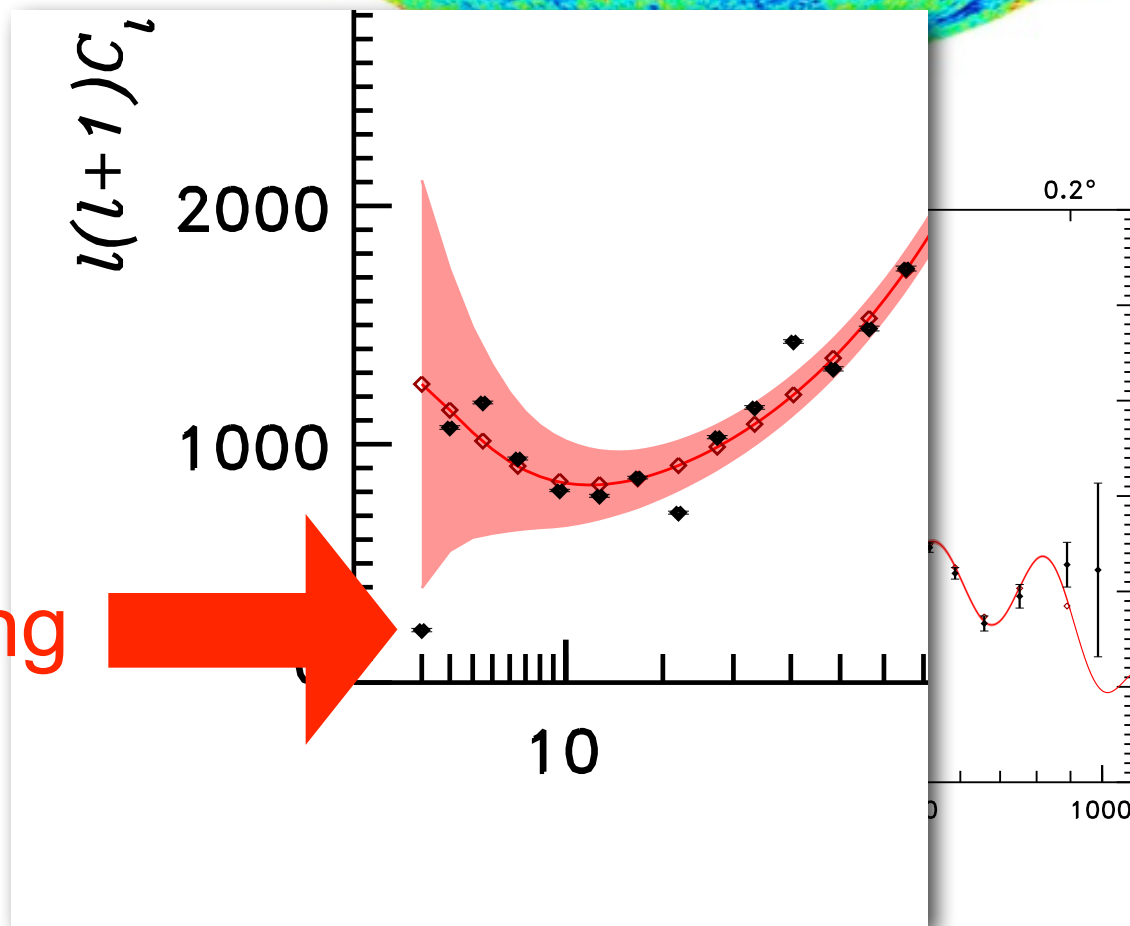
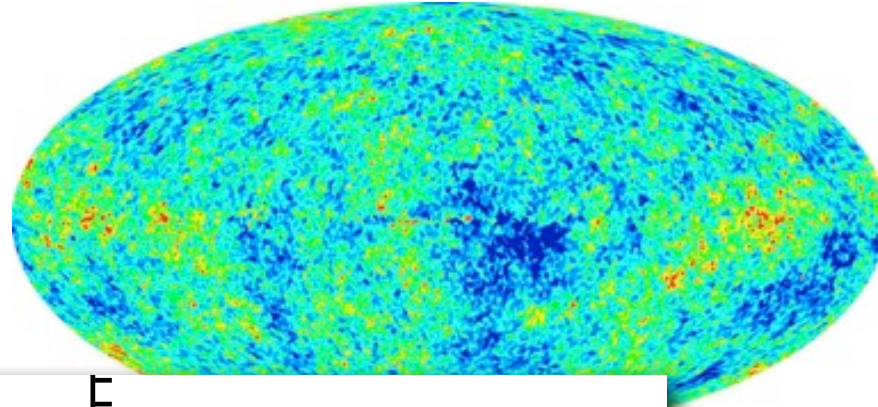
In the early 1990s limits on the CMB quadrupole were pushing the limits of cold dark matter cosmology...

...and then came COBE



So what about the quadrupole?

Tuning!



Outline

- Tuned SUSY
- Natural SUSY
 - BMSSM Higgs
 - Displaced vertices
 - R-parity violation
 - Hidden sector dark matter

Tuned SUSY

Occam's razor:

“Entities must not be multiplied beyond necessity”



MSSM with one tuning is arguably the simplest explanation of particle physics data

Implications:

- SUSY most likely “just around the corner”
- Keep looking for standard SUSY signals

Natural SUSY

Crucial to test...

Requirements:

$$\mu \lesssim 200 \text{ GeV}$$

$$m_{\tilde{t}} \lesssim 500 \text{ GeV} \quad \longleftrightarrow$$

$$m_{\tilde{g}} \lesssim 1.5 \text{ TeV}$$

EWino searches,
stop, sbottom searches

But also: Higgs sector beyond MSSM

BMSSM Higgs

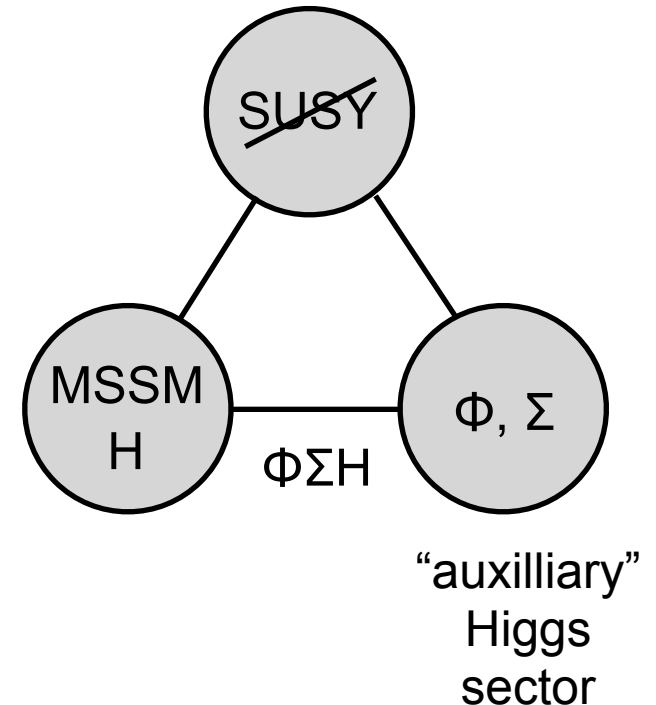
$$V_{\text{eff}} = m_H^2 |H|^2 + \lambda |H|^4 \quad \Rightarrow \quad (126 \text{ GeV})^2 = 2\lambda v^2$$

Need additional contributions to quartic

- NMSSM
- Non-decoupling D-terms

...or additional sources of EWSB

- “Induced” EWSB



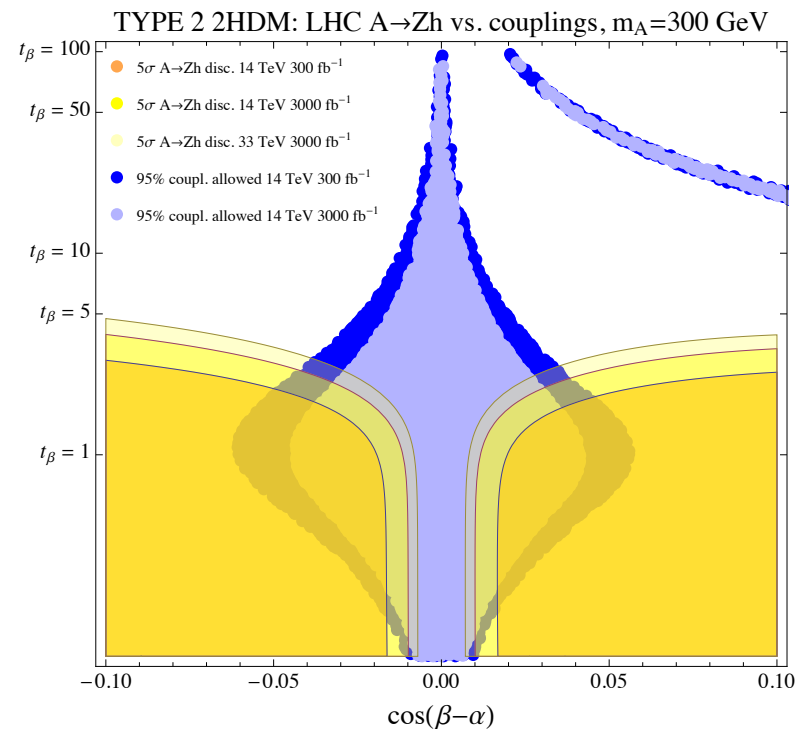
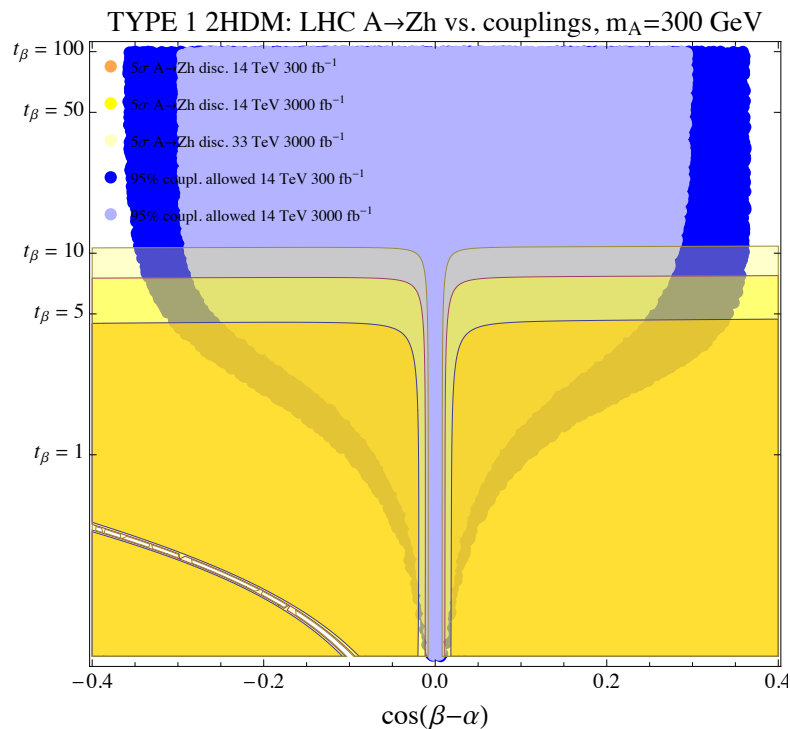
Motivate searches for additional Higgs bosons with significant mixing to 126 GeV state

BMSSM Higgs Signals

Models have many parameters, but signals described by simplified models with few parameters

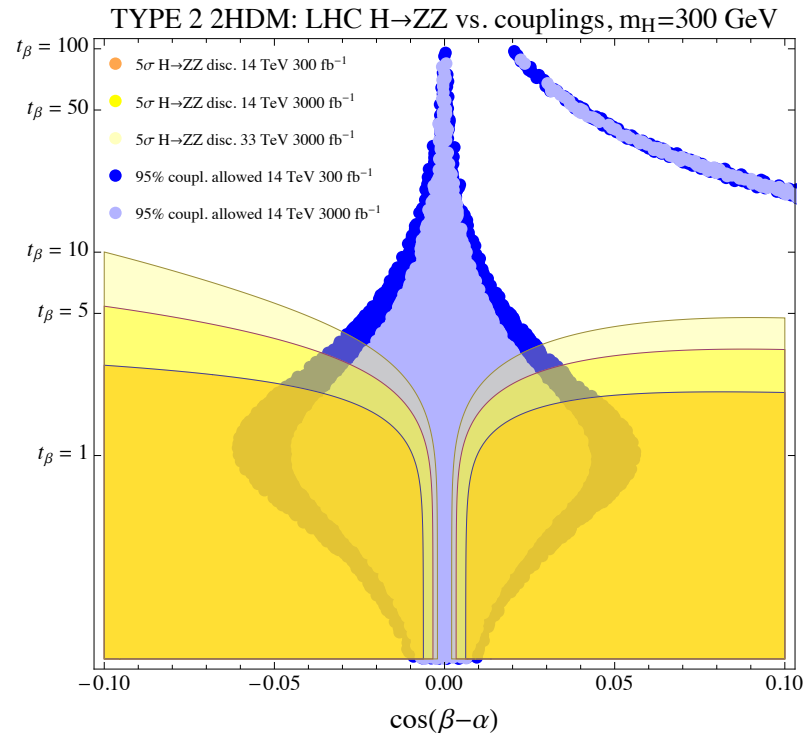
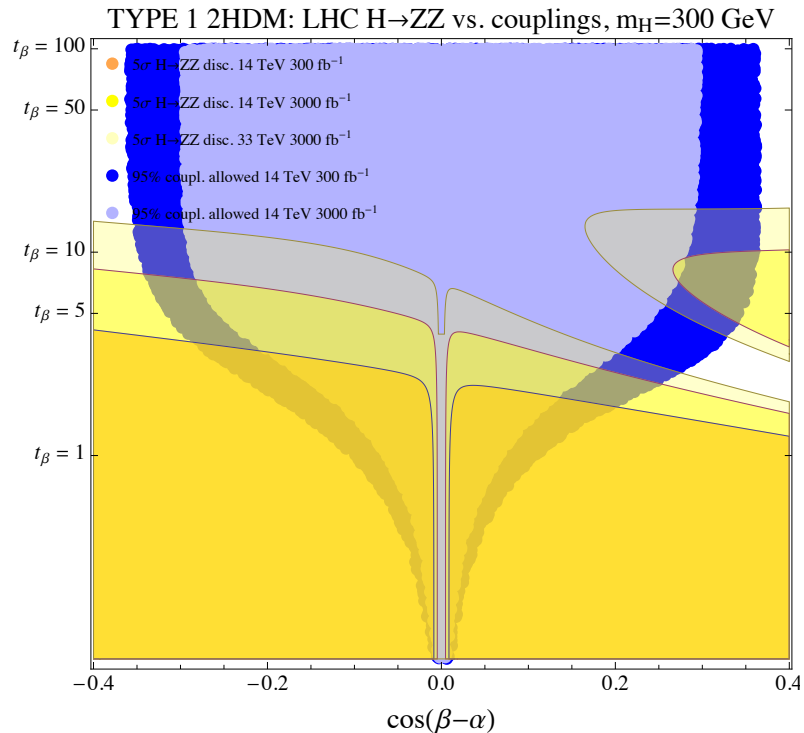
Novel feature: heavy Higgs fields can be far from decoupling limit

Example: $A \rightarrow Zh \rightarrow (\ell\ell)(bb)$ or $(\tau\tau)$



BMSSM Higgs Signals

Also: $H \rightarrow ZZ \rightarrow (\ell\ell)(\ell\ell)$



Other examples deserving further study:

$H/A \rightarrow tt, bb, \tau\tau$ $H \rightarrow hh$

$H \rightarrow AZ \rightarrow ZZh$

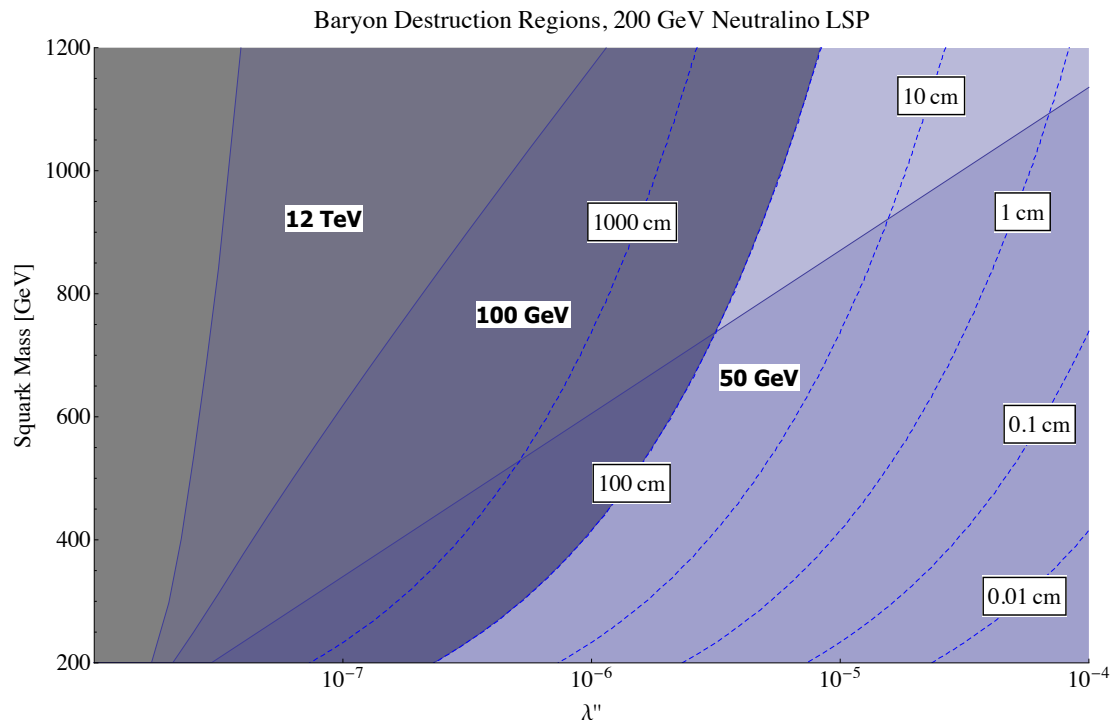
R-Parity Violation

UDD R-parity violation is often considered as a way to “hide” natural SUSY

Cosmology motivates displaced vertices

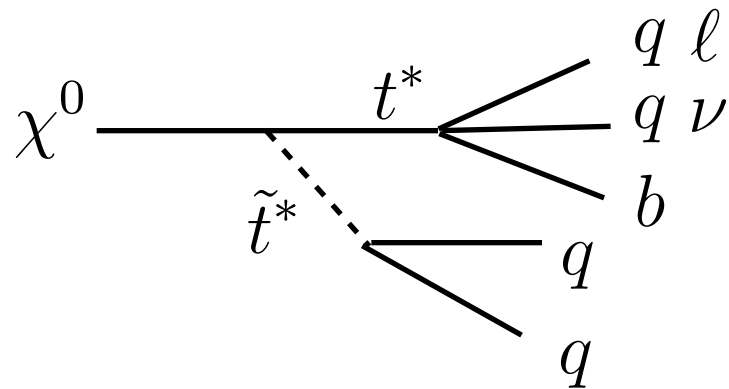
Barry, Graham, Rajendran [arXiv:1310.3853](https://arxiv.org/abs/1310.3853)

UDD destroys baryon asymmetry in early universe
⇒ requires low scale baryogenesis

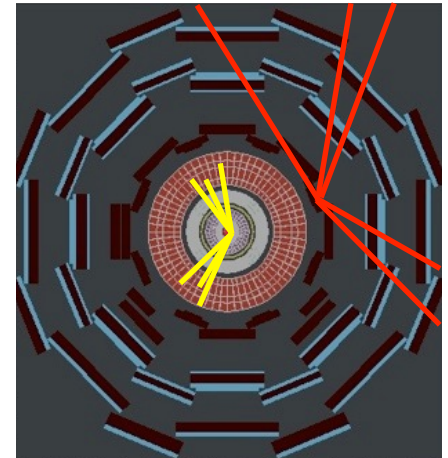


Displaced RPV

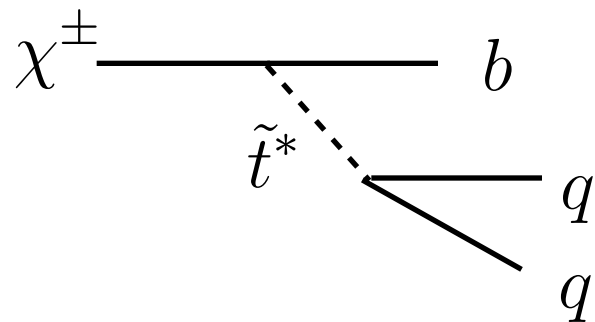
NLSP = neutralino



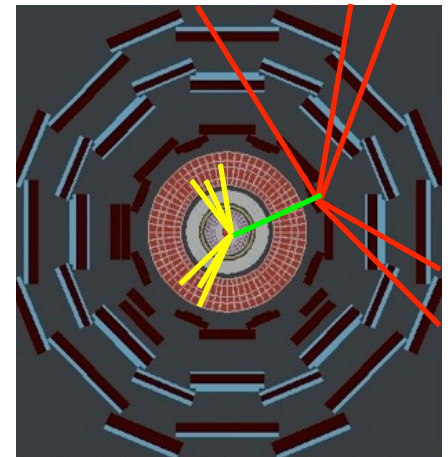
\Rightarrow displaced vertex



NLSP = chargino



\Rightarrow exploding track



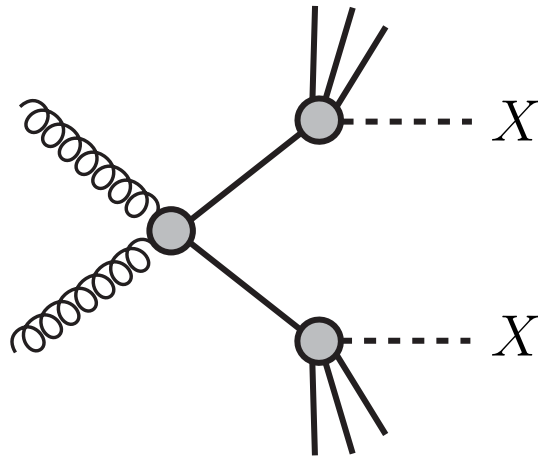
BMSSM Dark Matter

Dark matter exists

It requires physics beyond the standard model

But can we learn it's identity?

SUSY WIMP:



In the absence of signals in direct detection and collider searches, look at plausible mechanisms beyond WIMPs

⇒ new signals?

Hidden Sector Dark Matter

$$\Delta\mathcal{L} = \epsilon \mathcal{O}_{\text{vis}} \mathcal{O}_{\text{hidden}}$$

LSP in hidden sector \Rightarrow may be dark matter

LOSP (lightest ordinary supersymmetric particle)
decays to dark matter, may be long-lived

- SuperWIMP (Feng, Rajaraman, Takayama 2003)

LOSP dominates universe, then decays

$$\Omega_X = \frac{m_X}{m_{\text{LOSP}}} \Omega_{\text{LOSP}}$$

$X =$ gravitino or ...

Freeze-in Dark Matter

(Hall, Jedamzik, March-Russell, West 2009)

Relic density generated by LOSP decay

X never in equilibrium \Rightarrow no inverse decays

$$\Omega_X \simeq \frac{10^{27}}{g_*^{3/2}(m_{\text{LOSP}})} \frac{m_X \Gamma(\text{LOSP} \rightarrow X)}{m_{\text{LOSP}}^2}$$

$$c\tau(\text{LOSP} \rightarrow X) \sim 0.5 \text{ m} \left(\frac{m_X}{\text{keV}} \right) \left(\frac{m_{\text{LOSP}}}{100 \text{ GeV}} \right)^{-2}$$

$m_X \gtrsim \text{keV}$ (structure formation)

\Rightarrow motivates highly displaced decays

Asymmetric Dark Matter

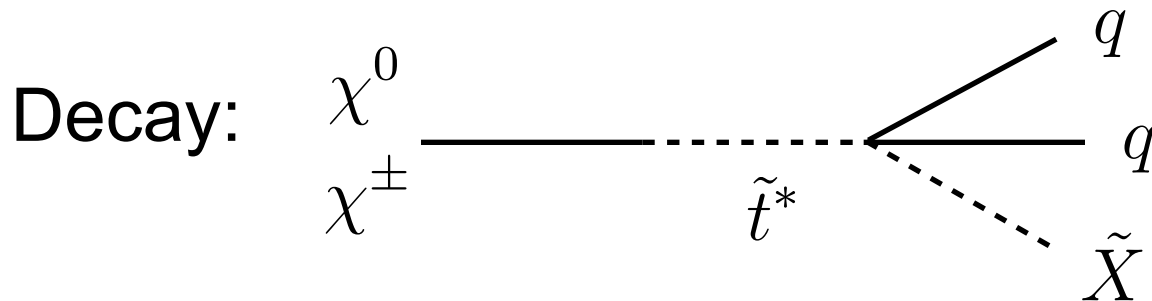
(Kaplan, ML, Zurek 2009)

$$\Delta W = \frac{1}{M^2} \underbrace{UDD}_{\text{visible}} \underbrace{X^2}_{\text{hidden}}$$

Transfers B-L asymmetry to dark matter

Freeze out: $T_f \gtrsim m_X \Rightarrow \Omega_X \sim \frac{m_X}{m_p} \Omega_B \Rightarrow m_X = 14 \text{ GeV}$

Valid for $M \gtrsim \text{TeV}$ due to rapid decoupling for $T < m_{\tilde{q}}$

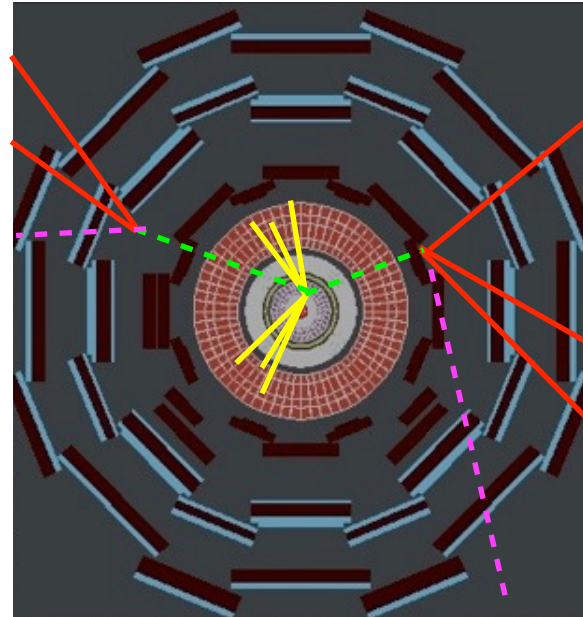
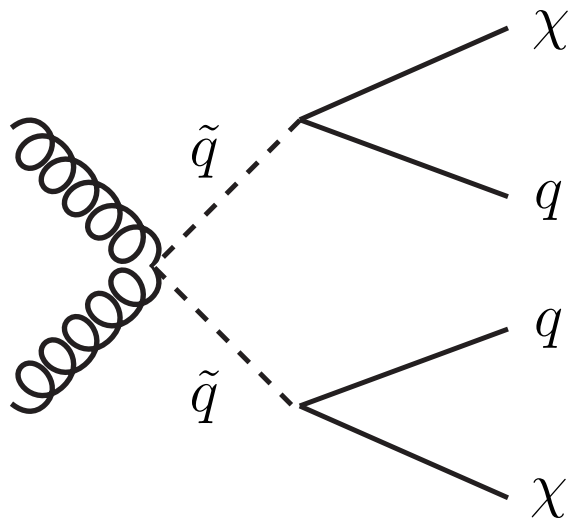


$$c\tau \sim \text{cm} \left(\frac{M}{\text{TeV}} \right)^4 \left(\frac{m_{\tilde{t}}}{500 \text{ GeV}} \right)^4 \left(\frac{m_\chi}{100 \text{ GeV}} \right)^{-9}$$

Displaced Optimism

Displaced vertices are exciting discovery mode...

...and open the possibility of fully reconstructing SUSY events



LOSP = charged: SUSY fully reconstructed

LOSP = neutral: \vec{v}_{LOSP} from vertex position + timing

Constraining/measuring dark matter mass possible

Conclusions

- Impressive breadth and depth of SUSY searches
- Both tuned and natural versions of SUSY are still plausible and important to test
- BMSSM Higgs and displaced vertices are important to search for

Decoupling Limit

Additional MSSM Higgs generically near decoupling limit

$$\begin{pmatrix} H \\ H_{\perp} \end{pmatrix} = \begin{pmatrix} \sin \beta & \cos \beta \\ \cos \beta & -\sin \beta \end{pmatrix} \begin{pmatrix} H_u \\ H_d^{\dagger} \end{pmatrix} \quad \tan \beta = \frac{v_u}{v_d}$$

$$\underbrace{\langle H \rangle}_{\text{the Higgs field}} = \frac{v}{\sqrt{2}} \quad \langle H_{\perp} \rangle = 0$$

the Higgs field

126 GeV mass eigenstate $\simeq H_{\parallel}^0$

\Rightarrow additional Higgs fields decouple from W, Z
(but not fermions)

Simplified Models

Simplified models for BSM/BMSSM Higgs searches

For SUSY interpretation, assume decoupling

$$H = H_u \sin \beta + H_d^\dagger \cos \beta$$

H_\perp decoupled

⇒ SM Higgs with one additional parameter $\tan \beta$
governs couplings of H to fermions

...now add BMSSM Higgs

Higgs + Singlet

$$V_{\text{eff}} = m_H^2 |H|^2 + \lambda_H |H|^4 \\ + \frac{1}{2} m_\Phi \Phi^2 + \frac{1}{4} \lambda_\Phi \Phi^4 + \frac{1}{2} \lambda_{\Phi H} S^2 |\Phi|^2$$

6 parameters: $m_1, m_2, v, \tan \beta, \cos \gamma, \langle \Phi \rangle$

$$\begin{pmatrix} h \\ s \end{pmatrix} = \begin{pmatrix} \cos \gamma & \sin \gamma \\ -\sin \gamma & \cos \gamma \end{pmatrix}$$

Model 1: $\Phi = \text{CP even}$ $m_1 = 126 \text{ GeV}$

Production and decay of h_2 governed by m_2, β, γ

$$h_2 \rightarrow ZZ, WW, \bar{t}t, h_1 h_1$$

Model 2: $\Phi = \text{CP odd PNGB}$

$$m_2 = 126 \text{ GeV}$$

$$\langle \Phi \rangle = 0$$

$$h_1 \rightarrow aa \quad a \rightarrow \bar{b}b, \tau^+ \tau^-$$