

Searching for Dark Matter in a Bubble Chamber

Hugh Lippincott, Fermilab

Seminar at University of California, Davis
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COUPOP

Collaboration

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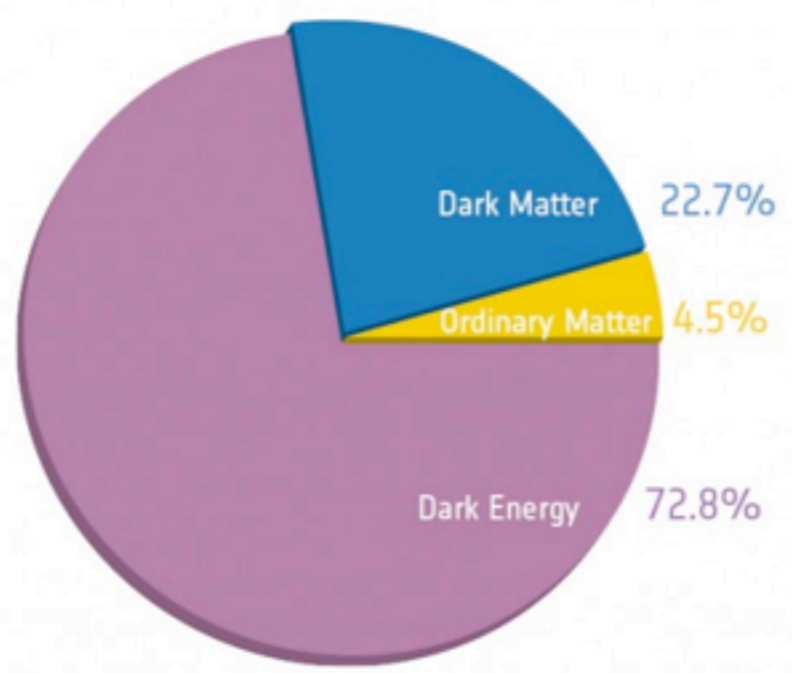
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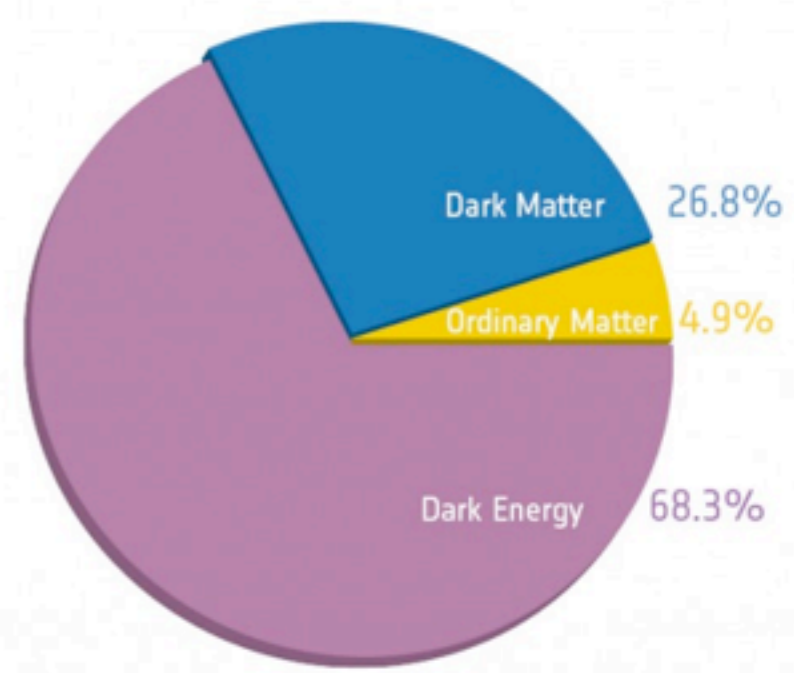


There is pretty strong consensus regarding how much stuff there is in the universe

By that same consensus, we only understand 5% of it



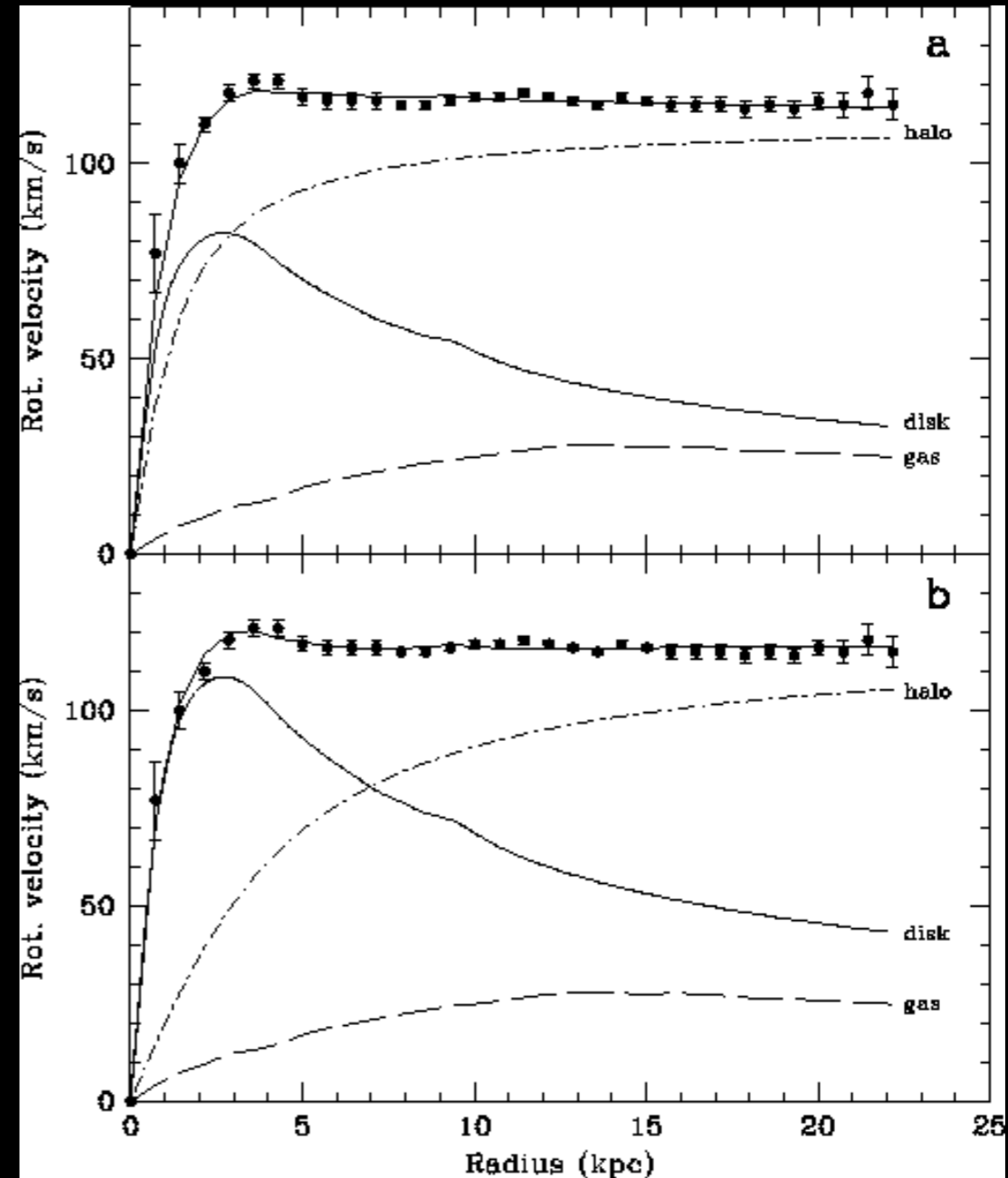
Before Planck



After Planck

Dark matter - evidence?

- Galaxy rotation curves



Dark matter - evidence?

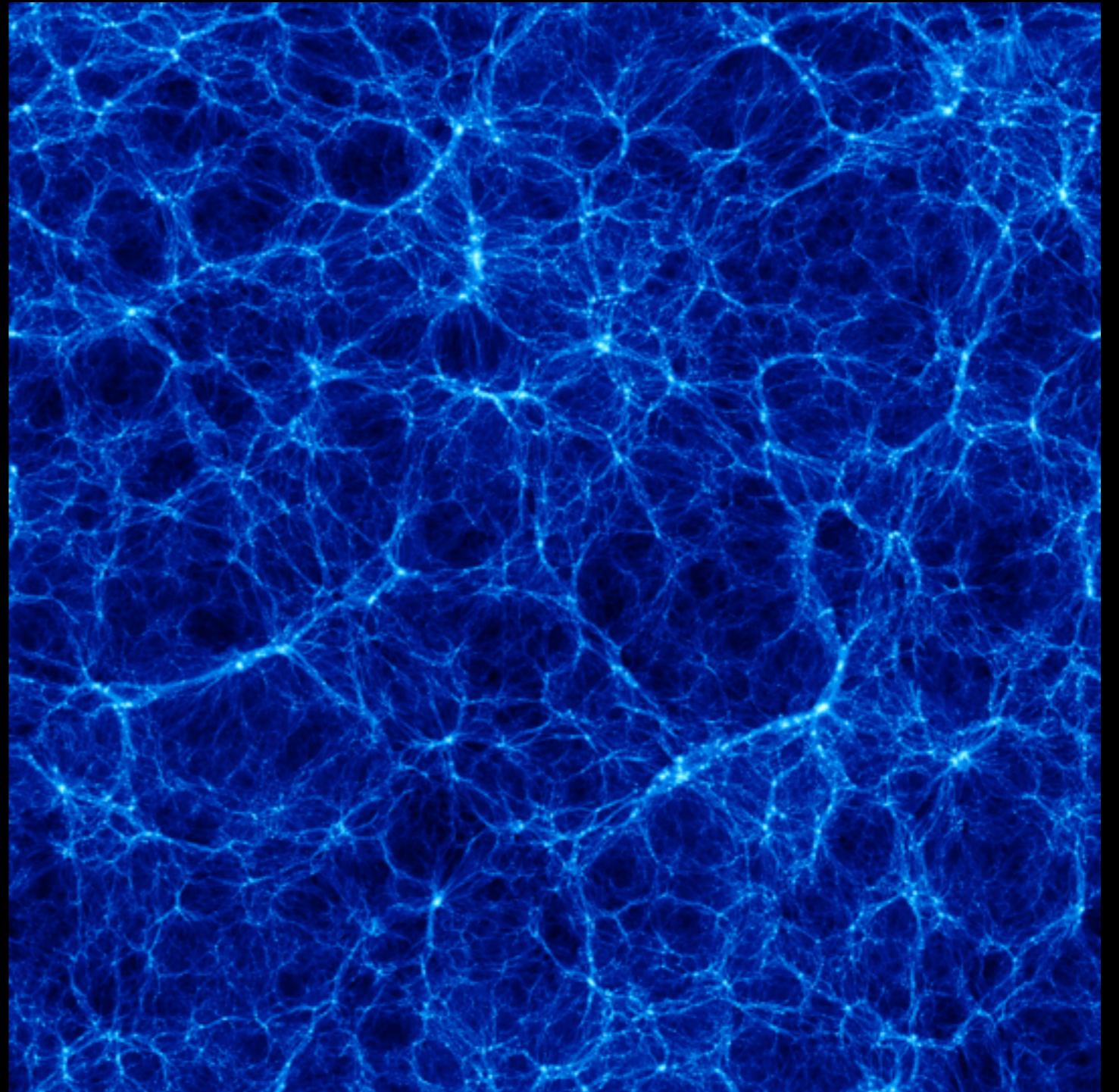
- Galaxy rotation curves
- Galaxy clusters



Fritz Zwicky, 1930

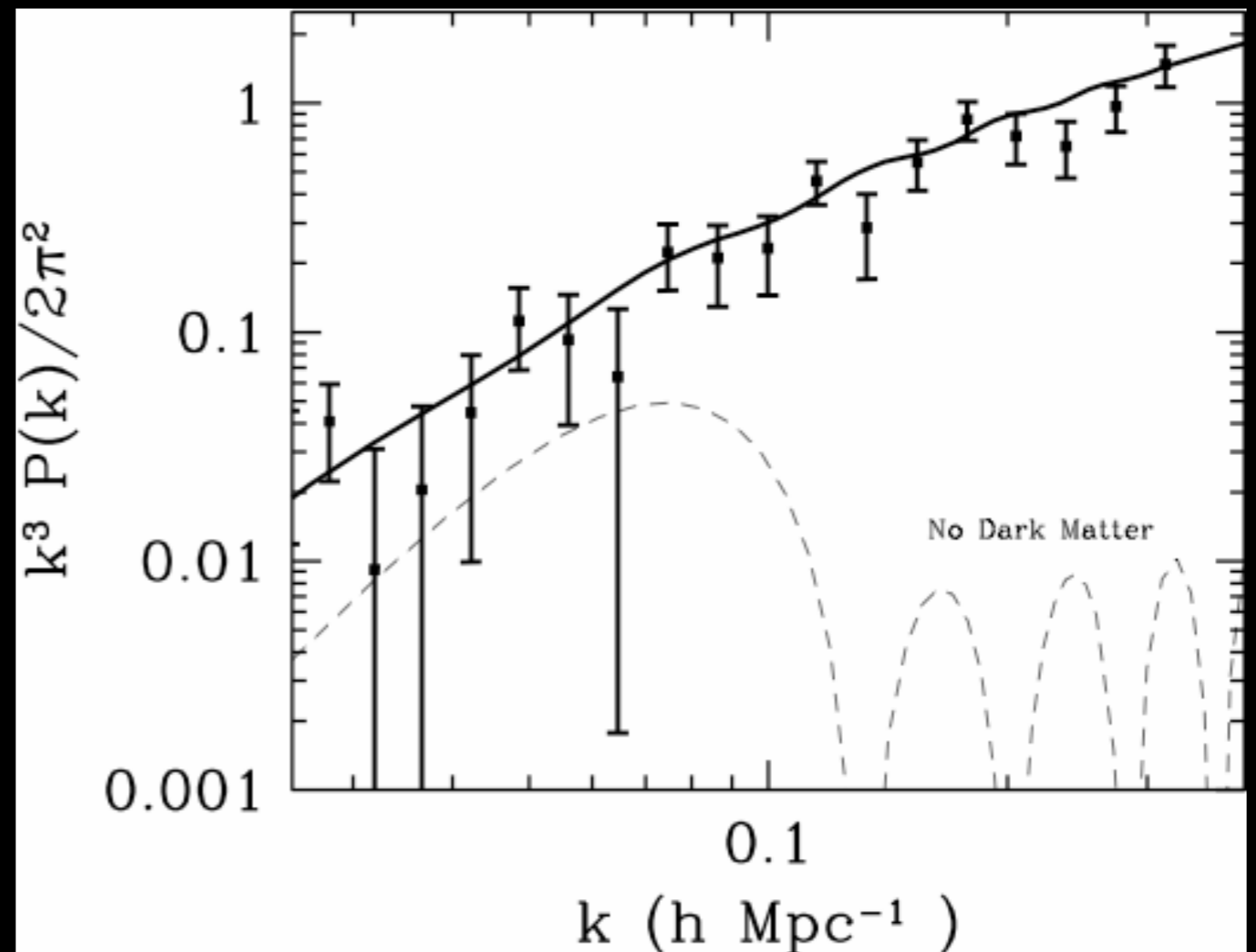
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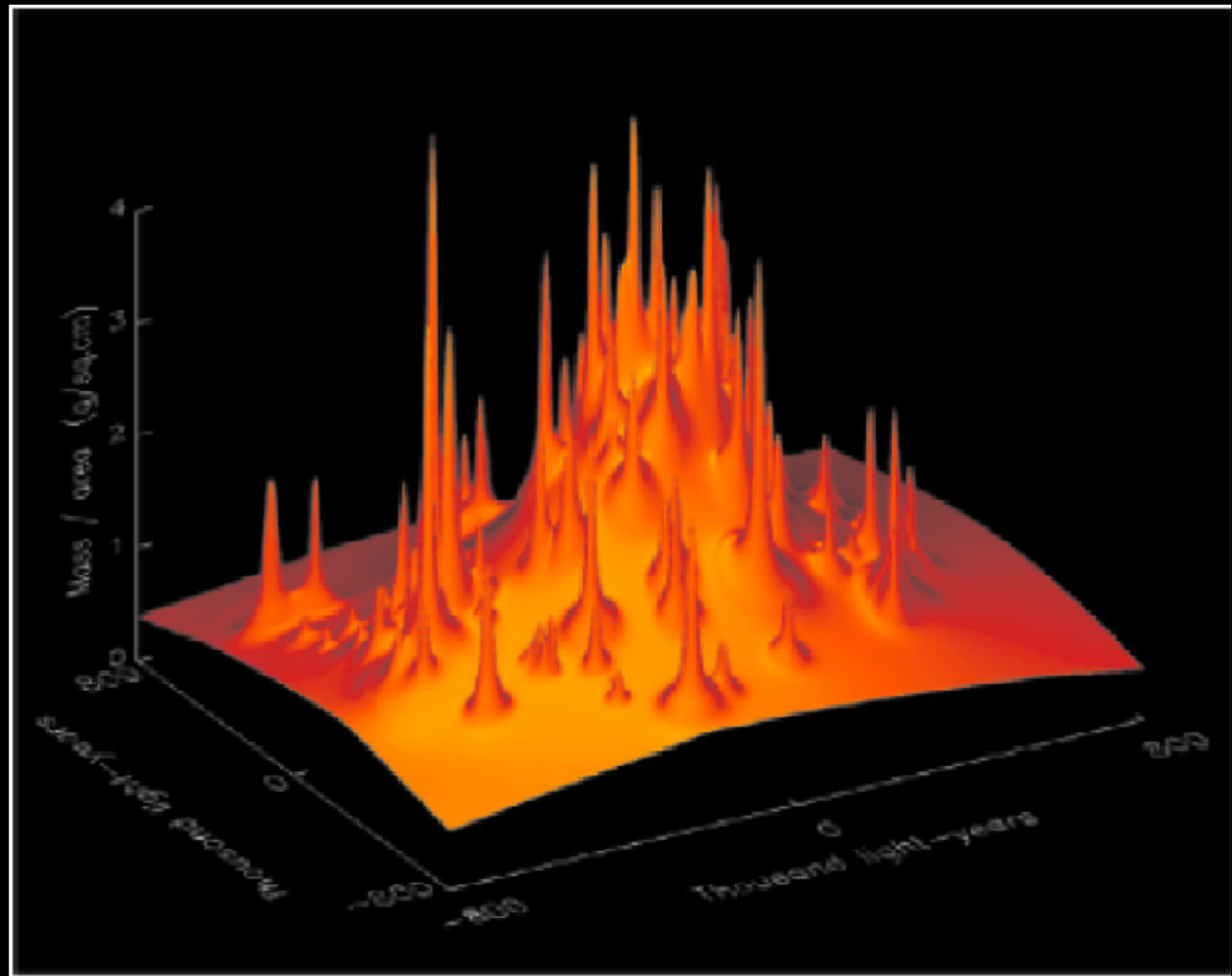
Dark matter - evidence?

- Galaxy rotation curves
- Galaxy clusters
- Gravitational lensing



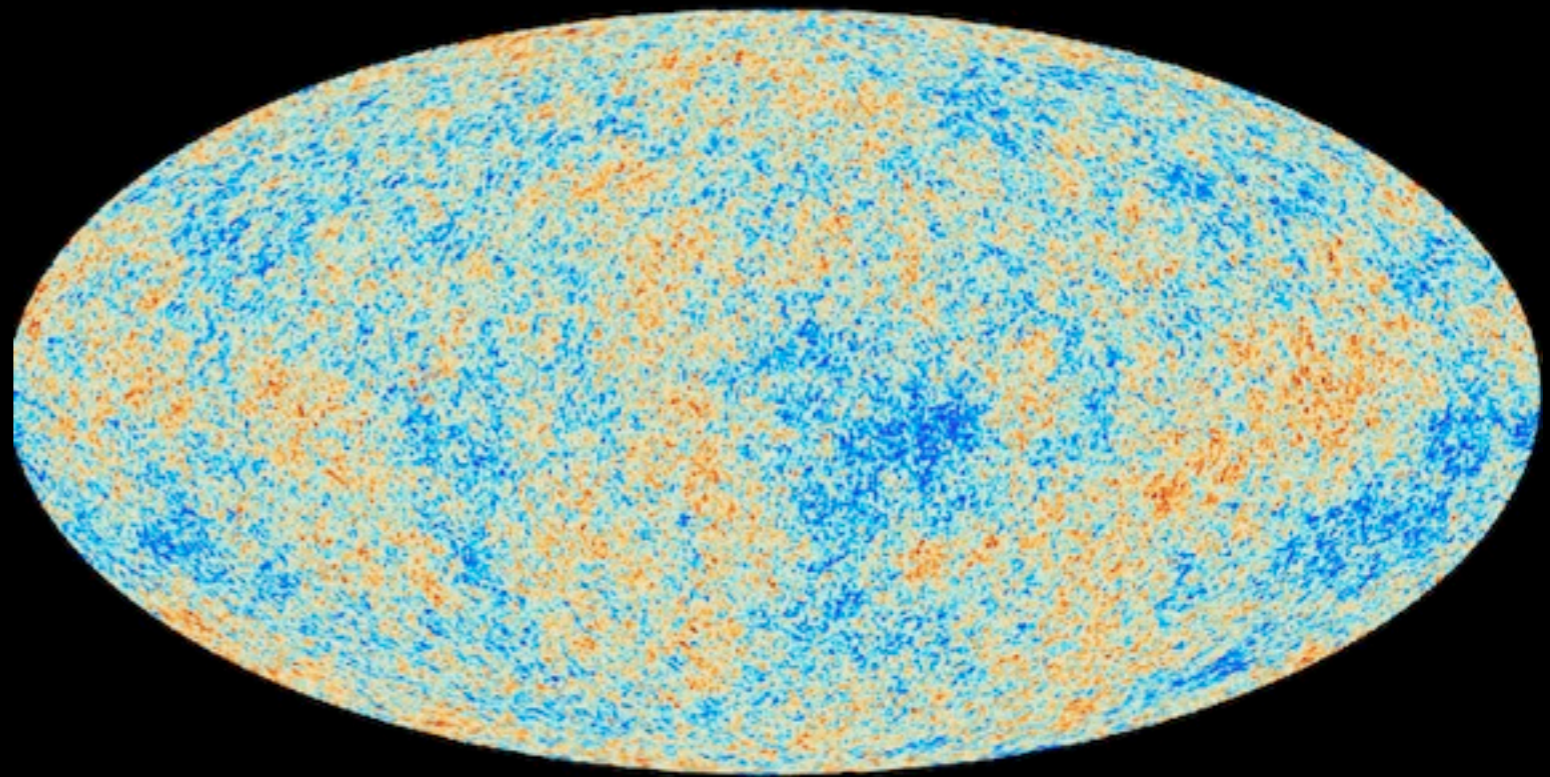
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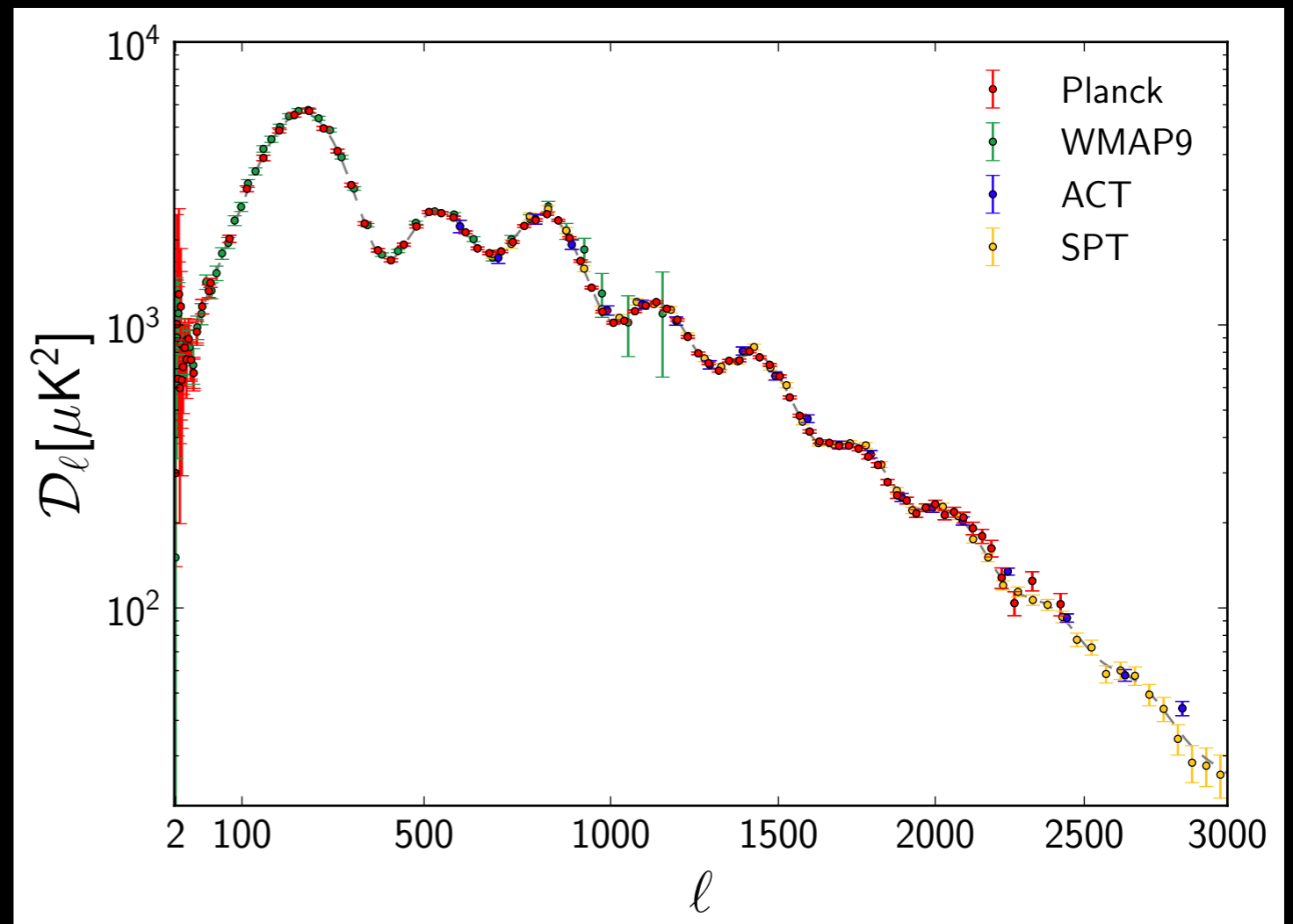
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Dark matter - evidence?

- Galaxy rotation curves
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- Gravitational lensing
- Cosmic microwave background
- Galactic collisions



So what is it?

- We know it interacts gravitationally
- It is “dark” - should not interact with light or electromagnetism
- Nearly collisionless
- Slow

Axions

Champs

Kaluza-Klein particles

Many more

WIMPs, WIMPzillas,
Light WIMPS

MACHOs

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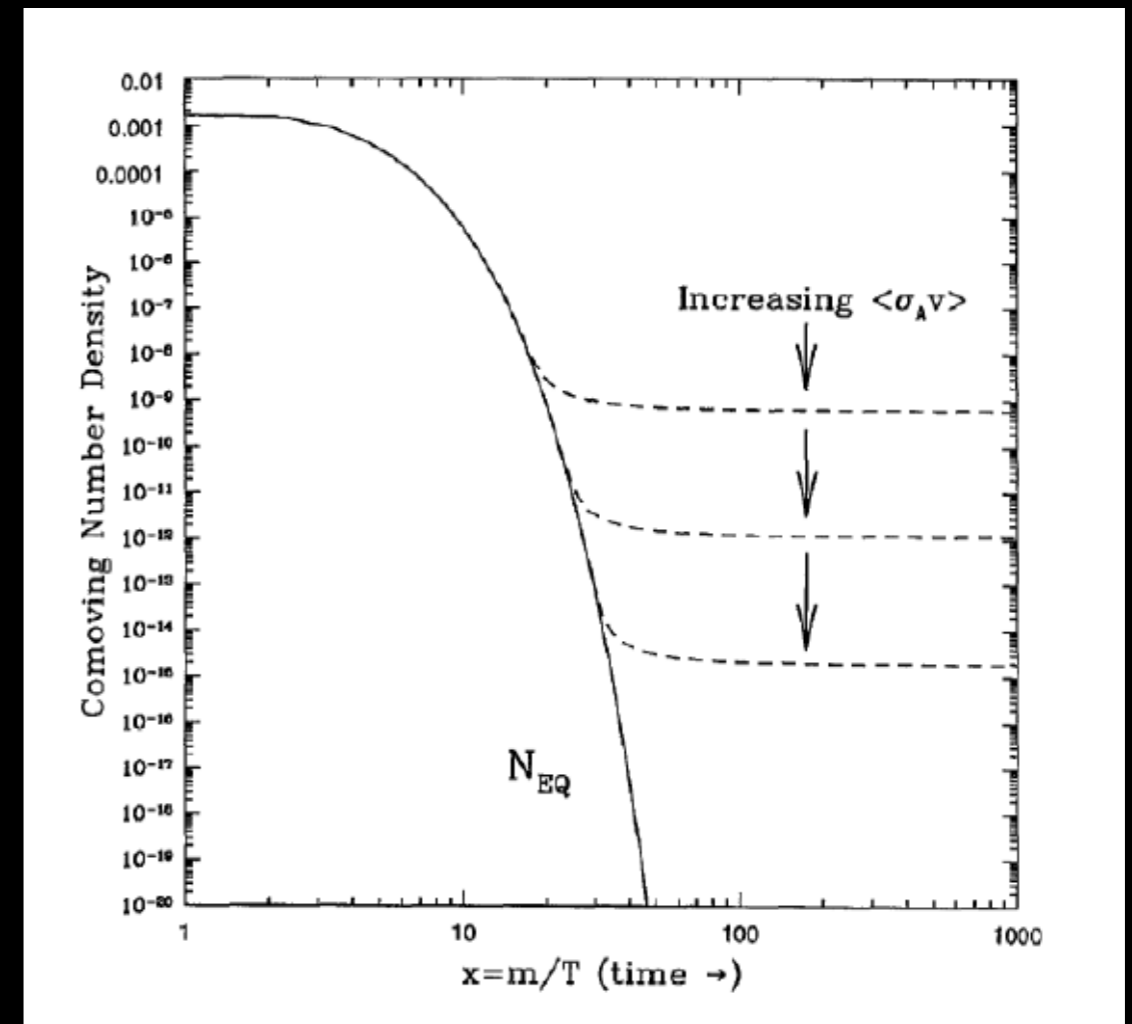


Beyond the Standard Model!



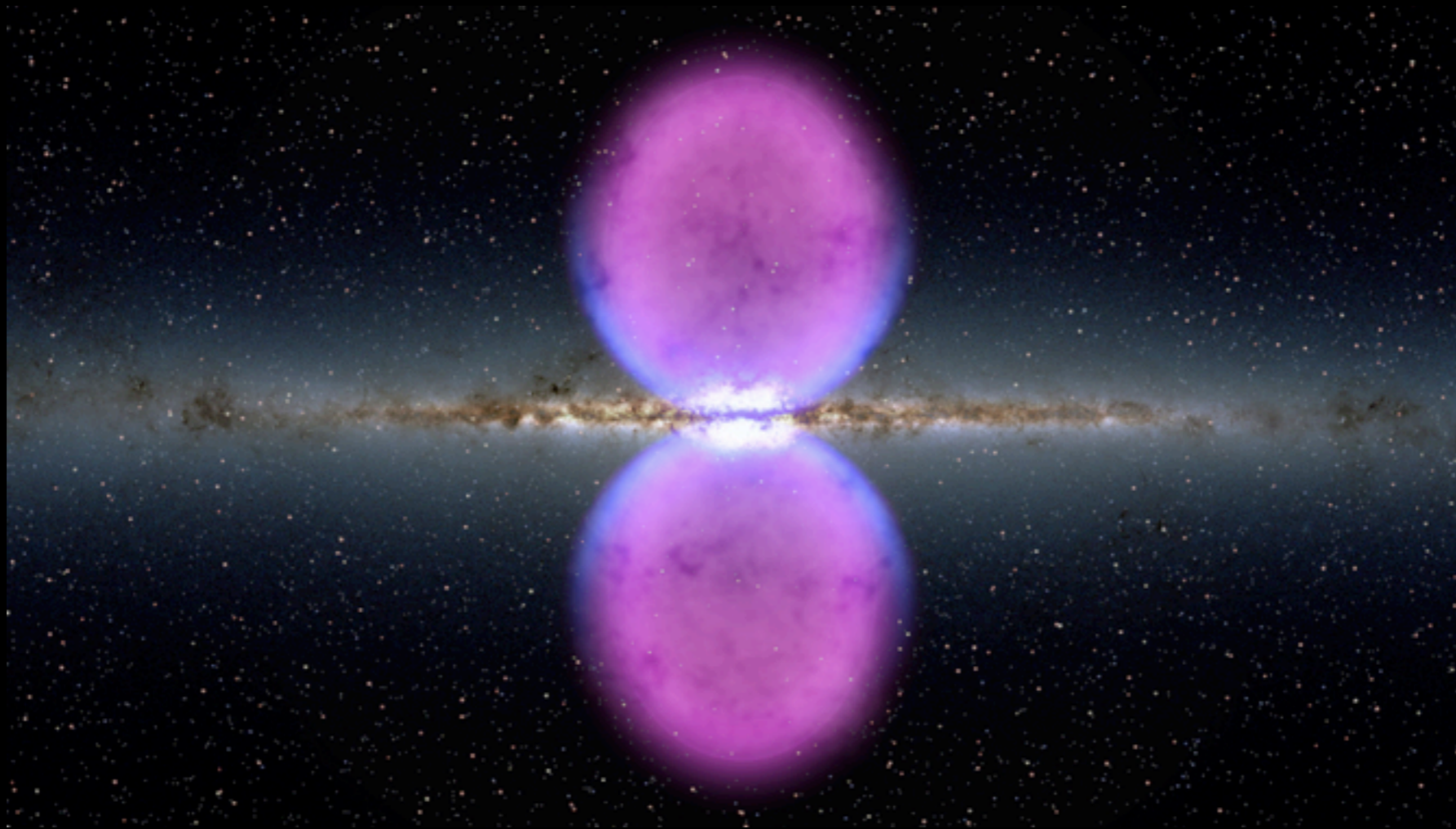
WIMPs

- Most discussed candidate is Weakly Interacting Massive Particle
 - Produced during big bang
 - Decouples from ordinary matter as the universe expands and cools
 - Still around today with densities of about a few per liter
- Supersymmetry produces a theoretical candidate (LSP), but others exist (e.g. Kaluza-Klein particles, ...)



How do we find it?

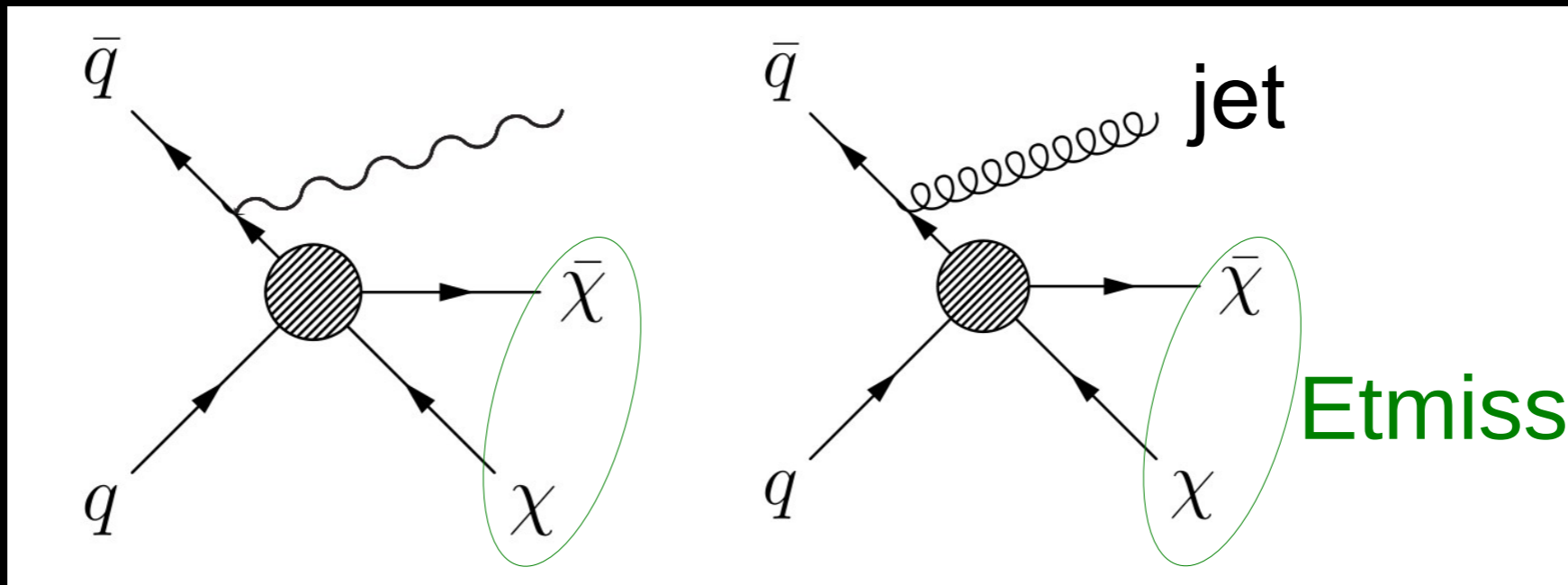
- Indirect - detect annihilation products from regions of high density like the sun or the center of the galaxy



Fermi bubbles, courtesy of NASA

How do we find it?

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- Accelerators - create a WIMP at the LHC
- Missing ET and monojet searches



How do we find it?

- Indirect - detect annihilation products from regions of high density like the sun or the center of the galaxy
- Accelerators - create a WIMP at the LHC
 - Missing ET and monojet searches
- Direct detection - WIMPs can scatter elastically with nuclei and the recoil can be detected

$$\frac{dR}{dQ} = \frac{\rho_0}{m_\chi} \times \frac{\sigma_0 A^2}{2\mu_p^2} \times F^2(Q) \times \int_{v_m} \frac{f(v)}{v} dv$$

Rate calculation

- ▶ The differential cross section (for spin-independent interactions) per kilogram of target mass per unit recoil energy is

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- ▶ The velocity distribution of dark matter in the galaxy - of order 30% uncertainty, and $v_m = \sqrt{Q/2m_r^2}$

The energy scale

- Energy of recoils is tens of keV
- Entirely driven by kinematics, elastic scattering of things with approximately similar masses (100 GeV) and $v \sim 0.001c$

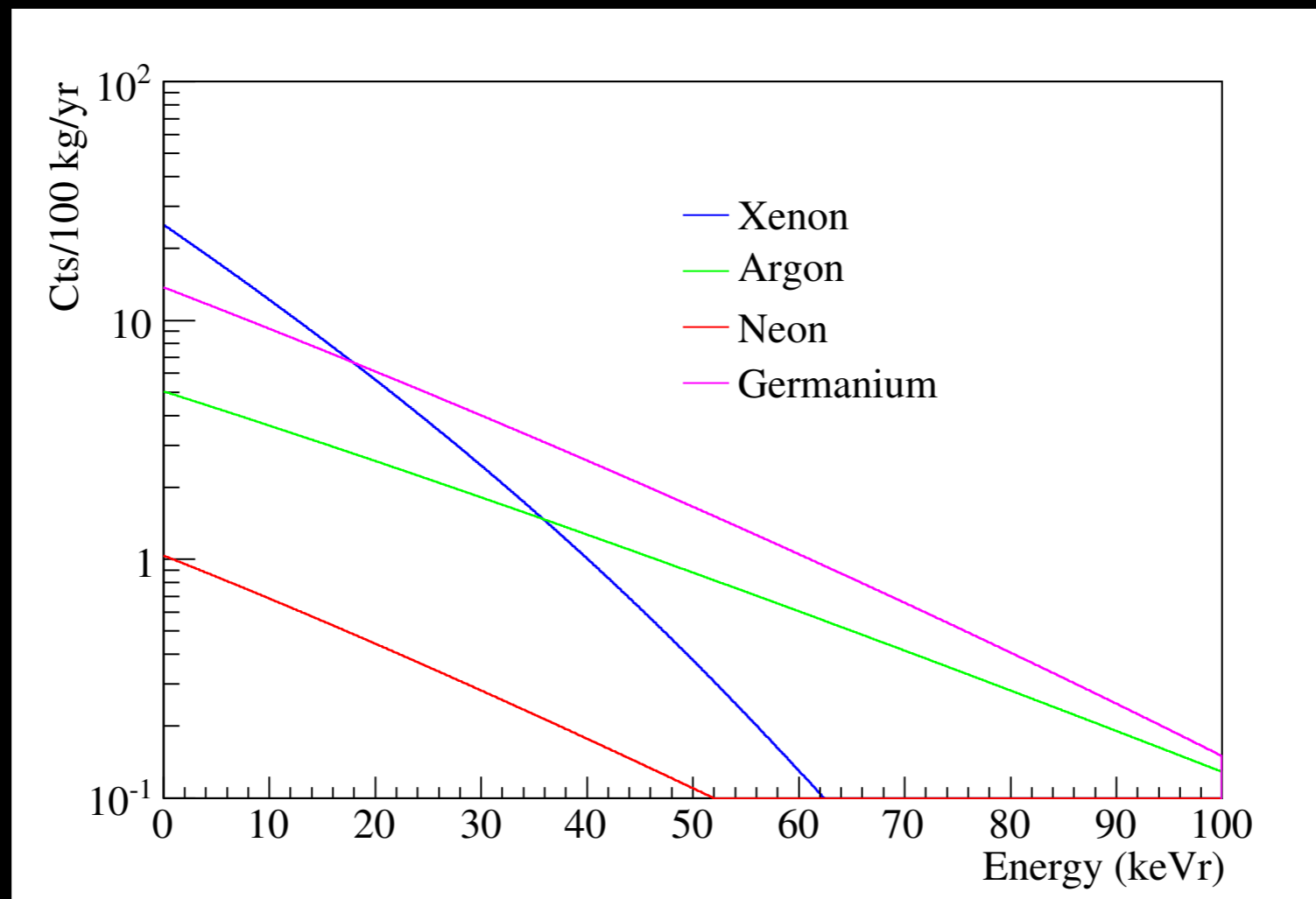
$$\frac{1}{2}m_N v_N^2 = \frac{1}{2} \times 100 \text{ GeV} \times 10^{-6} = 50 \text{ keV}$$



Rate calculation

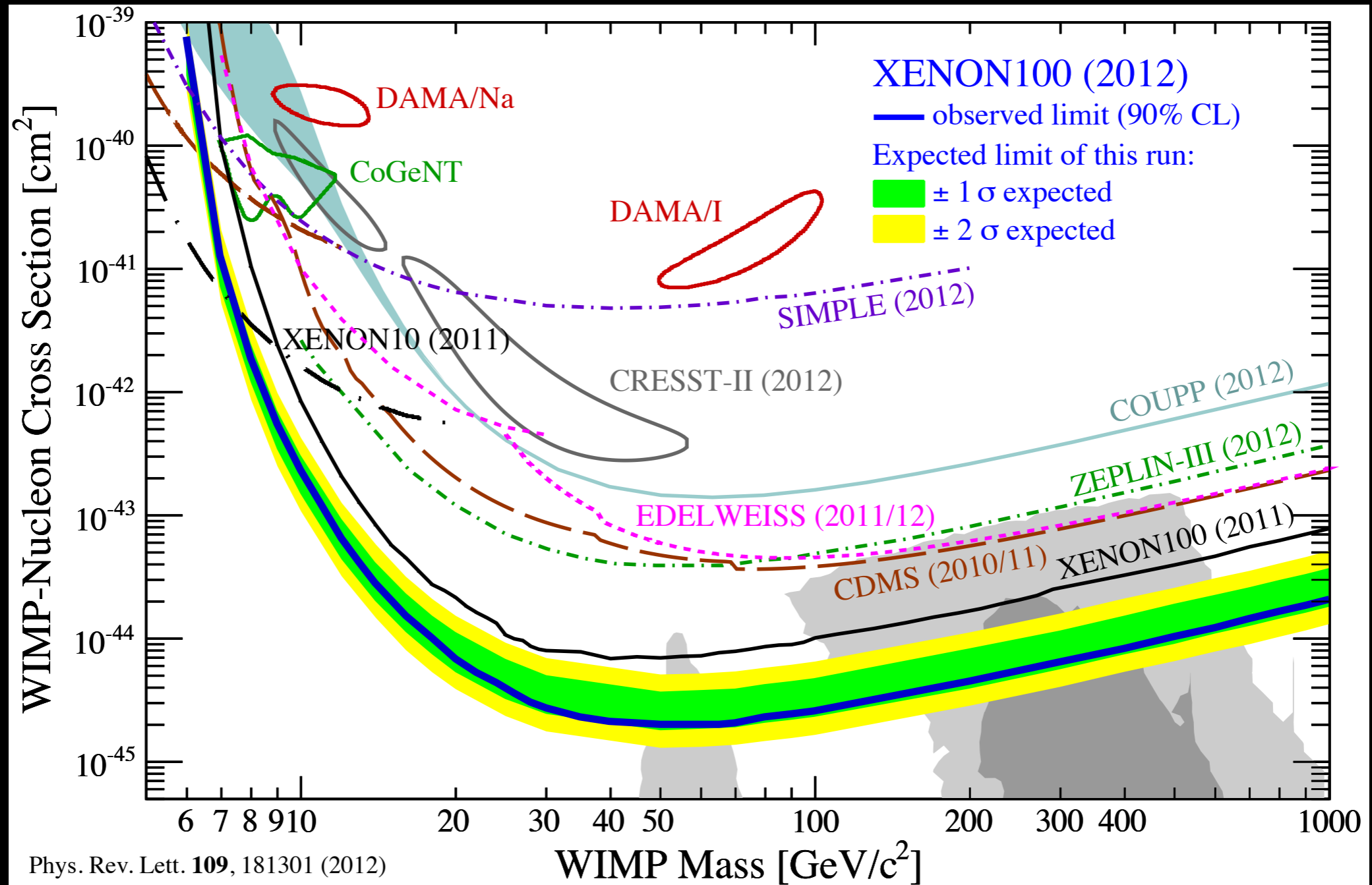
- ▶ Integrated rate above threshold, 100 GeV WIMP, $\sigma_0 = 10^{-45} \text{ cm}^2$

$$I = \int_{Q_{\text{thresh}}} dQ \frac{dR}{dQ} = \int_{Q_{\text{thresh}}} dQ \frac{\rho_0}{m_\chi} \frac{\sigma_0 A^2}{2\mu_p^2} F^2(Q) \int_{v_m} \frac{f(v)}{v} dv$$



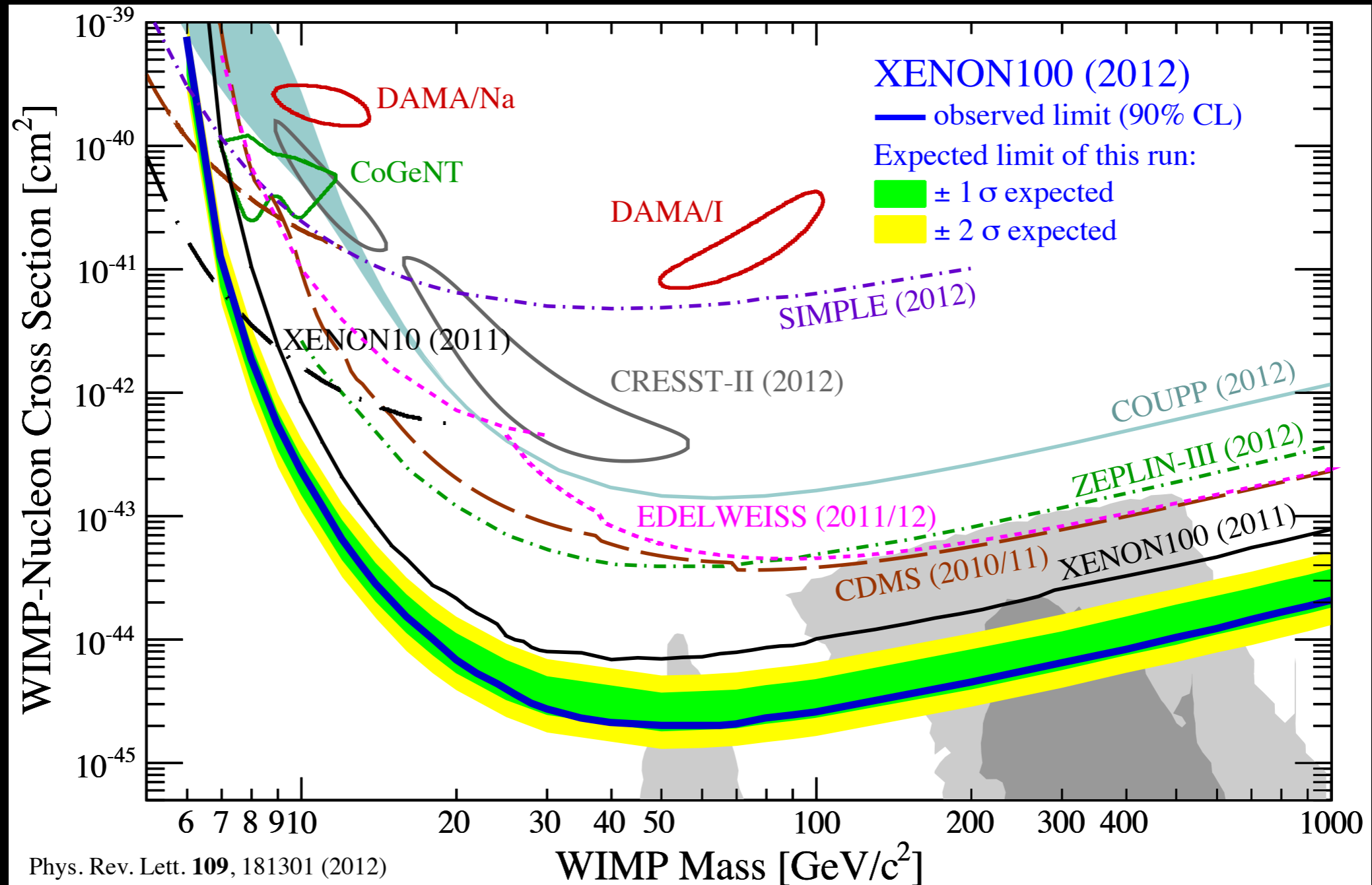
- ▶ Looking for a handful of events

The canonical plot



- Limited at low mass by detector threshold
- Limited at high mass by density

The canonical plot



- What happened to “weakly” interacting?
- Mediation via Z was excluded long ago ($\sim 10^{-39} \text{ cm}^2$), but only now are we probing Higgs exchange

So we look for WIMPs

- A few hundred just passed through us, and we might expect a handful of counts in a detector per year

So we look for WIMPs

- A few hundred just passed through us, and we might expect a handful of counts in a detector per year
- The problem is that background radioactivity is **everywhere!**



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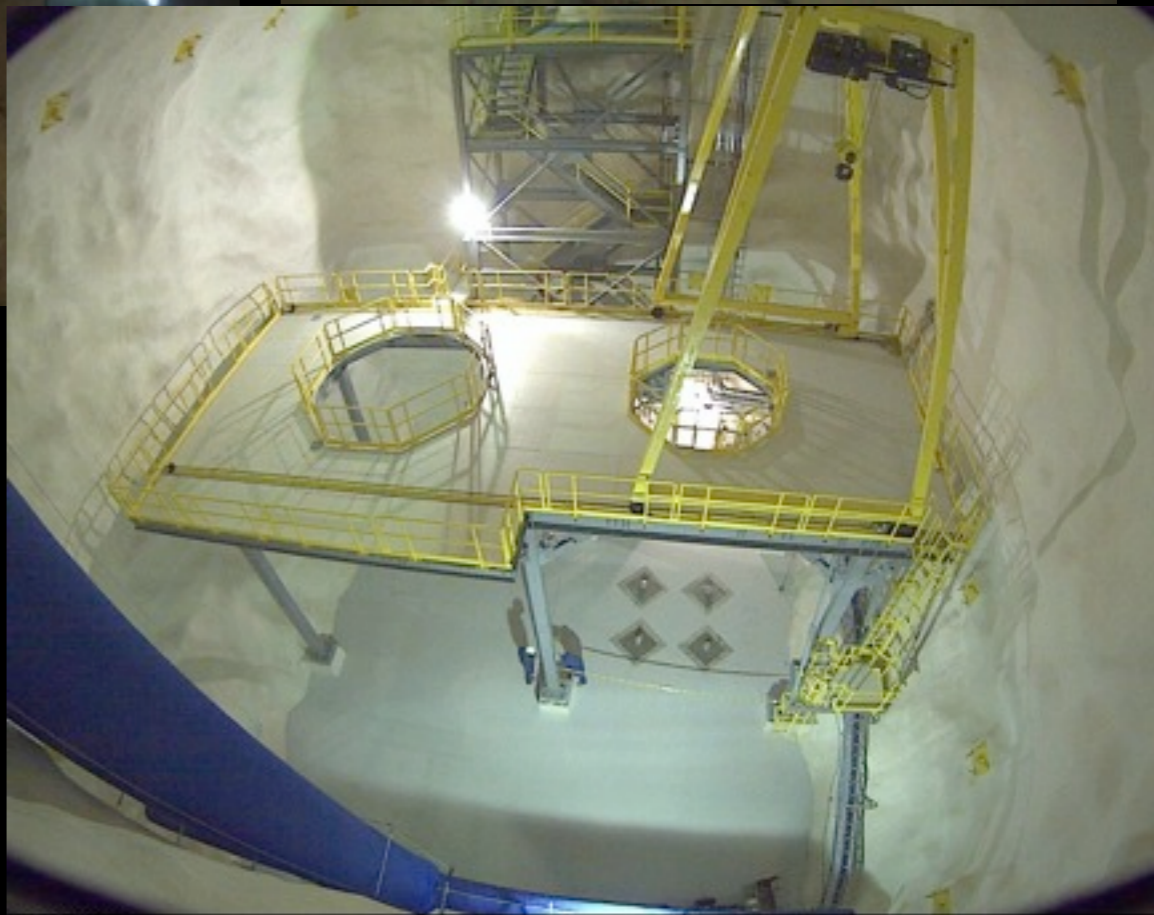
100 events/second/kg =
3,000,000,000,000 events/year
in a ton-scale experiment

Backgrounds!



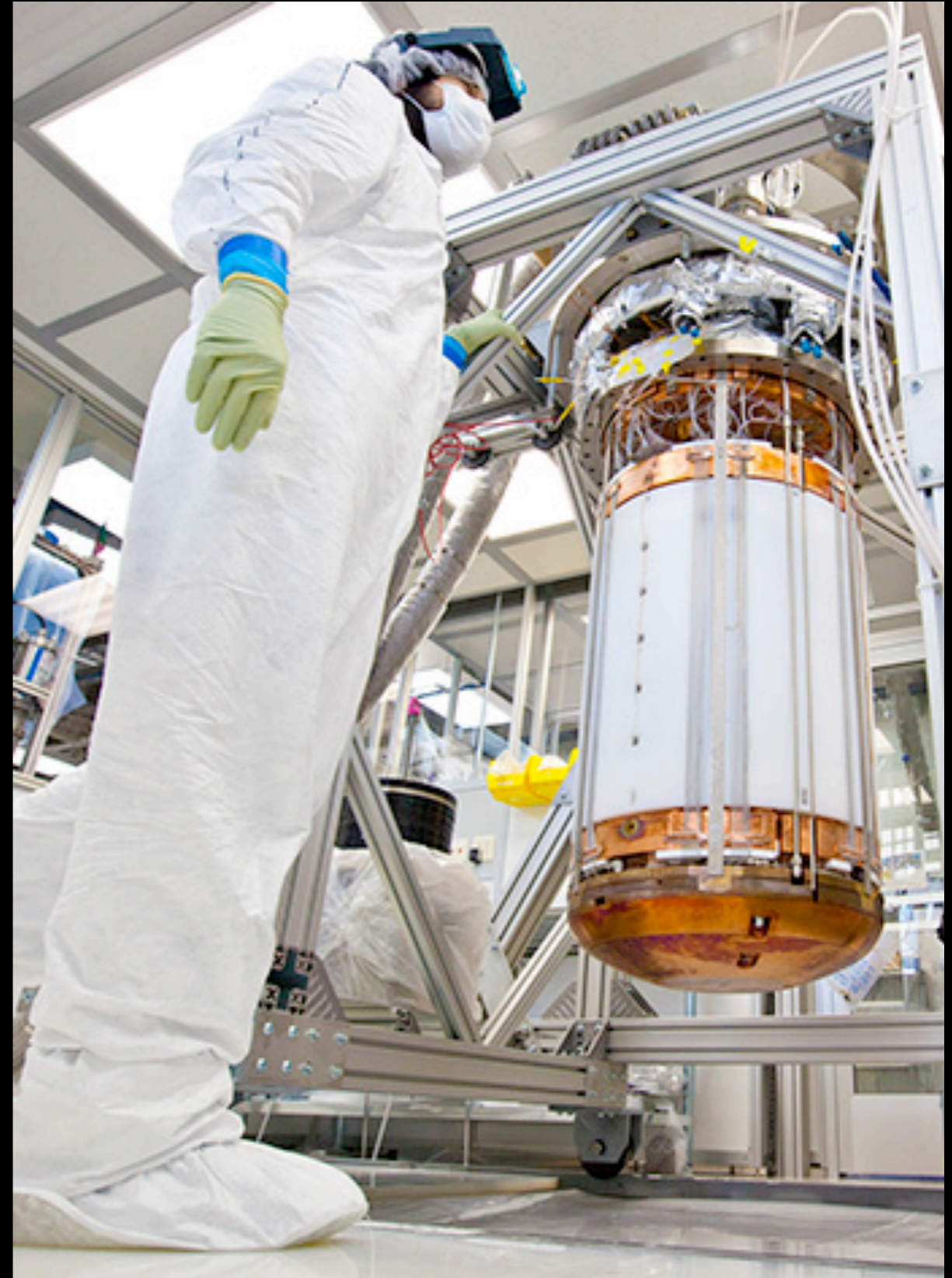
Background sources

- Cosmic rays are constantly streaming through
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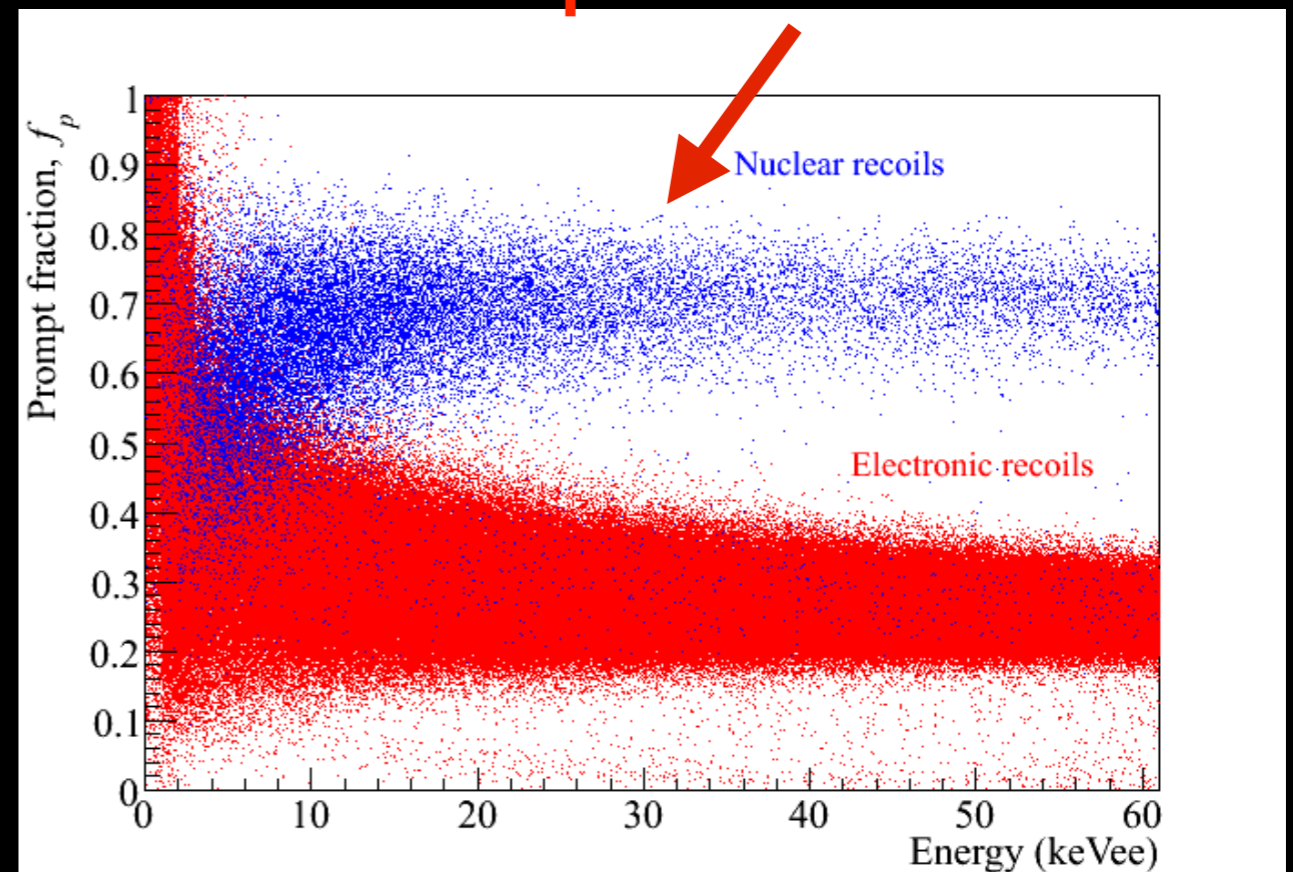
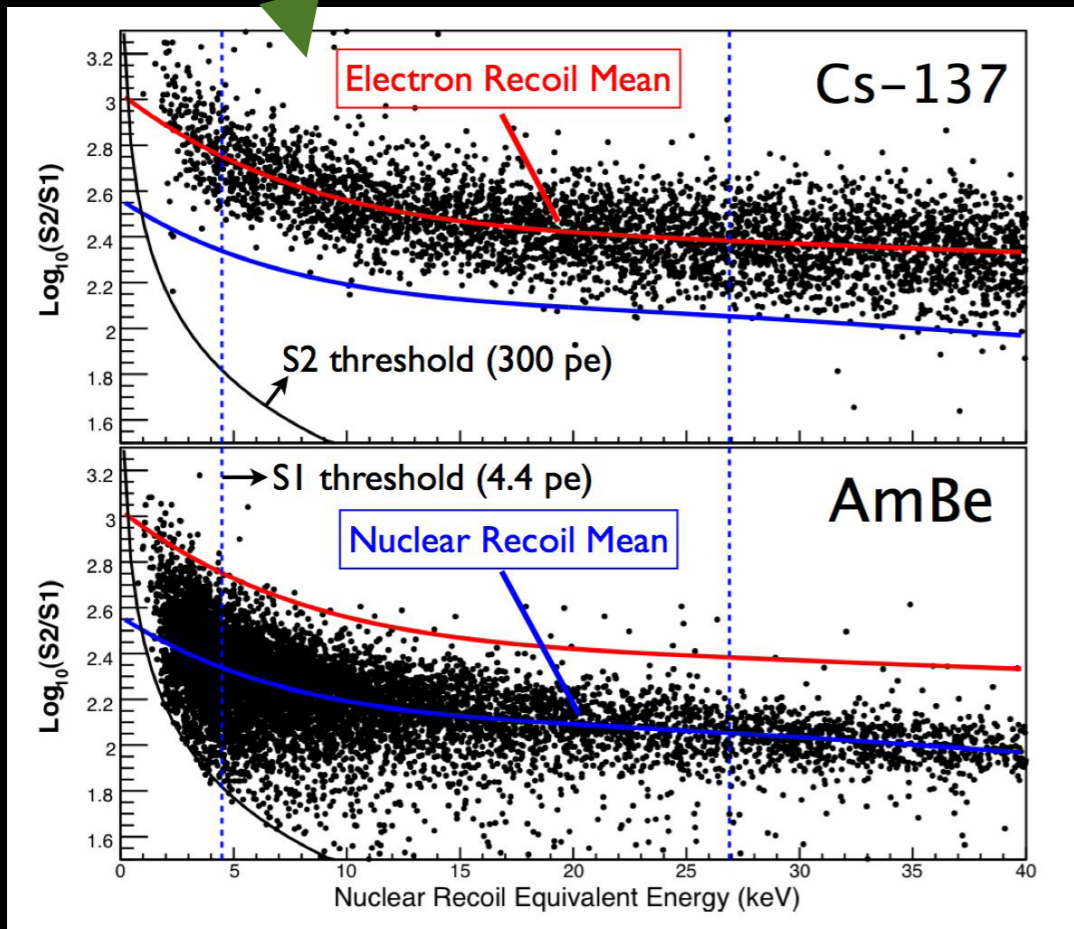
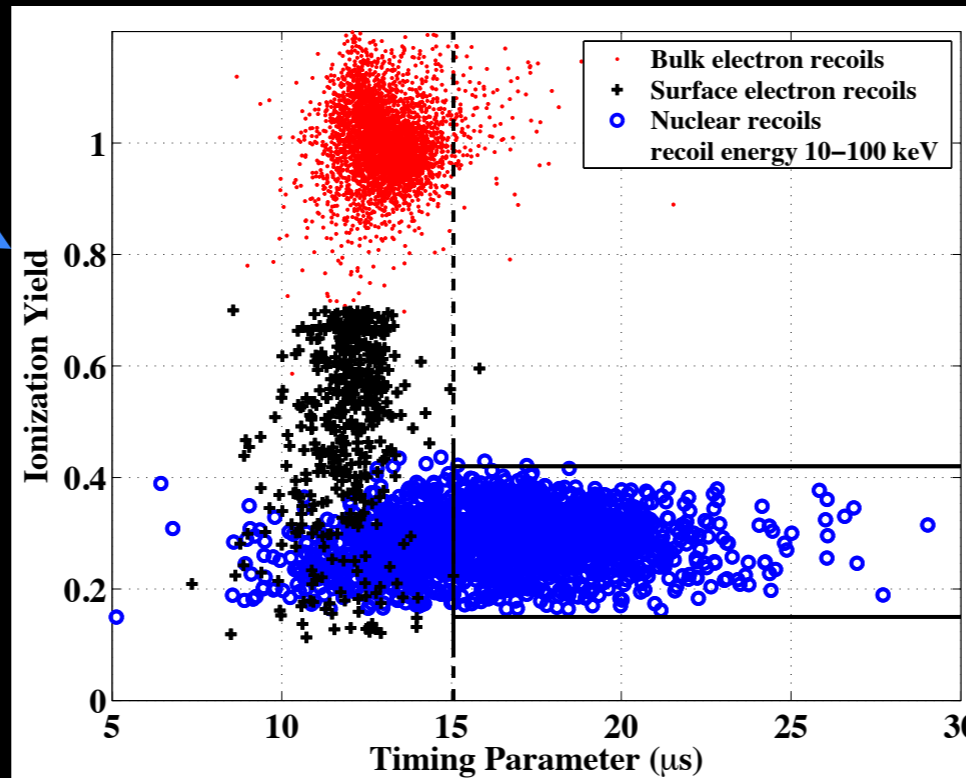
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- The detector itself - steel, glass, detector components
 - Self-shielding to leave a clean inner region
 - Discrimination - can you tell signal from background (gamma rays, alphas, neutrons, etc)?

CDMS - Charge to heat

Xenon TPCs - Charge to light

Electronic recoils (gammas) vs. nuclear recoils (WIMPs)

Argon - Pulse shape discrimination



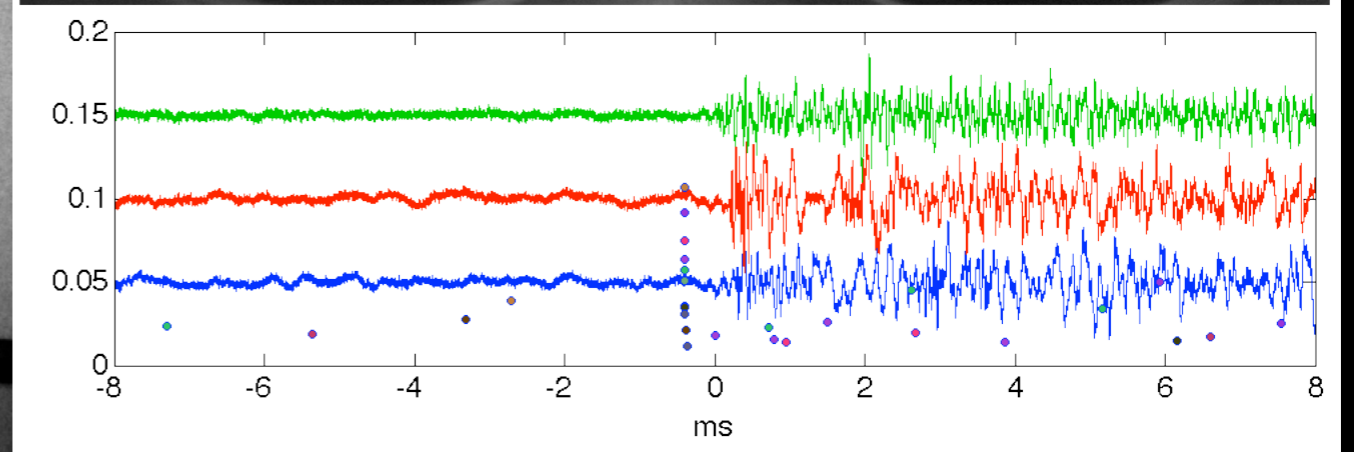
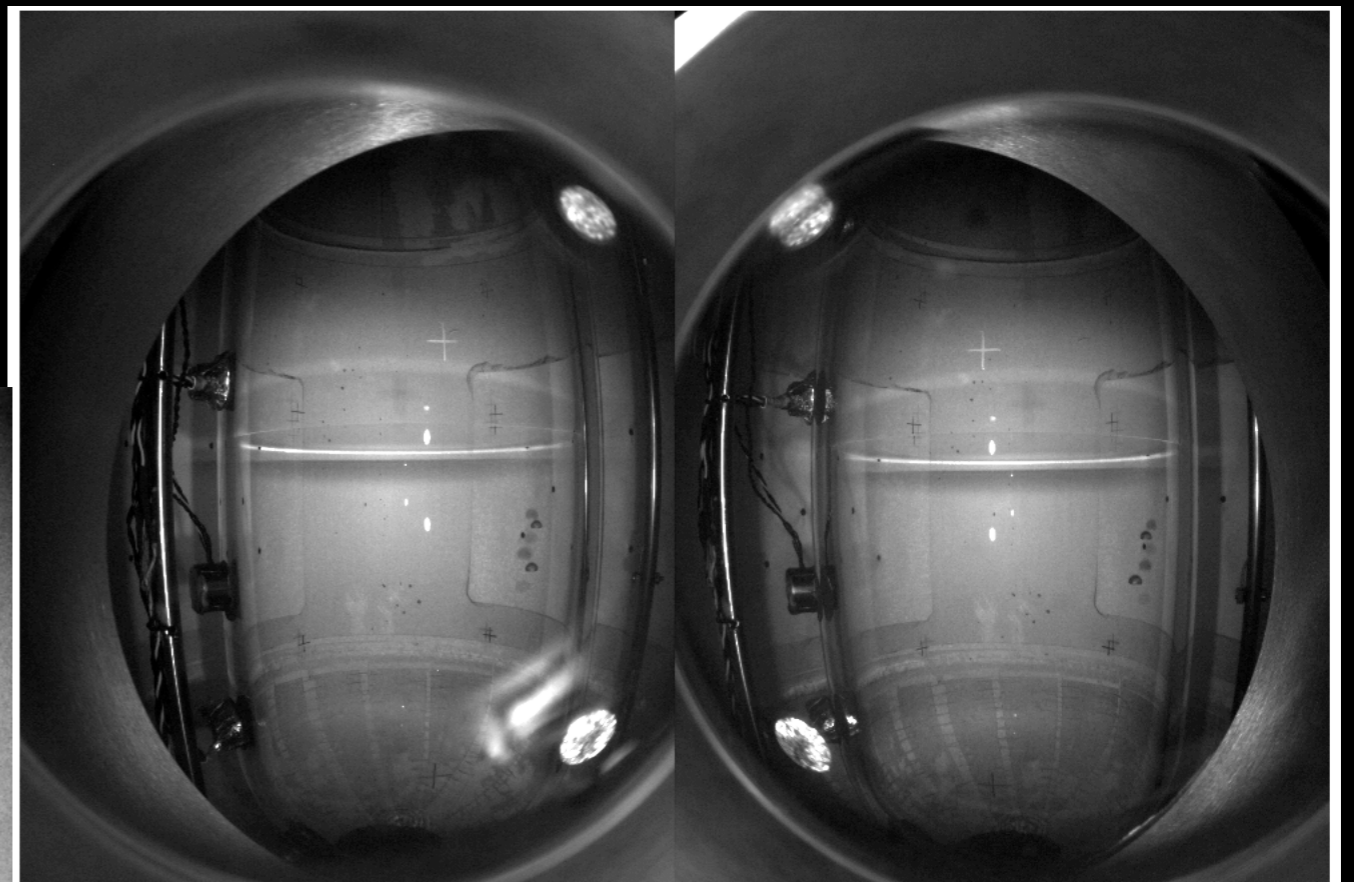
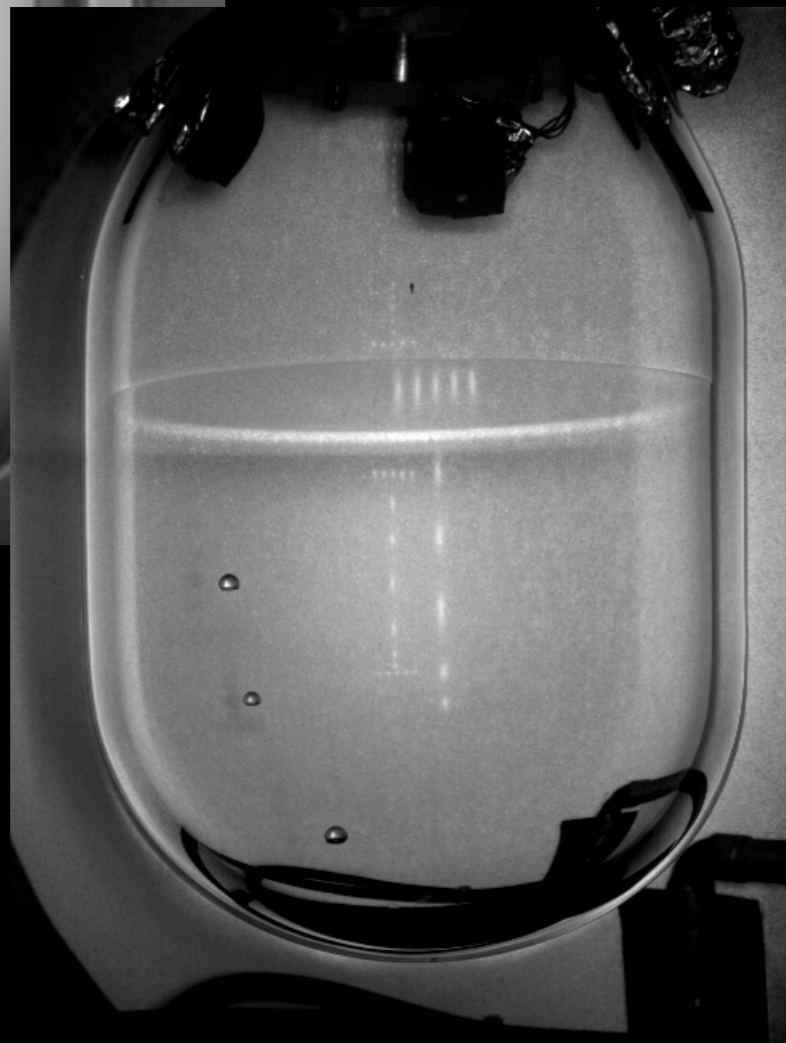
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Bubble Chambers!

Chicagoland Observatory for Underground Particle Physics (COUPP)

[Some debate over the pronunciation (should the Ps be silent?)]

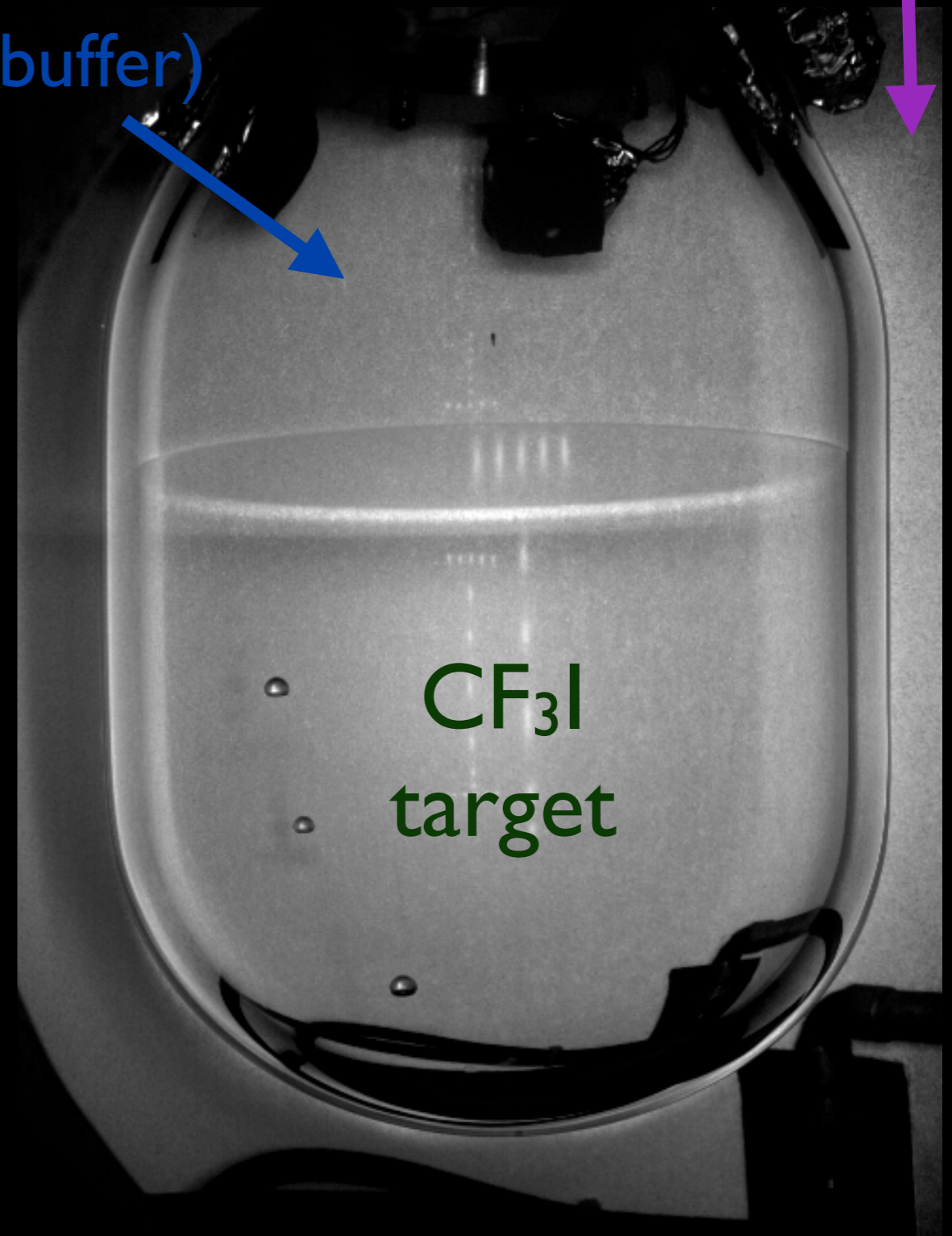


COUPP bubble chamber

- Pressure expansion creates superheated fluid, CF_3I
 - **F** for spin-dependent
 - **I** for spin-independent
 - Alternatives - e.g. C_3F_8
- Particle interactions nucleate bubbles
- Cameras see bubbles
- Recompress chamber to reset

Water
(buffer)

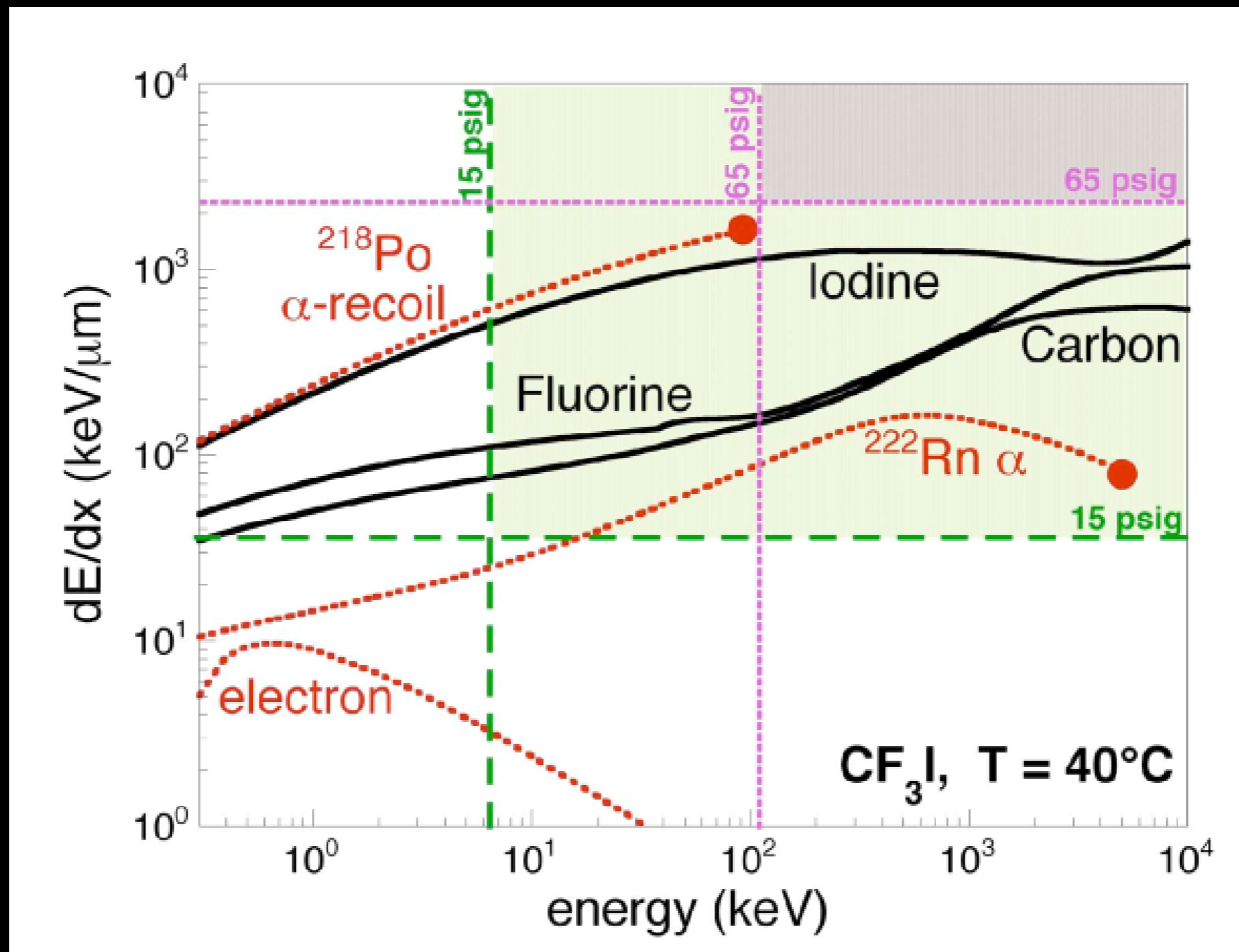
Propylene glycol
(hydraulic fluid)




Why bubble chambers?

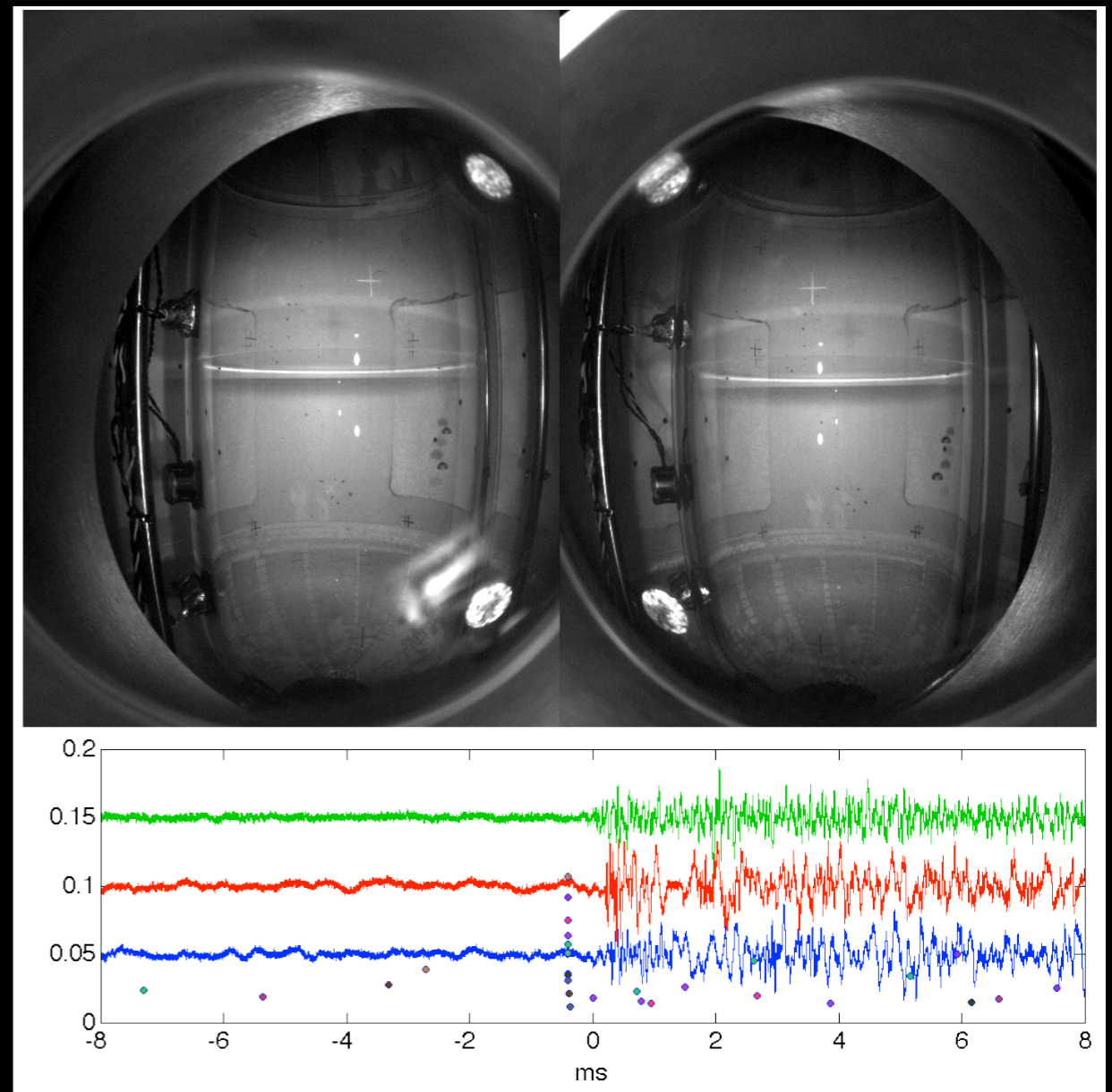
- To form a bubble requires two things
 - Enough energy
 - Enough energy density - length scale must be comparable to the critical bubble size
- By choosing superheat parameters appropriately (temperature and pressure), bubble chambers are blind to electronic recoils

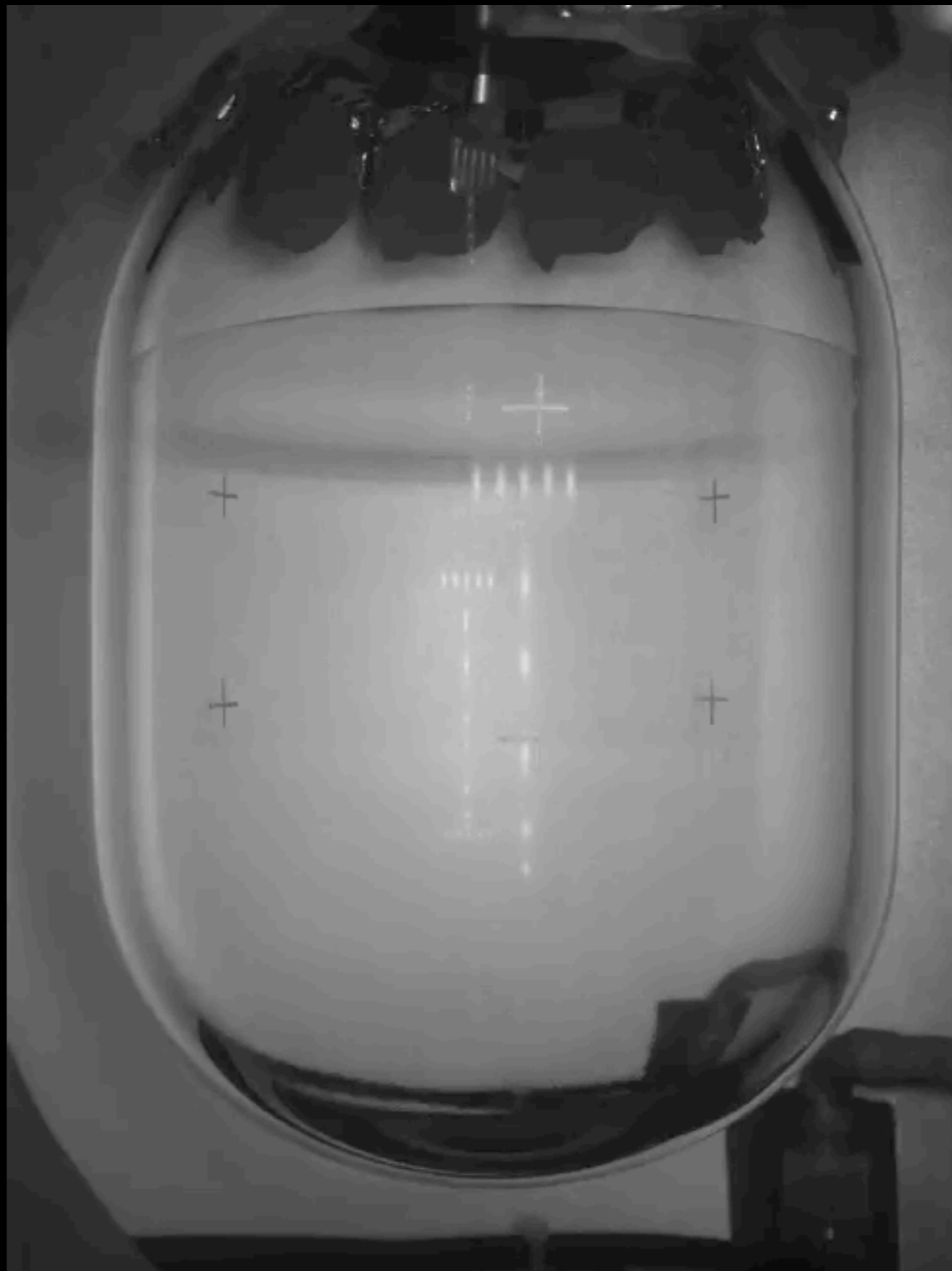
Why bubble chambers?



Why bubble chambers

- Easy to identify multiple scattering events  Neutron backgrounds
- Easy DAQ and analysis chain
 - Cameras
 - Piezos
- No PMTs, no cryogenics



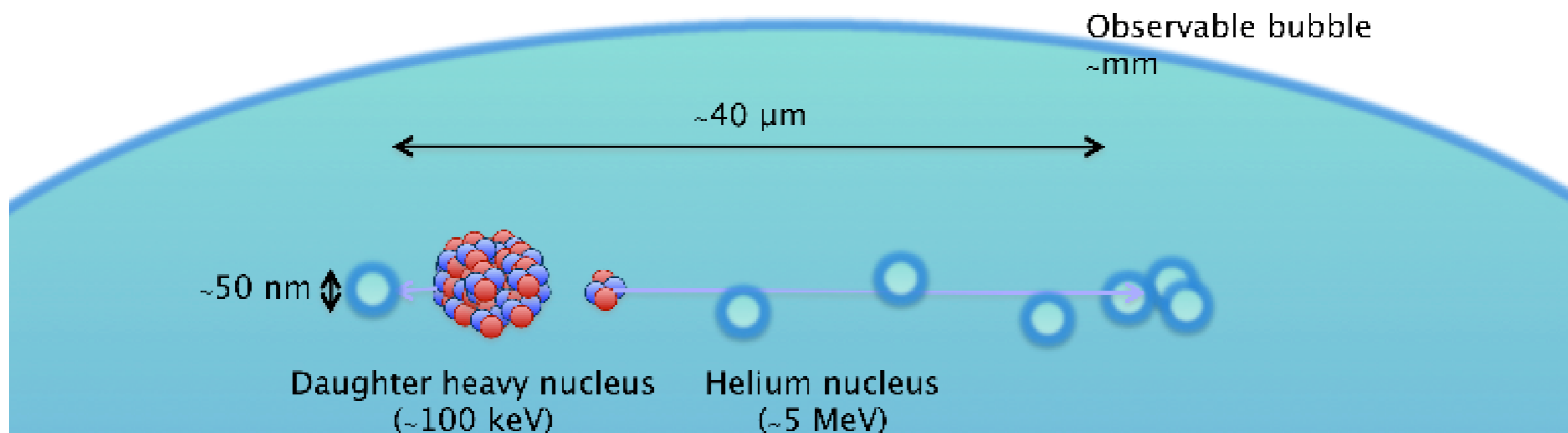


Why not bubble chambers?

- Threshold detectors - no energy resolution
 - Harder to distinguish some backgrounds, less information about any potential signal
 - Alphas (several MeV) were a big concern
 - Energy threshold calibrations are hard and important
- Bubble chambers are slow - about 30 s of deadtime for every event
 - Overall rate must be low

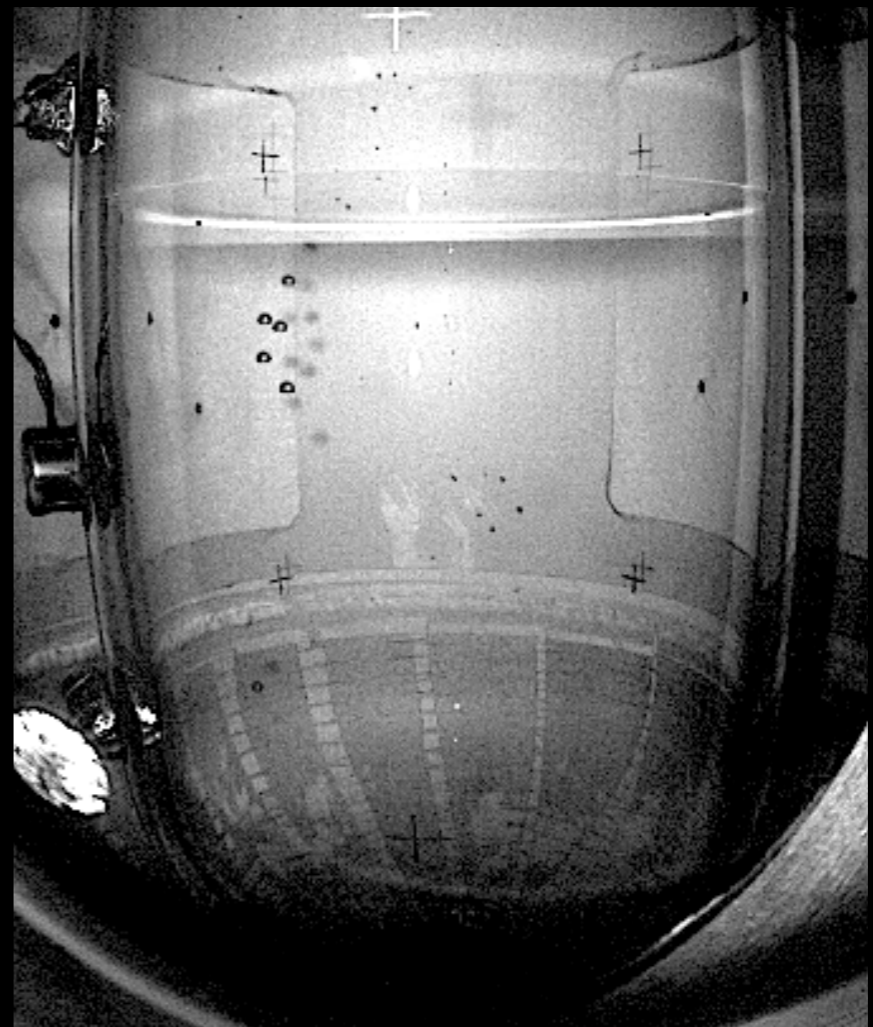
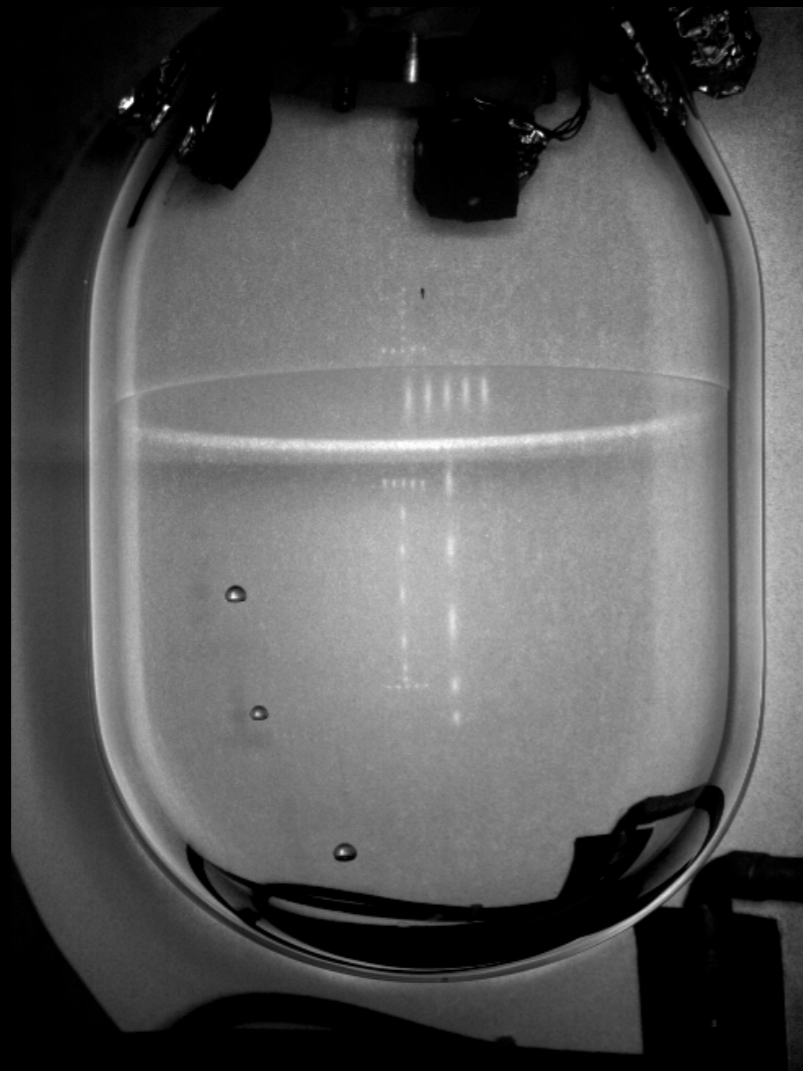
About those alphas

- Discovery of acoustic discrimination against alphas by PICASSO (Aubin et al, New J. Phys 10:103017, 2008)
- Alphas deposit energy over tens of microns
- Nuclear recoils deposit theirs in tens of nanometers
- In COUPP bubble chambers, alphas are several times louder

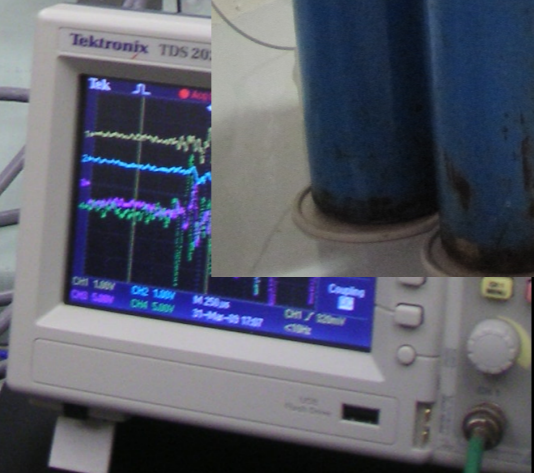


The COUPP program

- COUPP4: A 2-liter chamber operating at SNOLAB since 2010
- COUPP60: Up to 40 liters, commissioning at SNOLAB now
- COUPP500: Ton scale detector, funded by NSF and DOE, at SNOLAB in 2015?

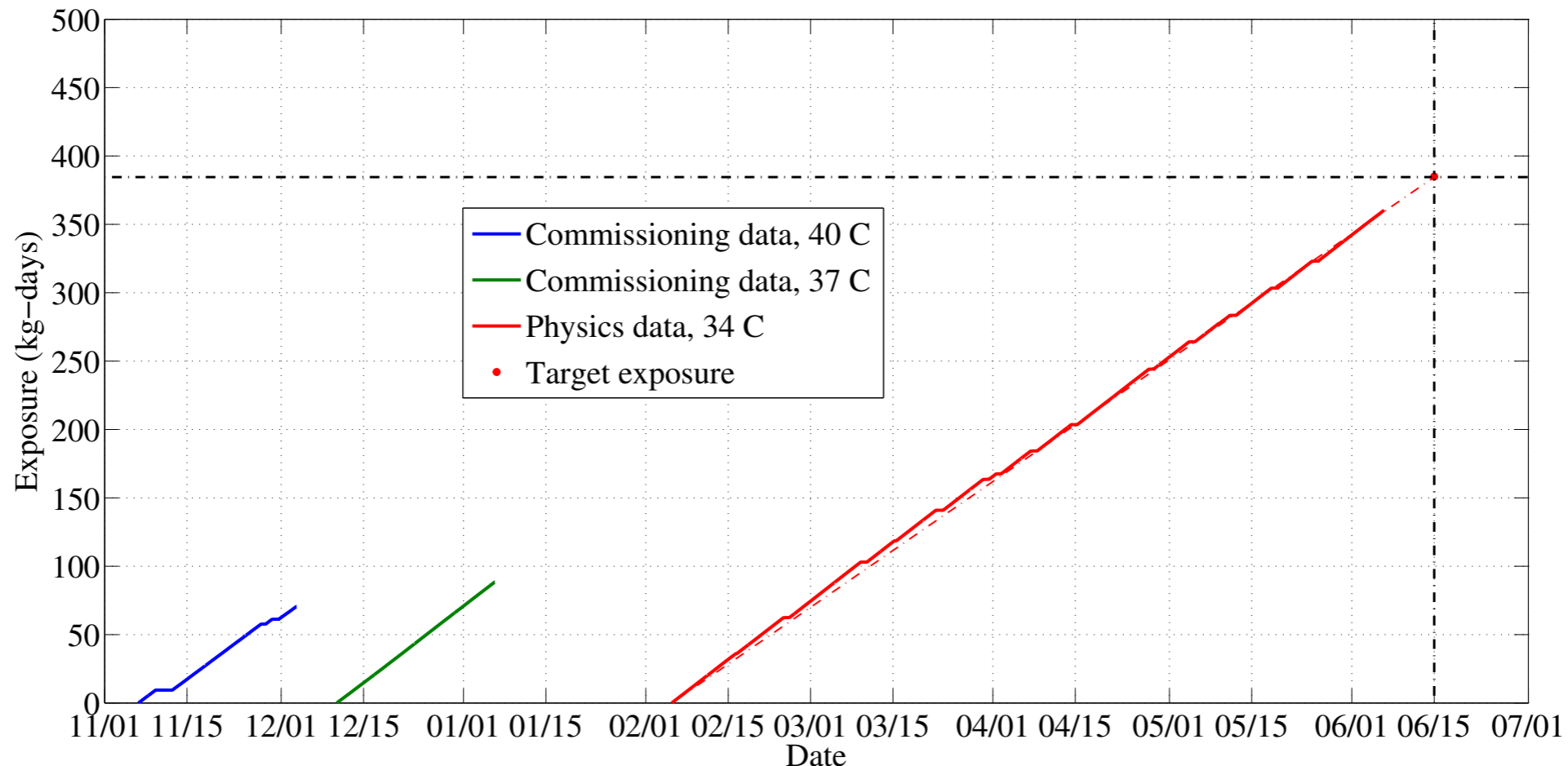


COUPP-4



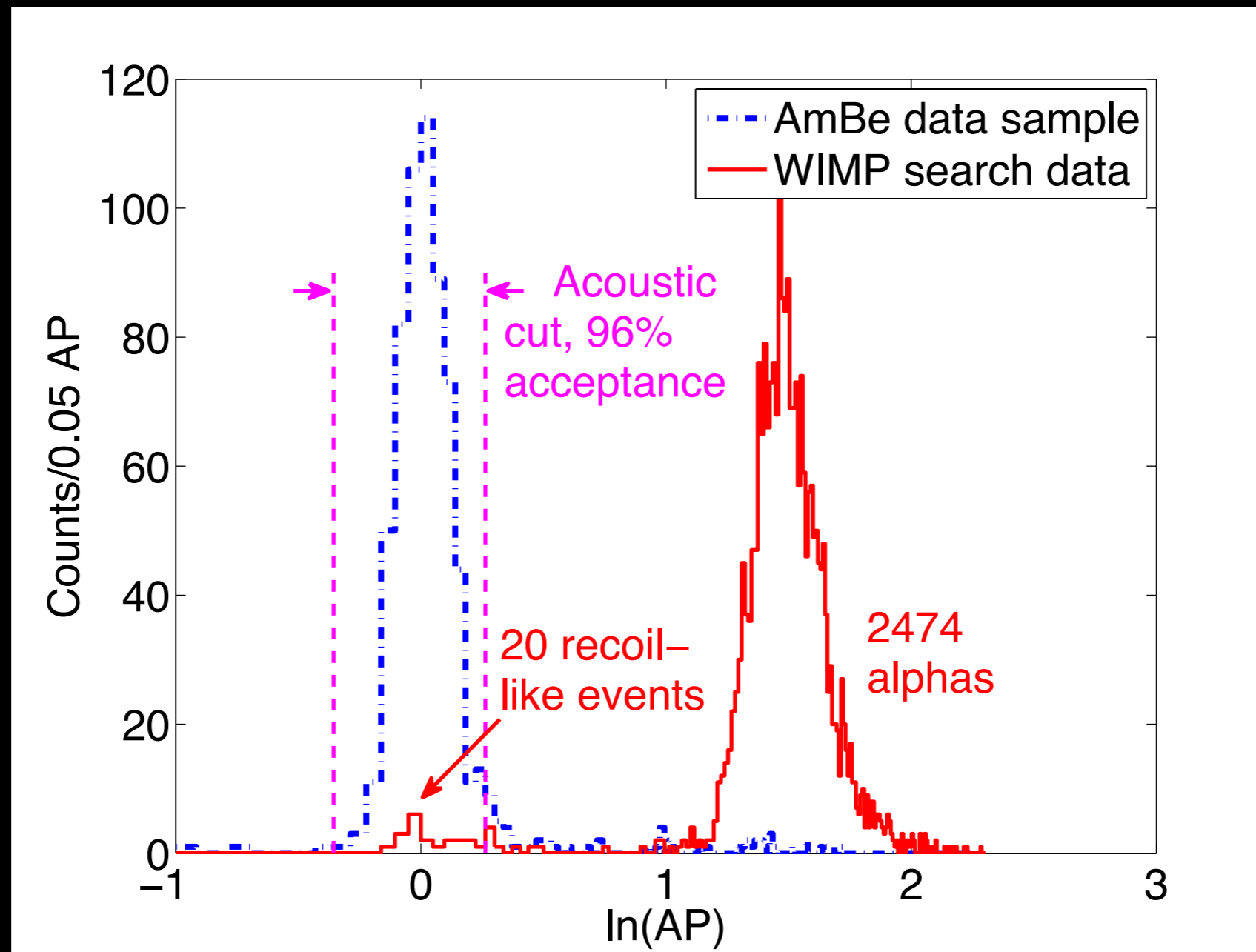
COUPP4: First run 2010-2011

- 17.4 live days at 8 keV threshold
- 21.9 live days at 11 keV threshold
- 97.3 live days at 16 keV threshold
- 79% acceptance for nuclear recoils after all cuts (including fiducial)

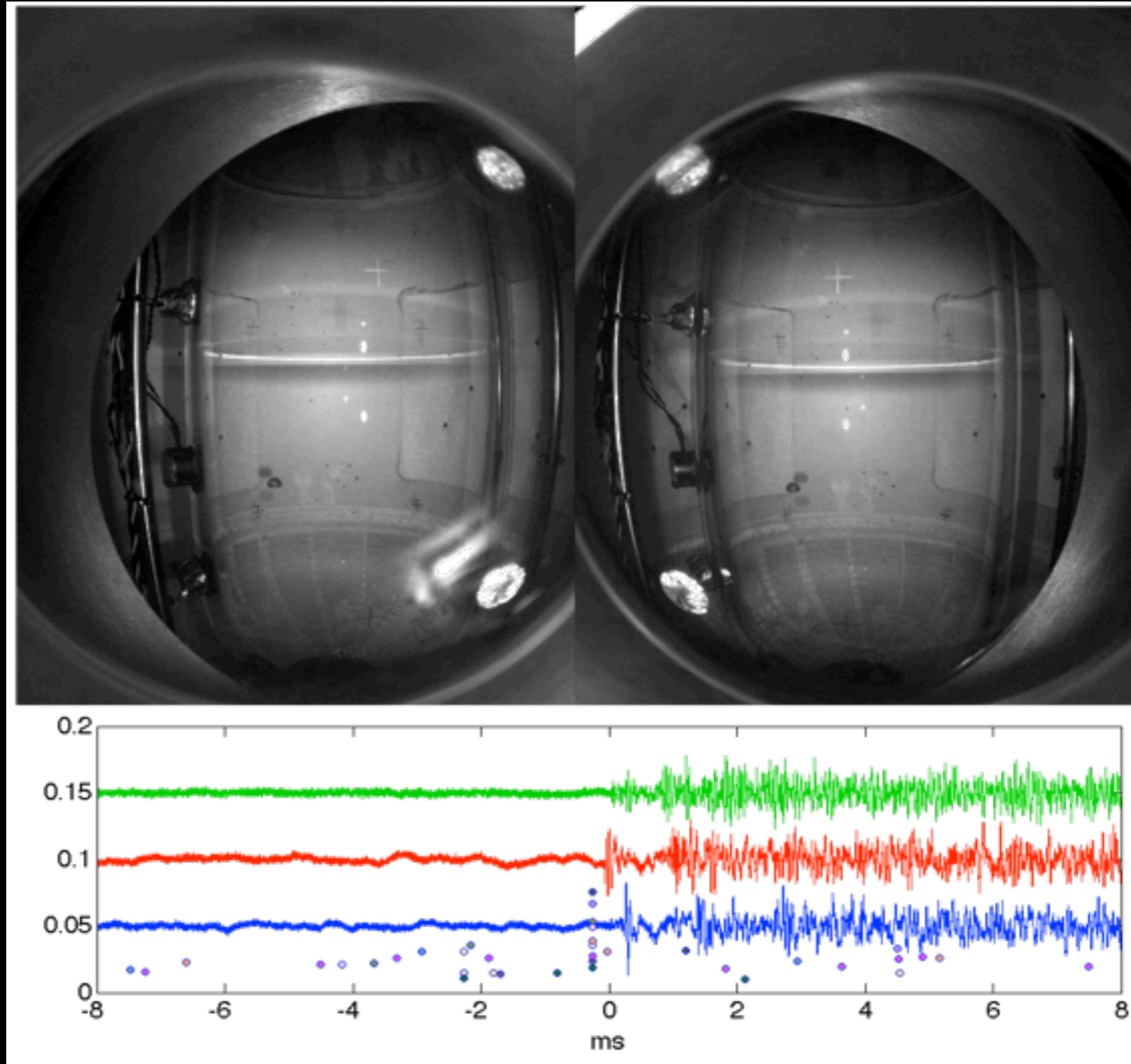


COUPP4: Acoustic discrimination

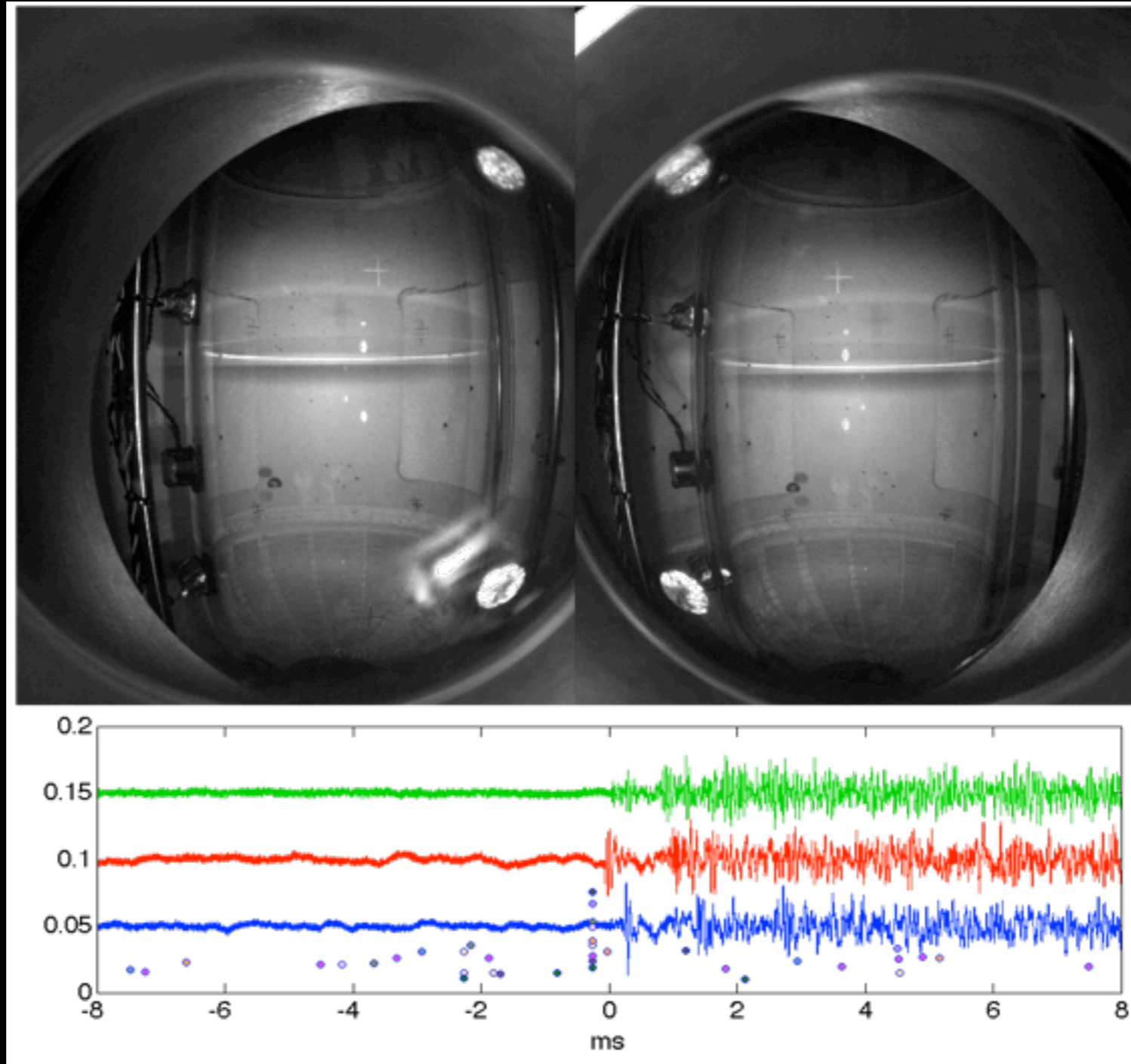
- Better than 99.3% rejection against alphas at 16 keV threshold
- Limited by statistics, and backgrounds



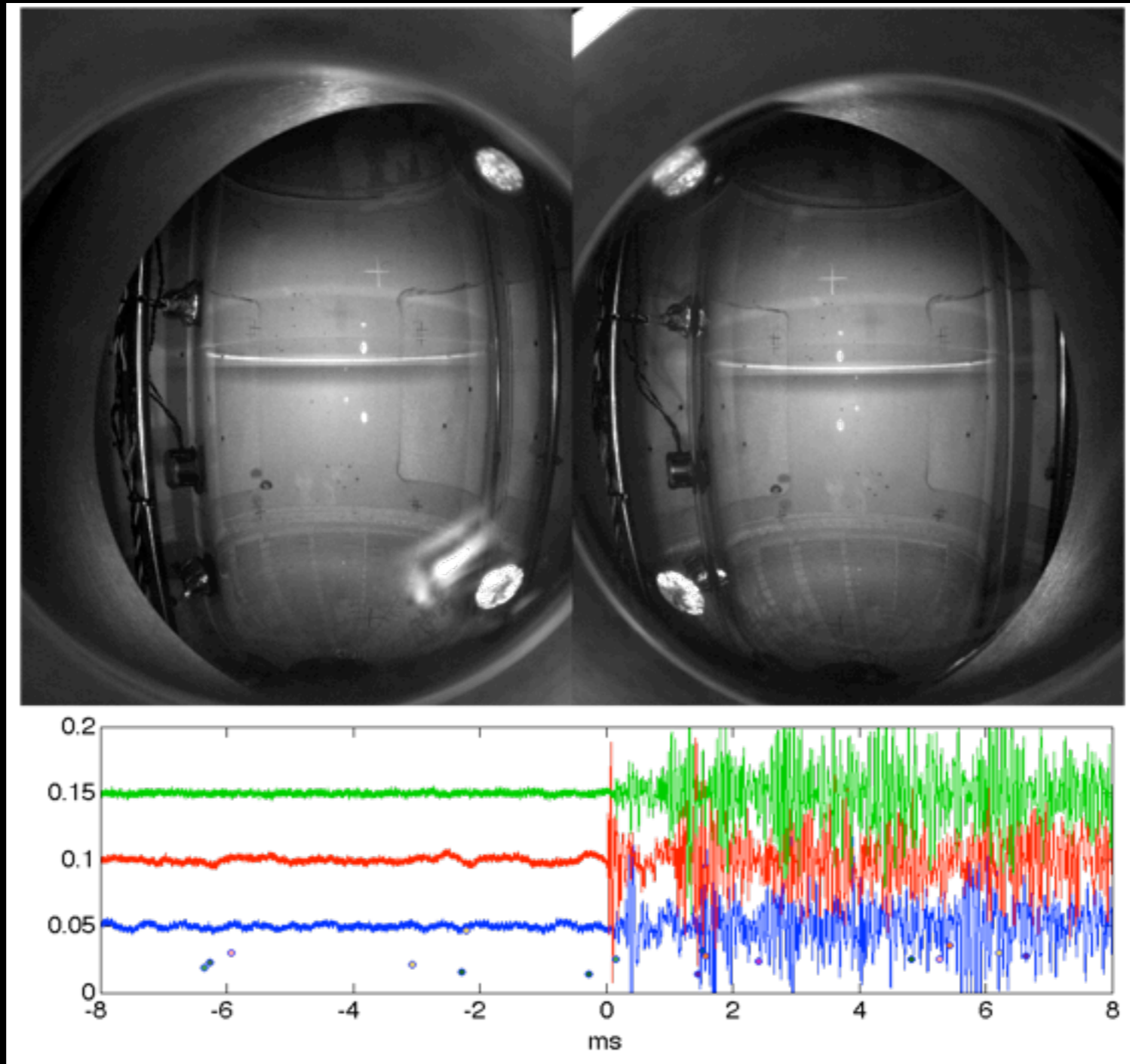
This is what dark matter would sound like



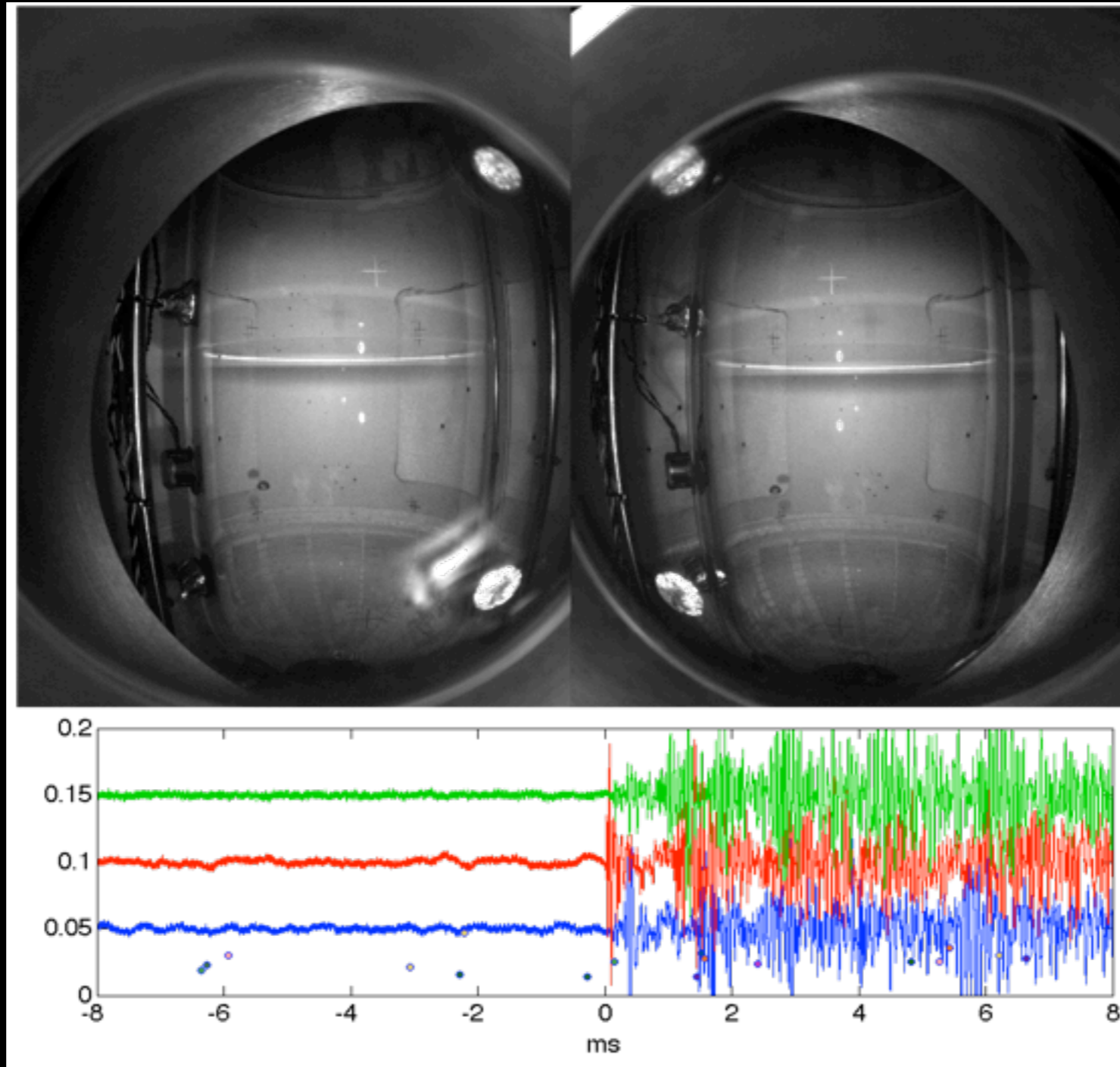
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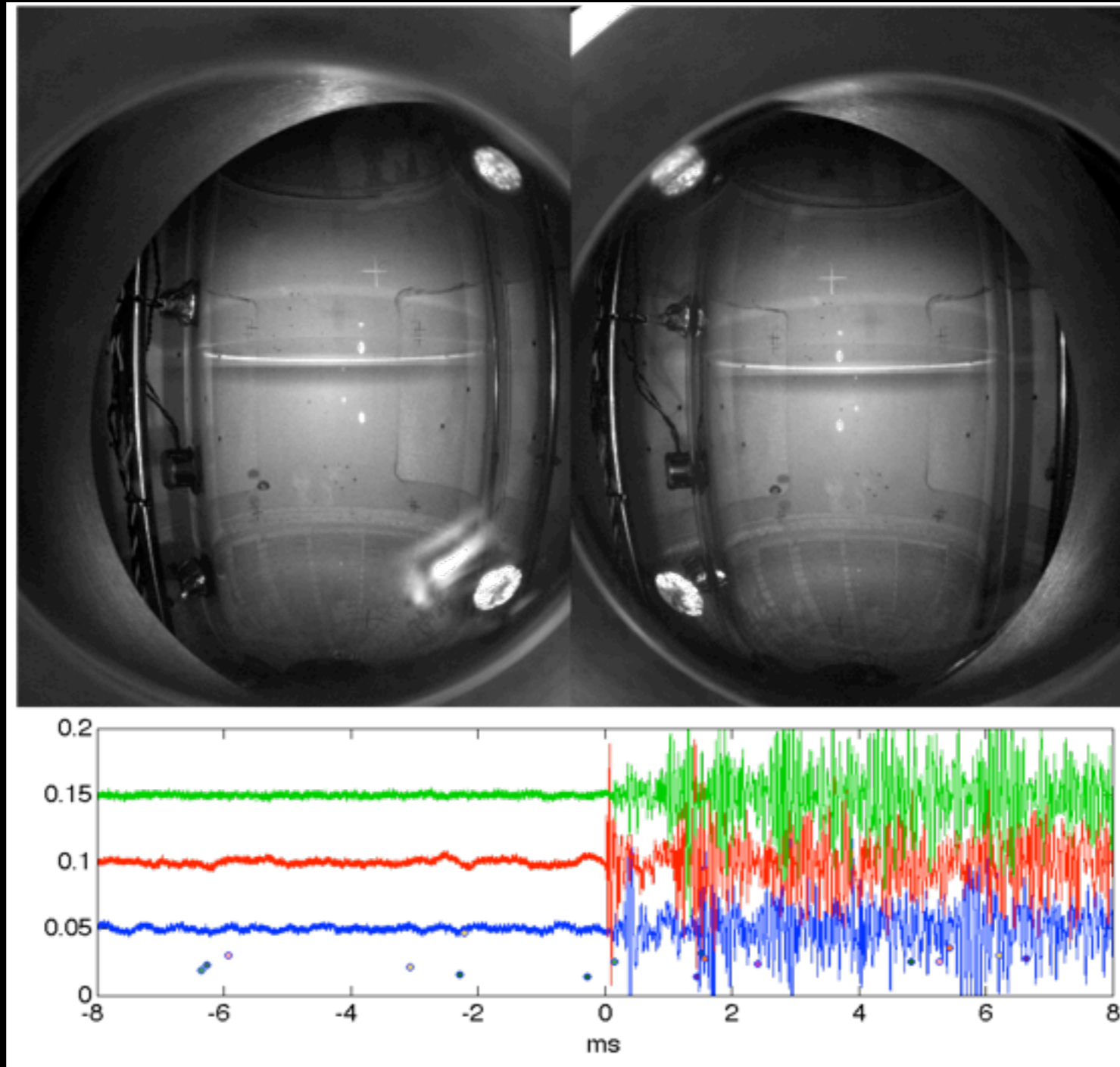
This is what an alpha sounds like



This is what an alpha sounds like



Both together, just to hear the difference



COUPP4: Results and sensitivity



- 20 WIMP candidates (8 at 8 keV, 6 at 11 keV, 8 at 16 keV)
- 3 multiple bubble events imply neutrons

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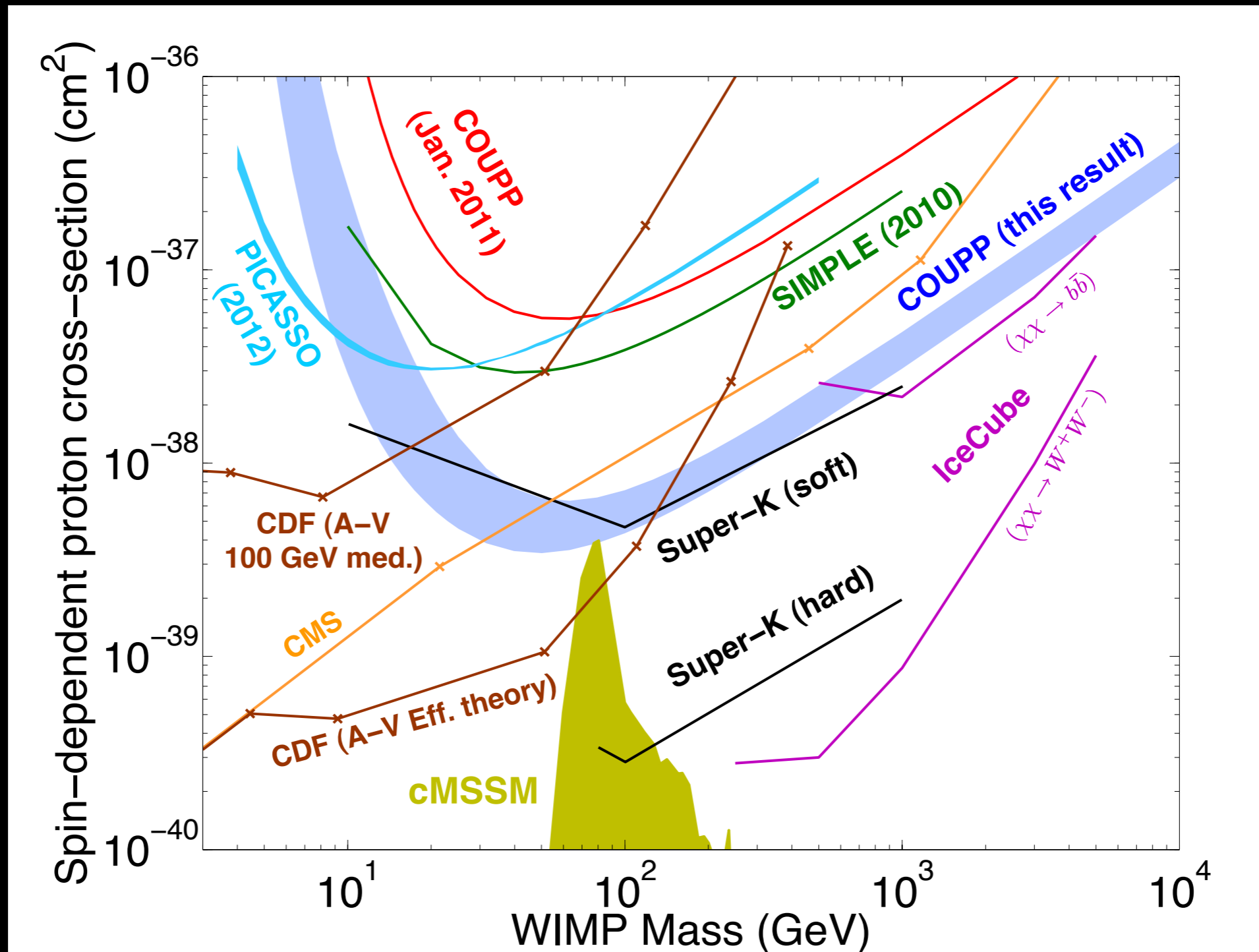
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- 20 WIMP candidates (8 at 8 keV, 6 at 11 keV, 8 at 16 keV)
- 3 multiple bubble events imply neutrons
- U,Th in the piezo-acoustic sensors and the viewports
- Remaining excess of singles at low threshold
- Time clustering
- Correlated with activity at water-CF₃I interface

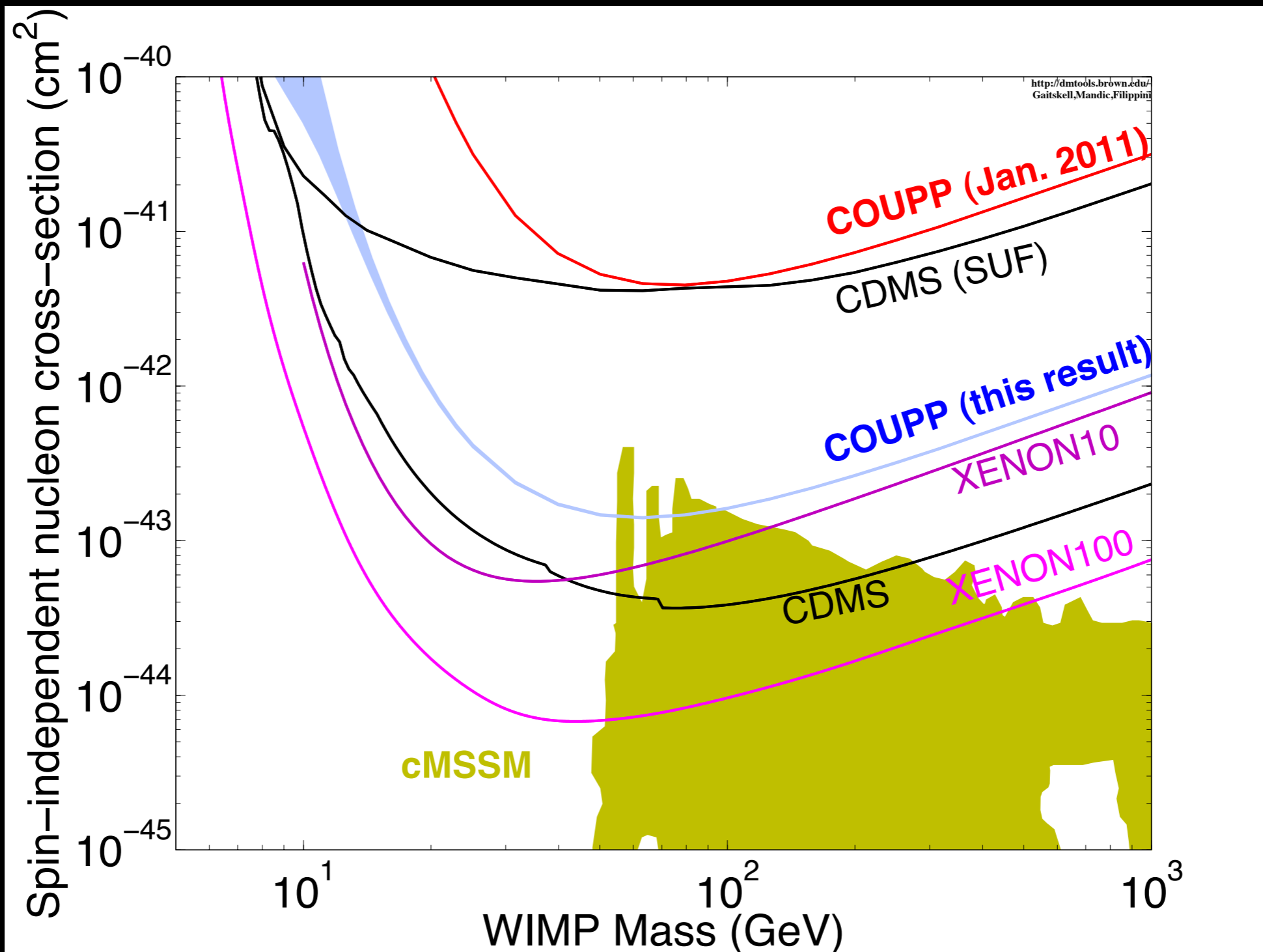
COUPP4: Results and sensitivity

- Given uncertainties on backgrounds, no background subtraction: PRD 86:052001 (2012)



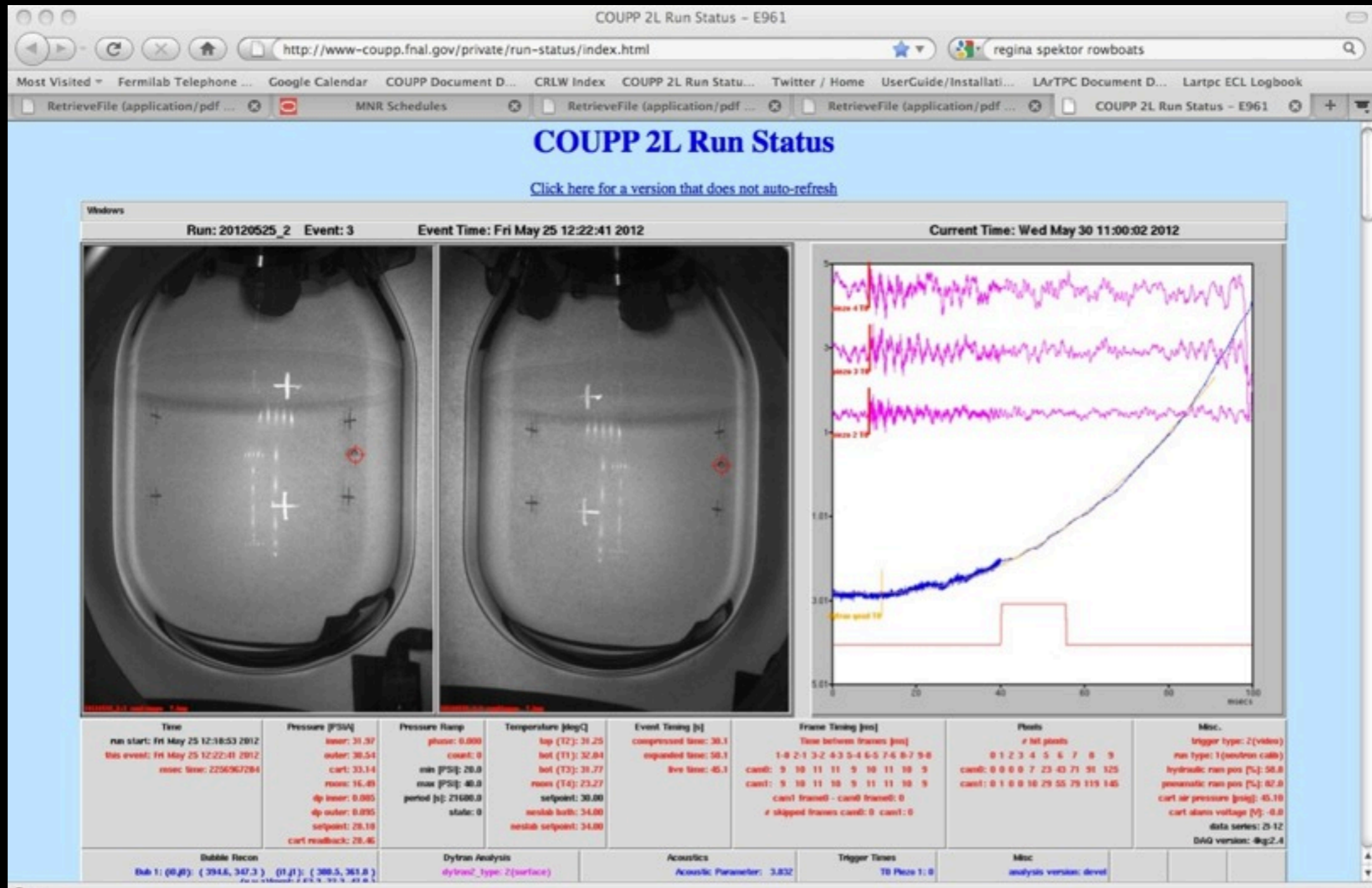
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COUPP4: Results and sensitivity

- Removed known neutron sources and improved fluid purification
- Second run ended last November



Detour: Threshold and efficiency

- Threshold determined from Seitz, Phys. of Fluids **1**, 2 (1958)

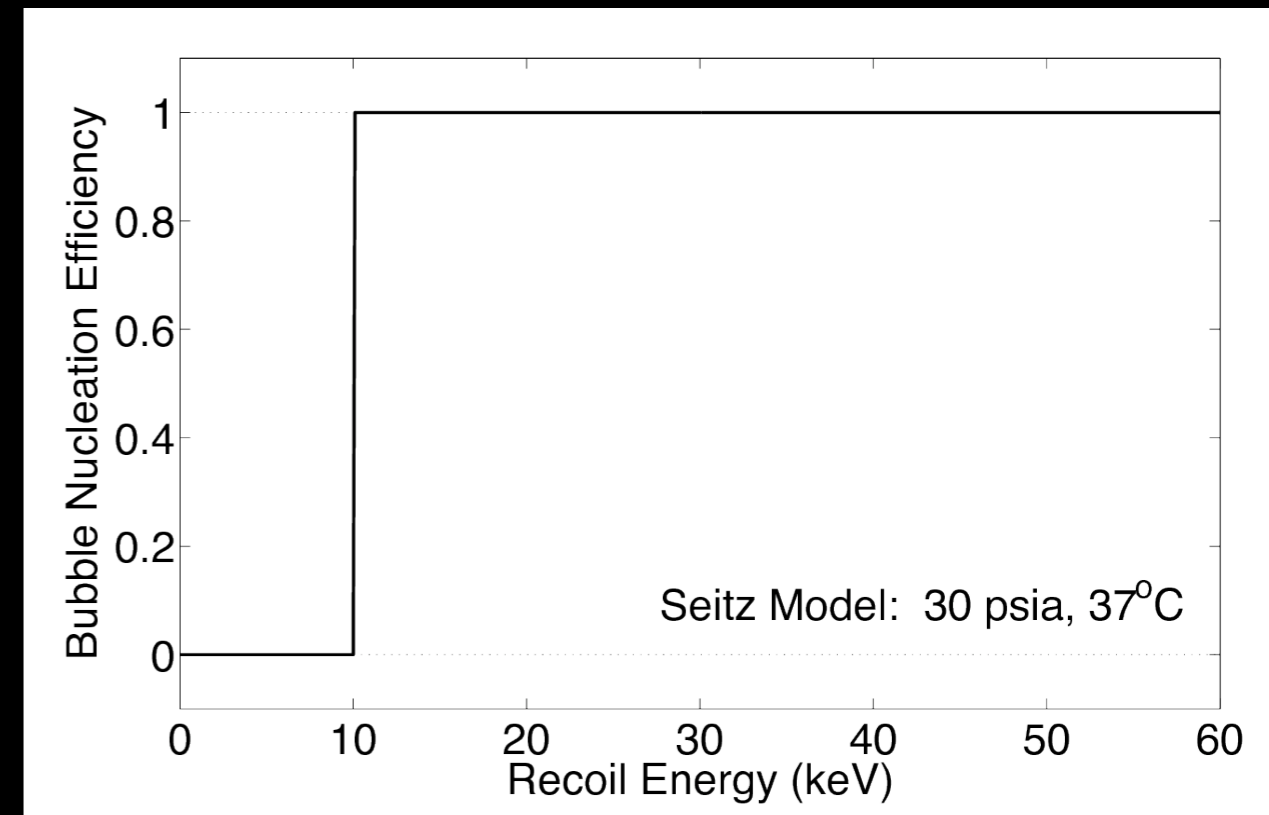
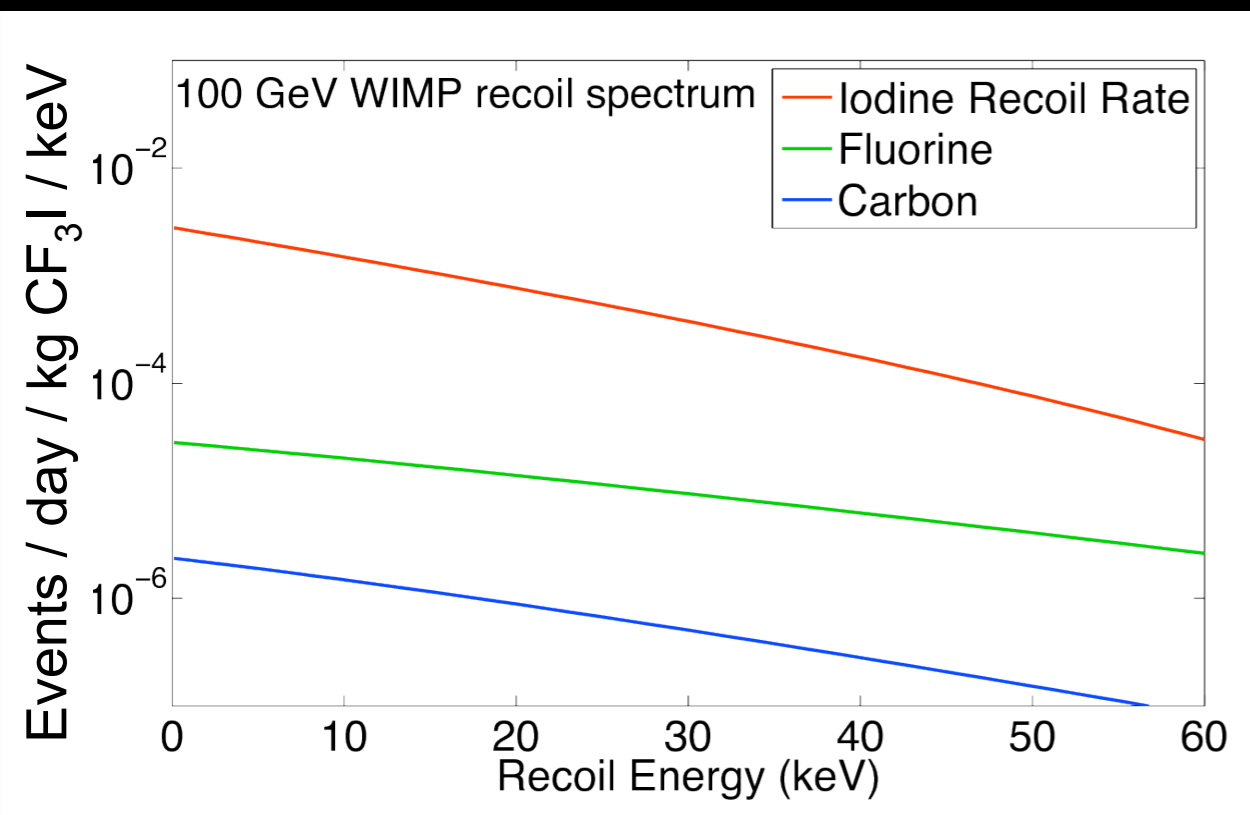
$$p_v - p_l = \frac{2\sigma}{r_c}$$
$$E_{th} = 4\pi r_c^2 \left(\sigma - T \frac{\partial \sigma}{\partial T} \right) + \frac{4}{3} \pi r_c^3 \rho_v h$$

Surface energy Latent heat

- Energy deposition E_{th} within length R_c will nucleate a bubble (Hot Spike model)
- Theory assumes a step function above threshold

Detour: Threshold and efficiency

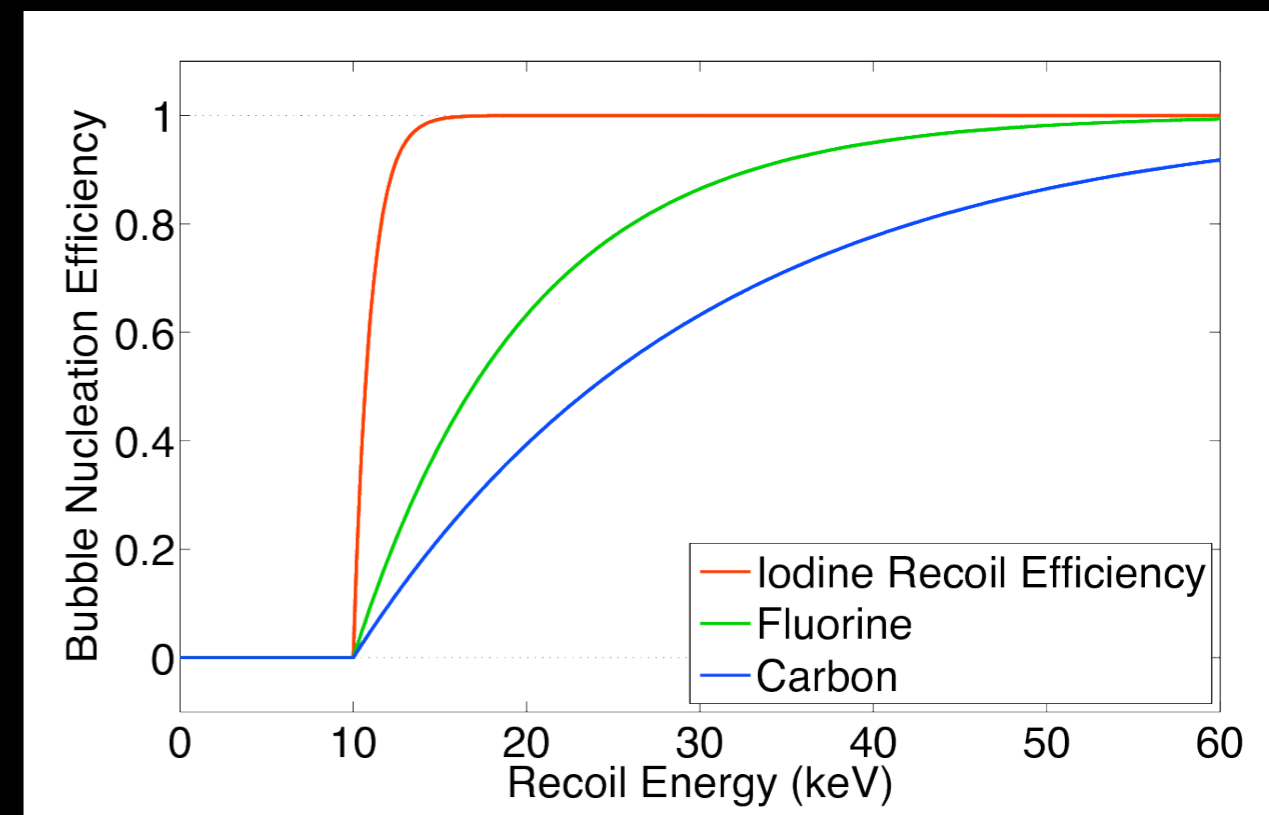
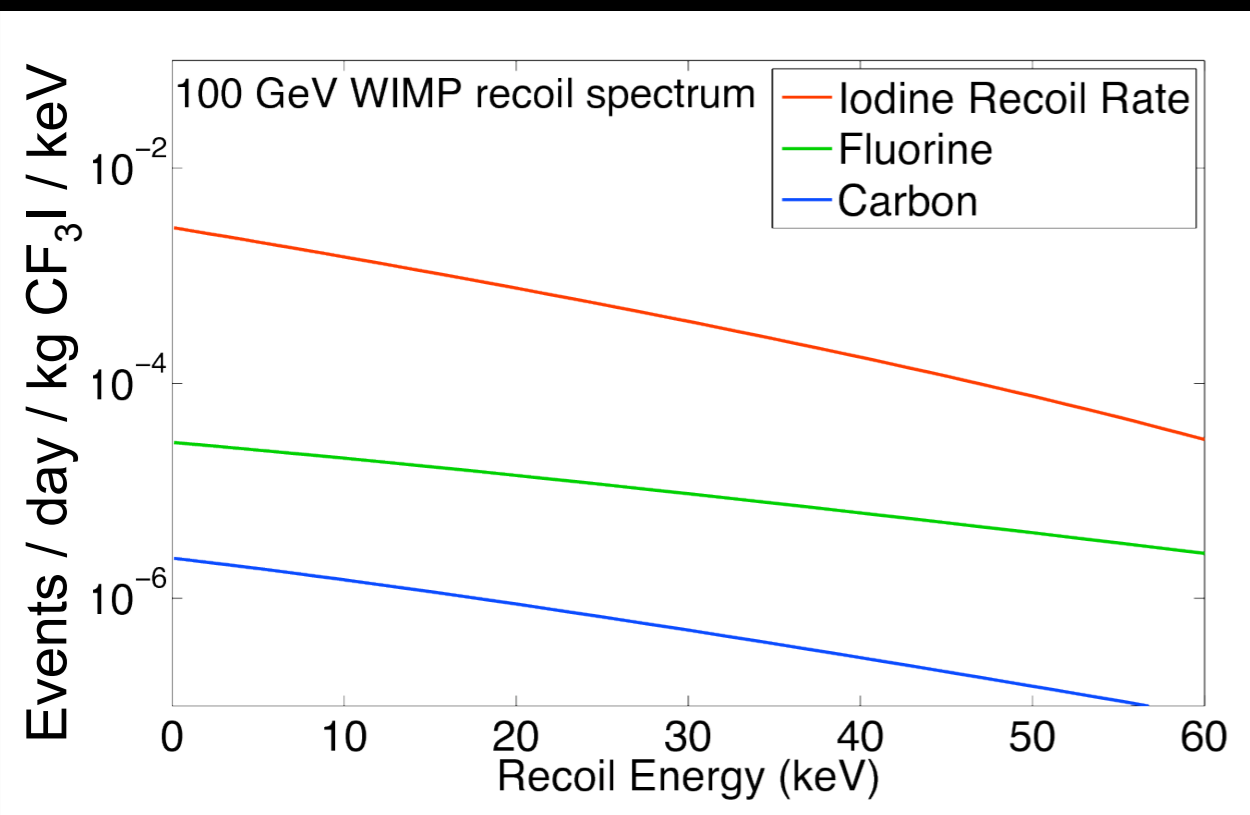
$$\text{Rate} = \int \text{WIMP recoil spectrum} \times \text{Bubble nucleation efficiency}$$



- Effect of threshold shape depends on target, WIMP mass

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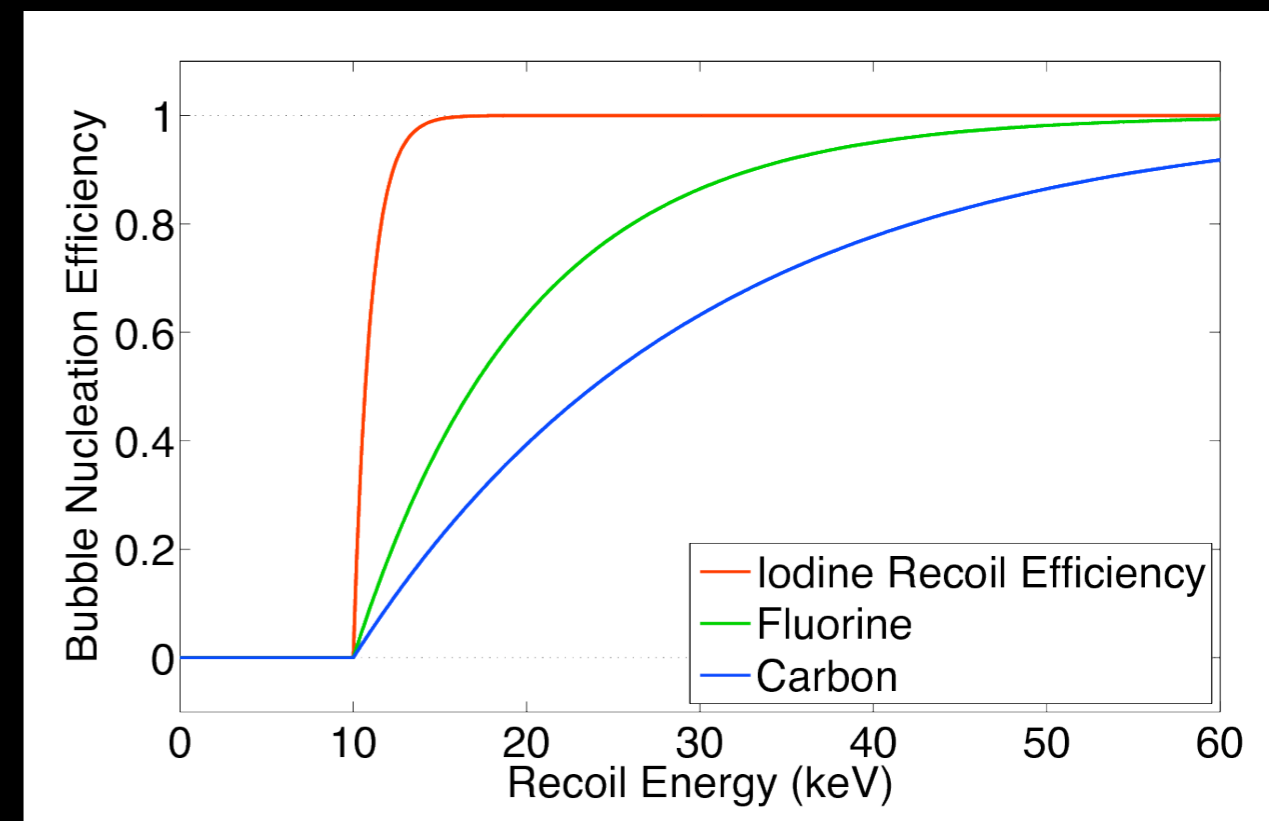
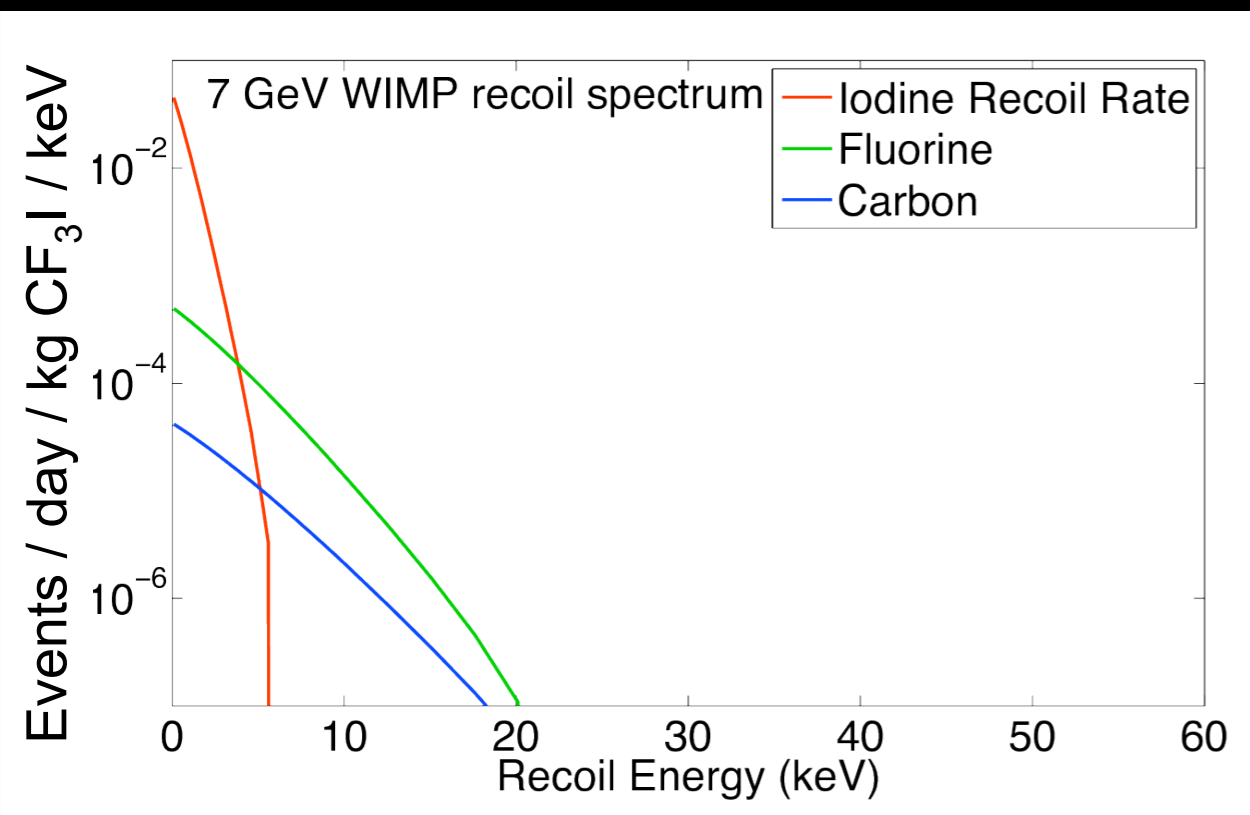
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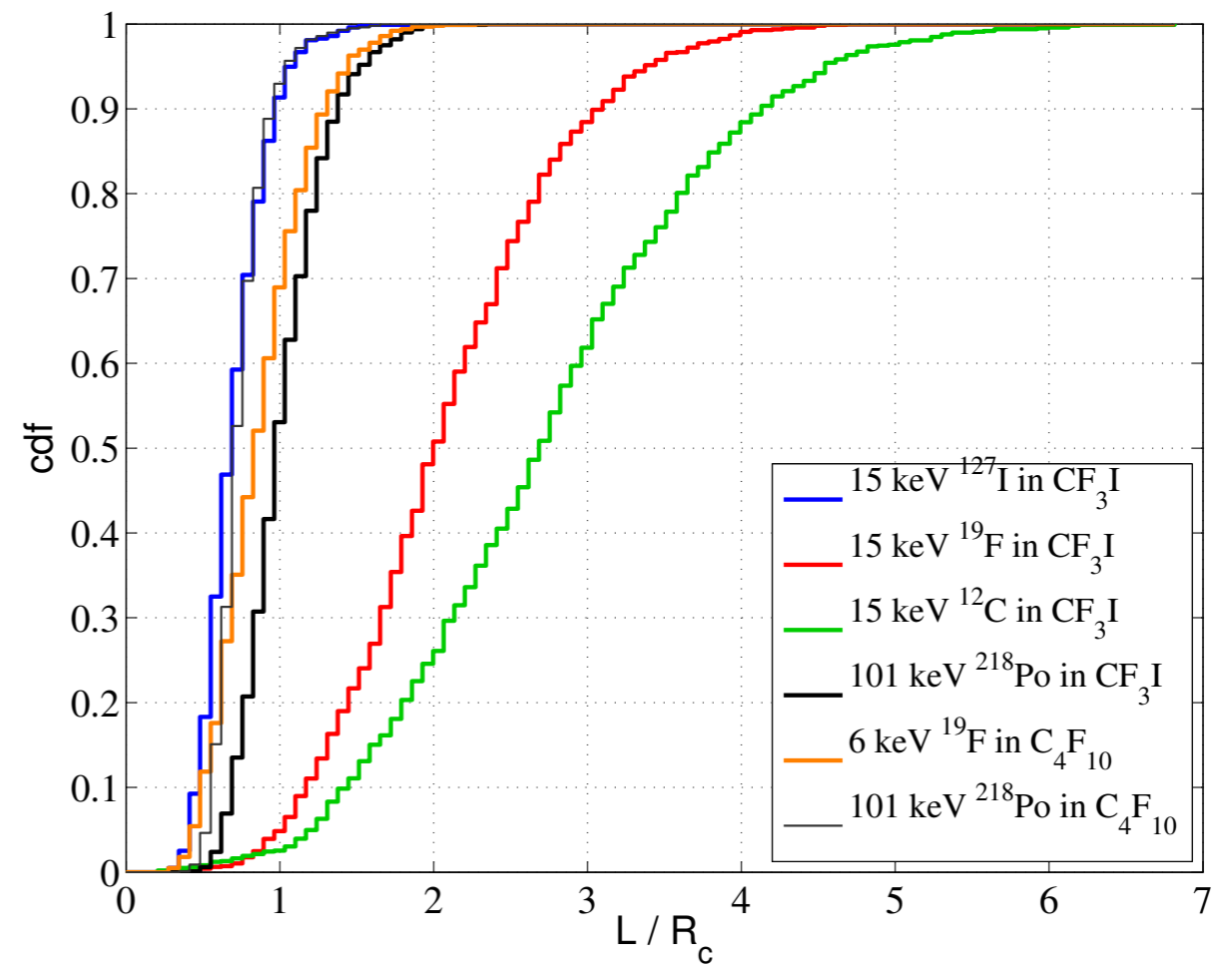
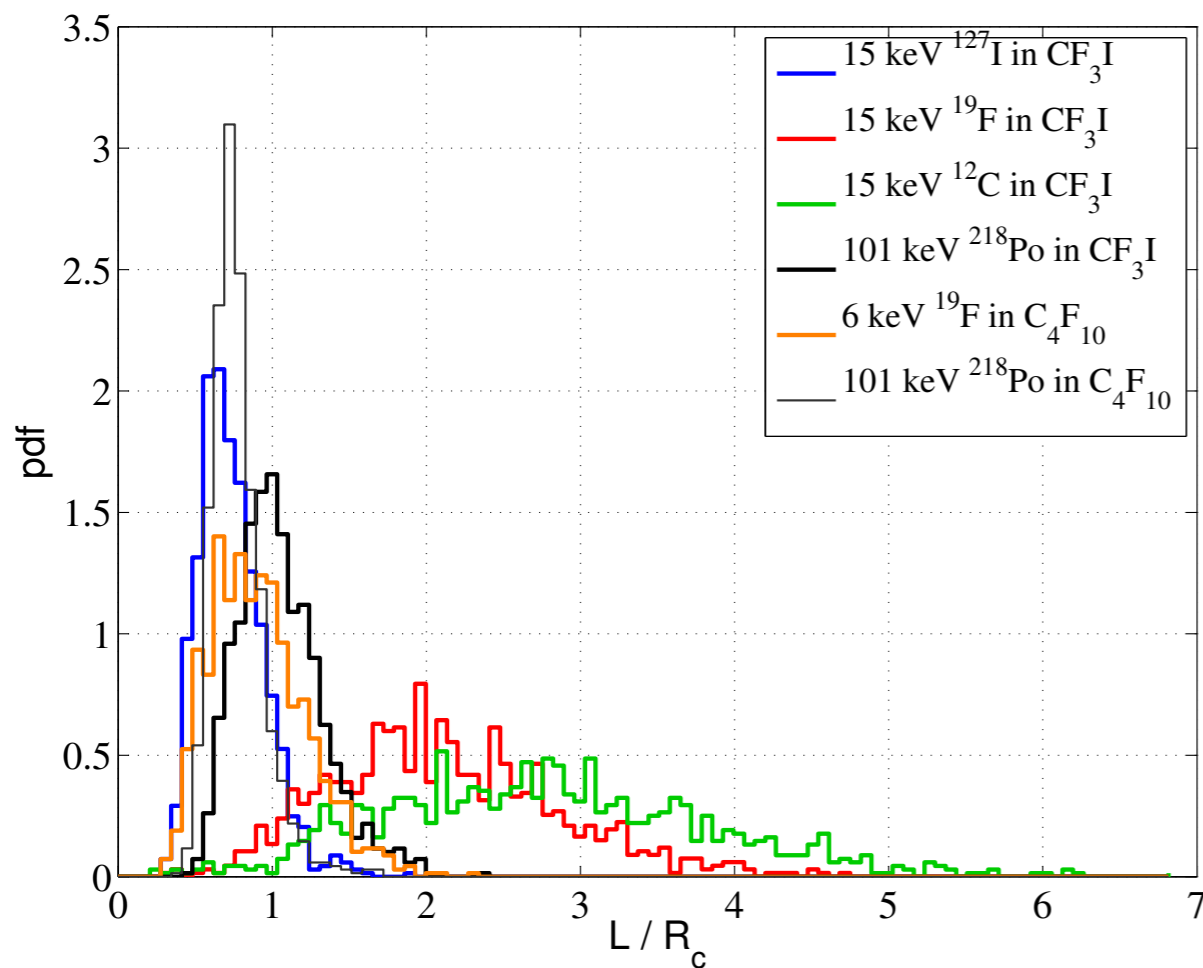
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$$E_{th} = 4\pi r_c^2 \left(\sigma - T \frac{\partial \sigma}{\partial T} \right) + \frac{4}{3} \pi r_c^3 \rho_v h$$

Surface energy Latent heat

- Energy deposition E_{th} within length r_c will nucleate a bubble (Hot Spike model)
- Theory assumes a step function above threshold
- Needs calibration

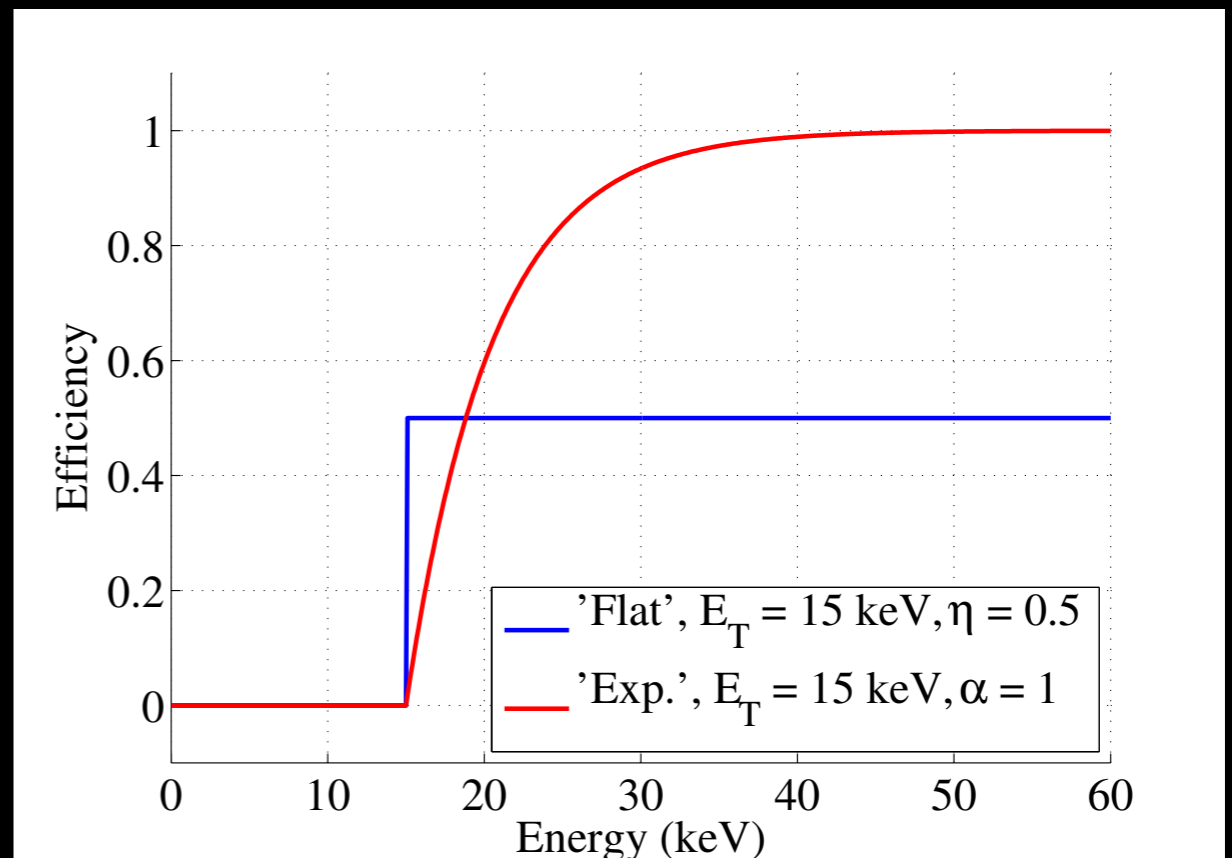
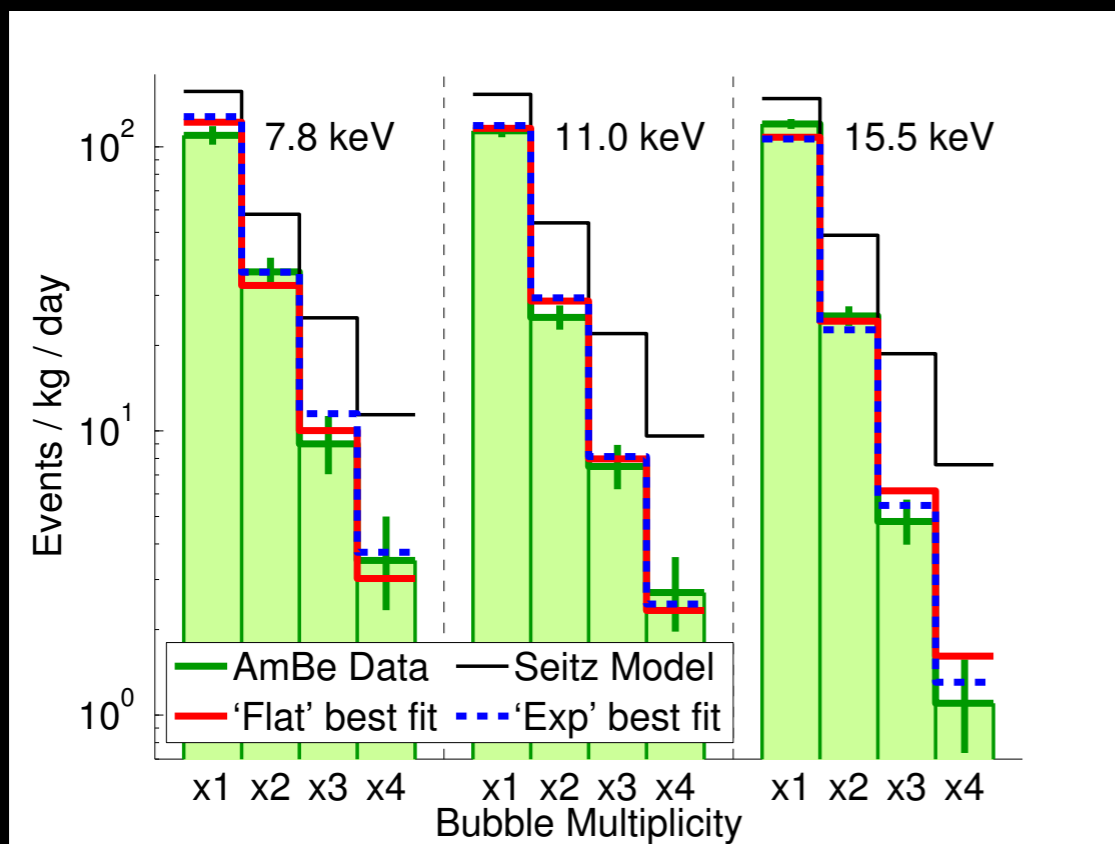
Detour: Threshold and efficiency

- Complicated by molecule, CF_3I
- Recall that the recoil track length L must be comparable to the bubble radius R_c



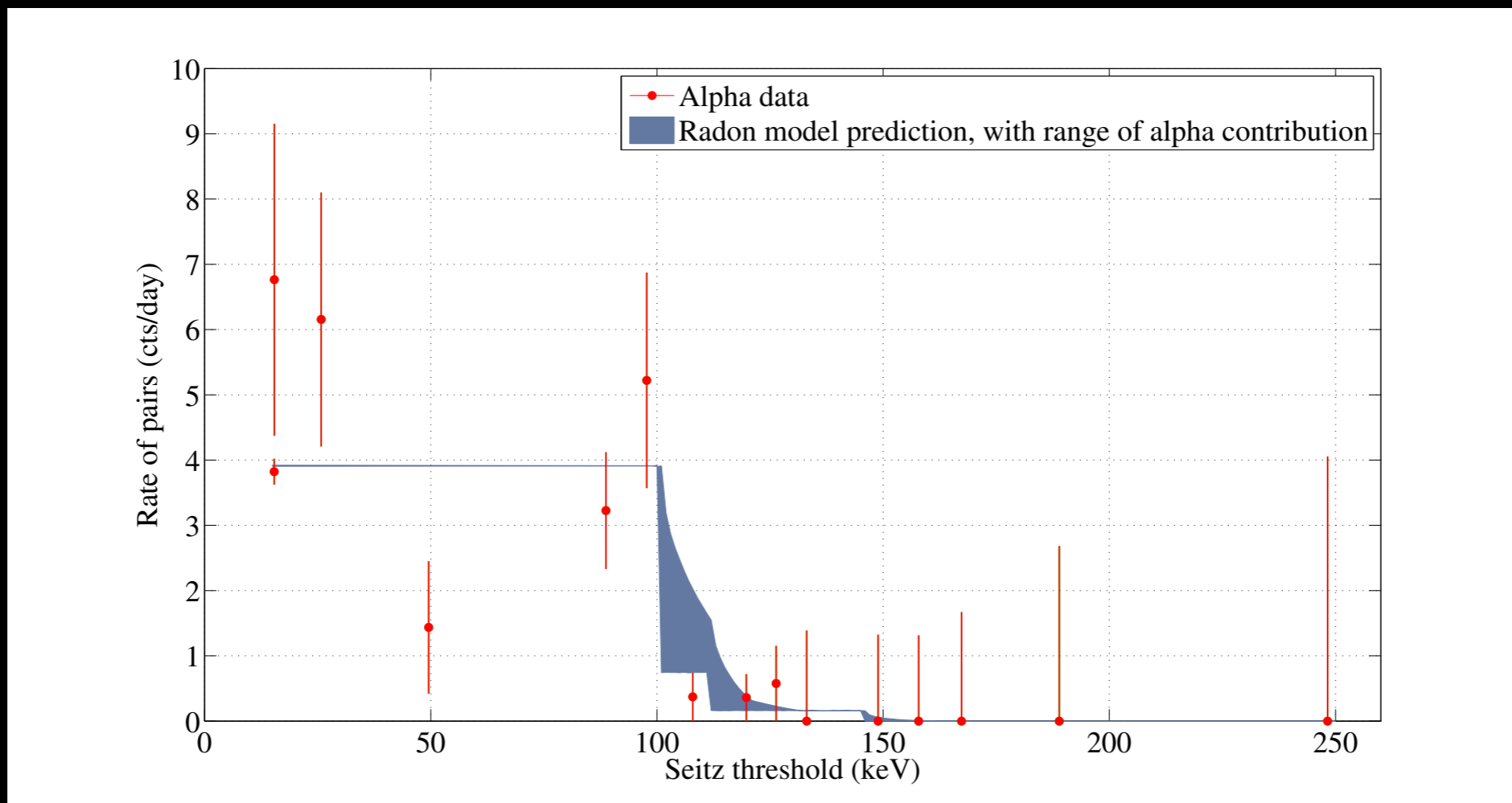
Carbon and fluorine

- Use neutron calibration sources at SNOLAB
- Compare MCNP-predicted rates of single, double, triple and quadruple bubble events with observation
- Data show a shortfall of events compared to simulation of the Seitz Model- i.e. the threshold is not a step function



What about iodine?

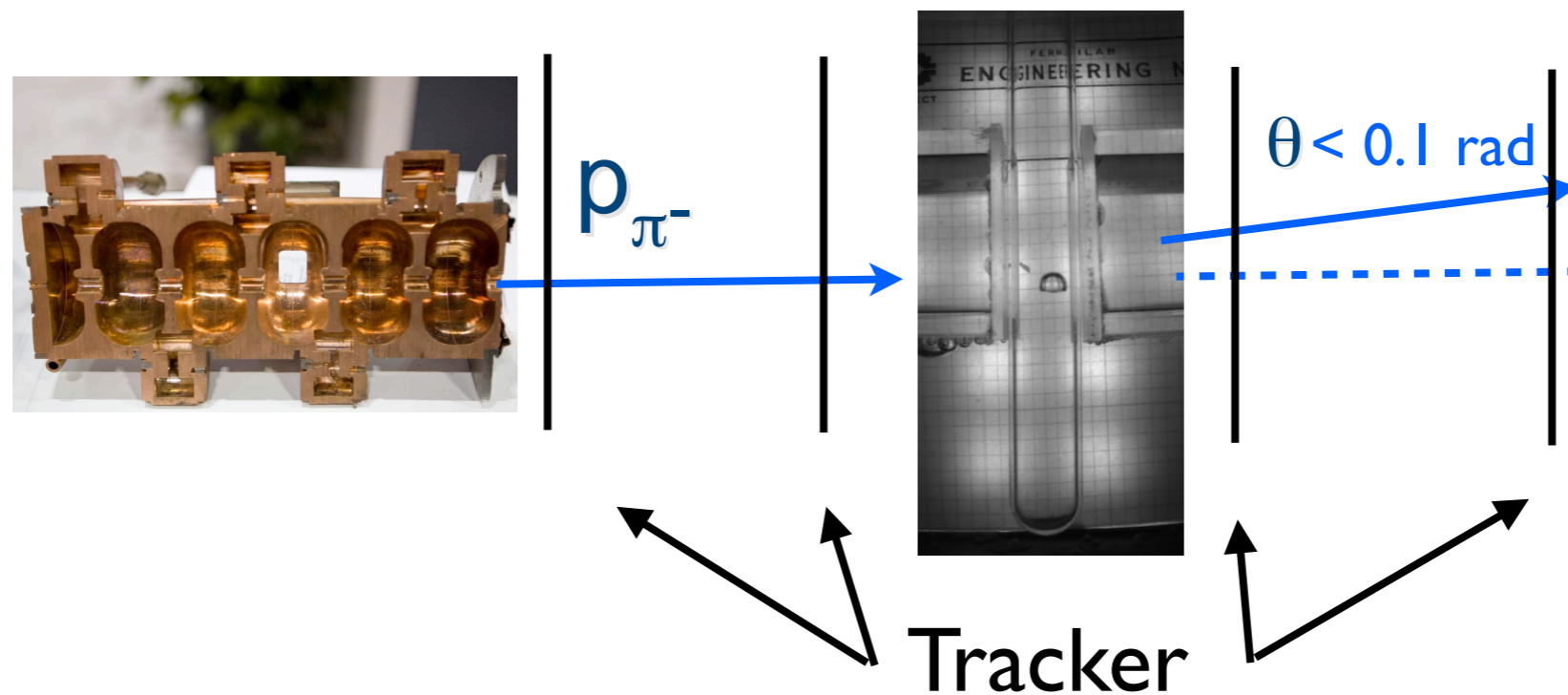
- Main sensitivity to spin independent dark matter from iodine
 - 85% of neutron source interactions are with C and F
- Heavy radon daughter nuclei are a proxy and are step-like



- We really need a direct calibration

COUPP Iodine Recoil Threshold Experiment

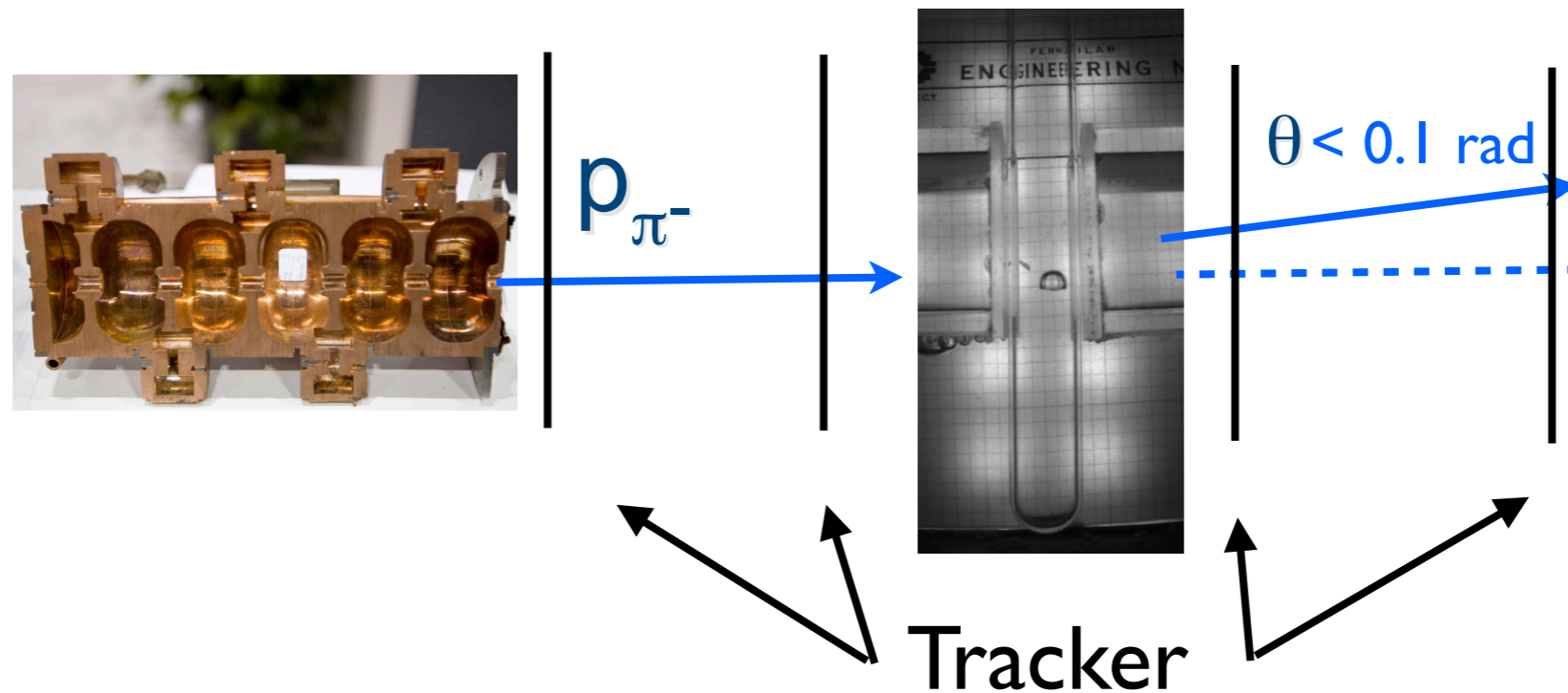
- Bubble chambers are insensitive to MIPs
- Elastic scattering of charged particles can be tracked with very high precision



$$T = E_{recoil} = \frac{(p\theta)^2}{2m_r}$$

COUPP Iodine Recoil Threshold Experiment

- Provides event by event energy information bubble chambers normally can't provide
- 75% of elastic scattering events with 12 GeV pions at energies relevant to dark matter involve iodine



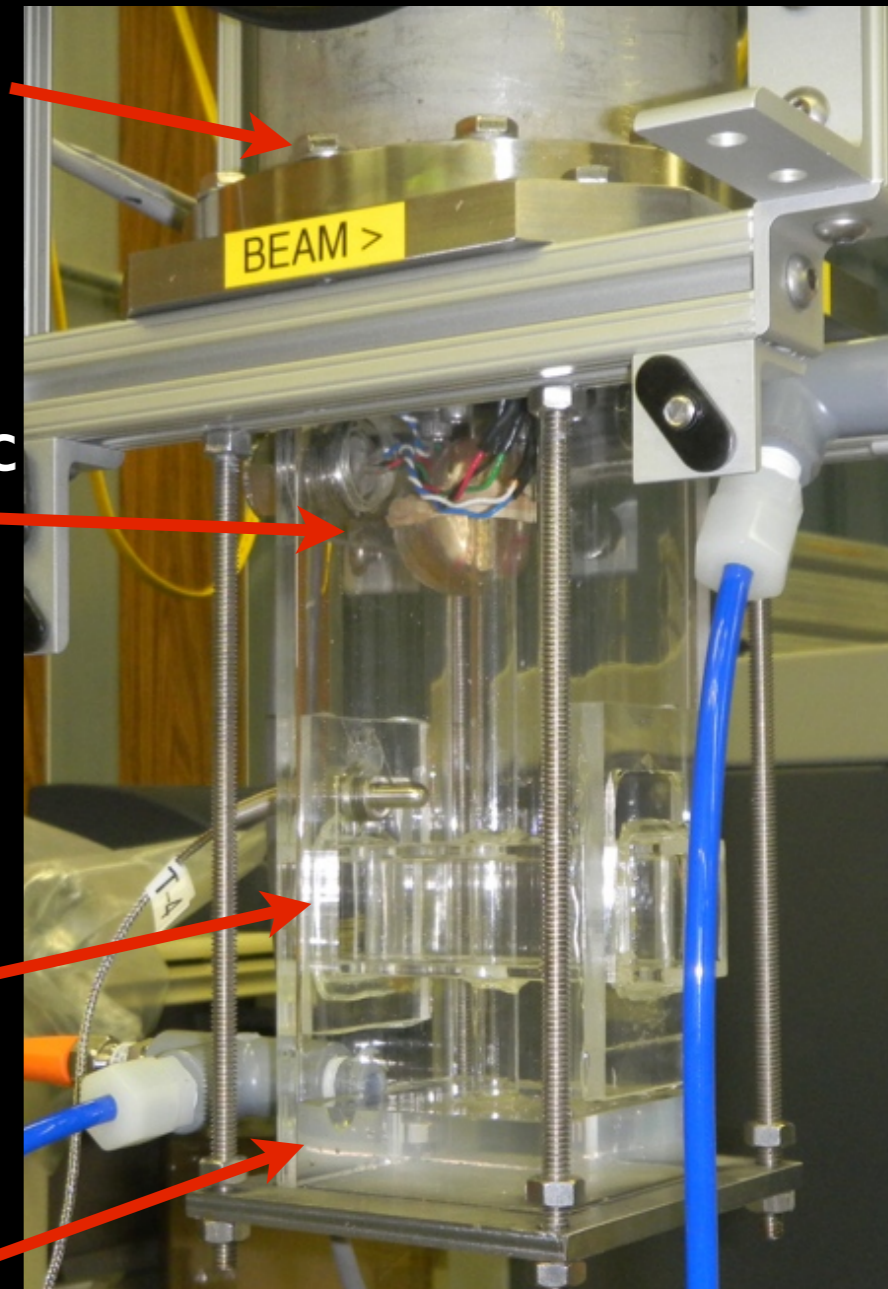
$$T = E_{recoil} = \frac{(p\theta)^2}{2m_r}$$

COUPP Iodine Recoil Threshold Experiment

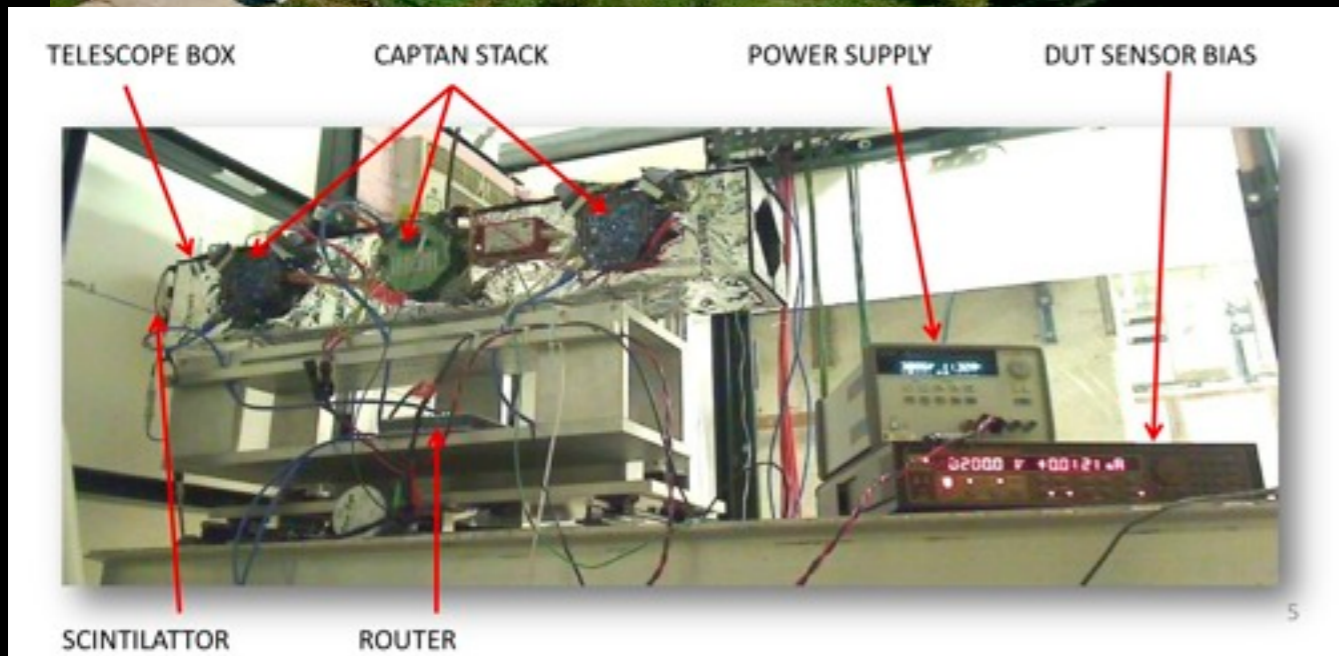
- Test beam at Fermilab with a silicon pixel telescope
- Designed a new test tube sized bubble chamber



Hydraulics
Piezo-acoustic sensor

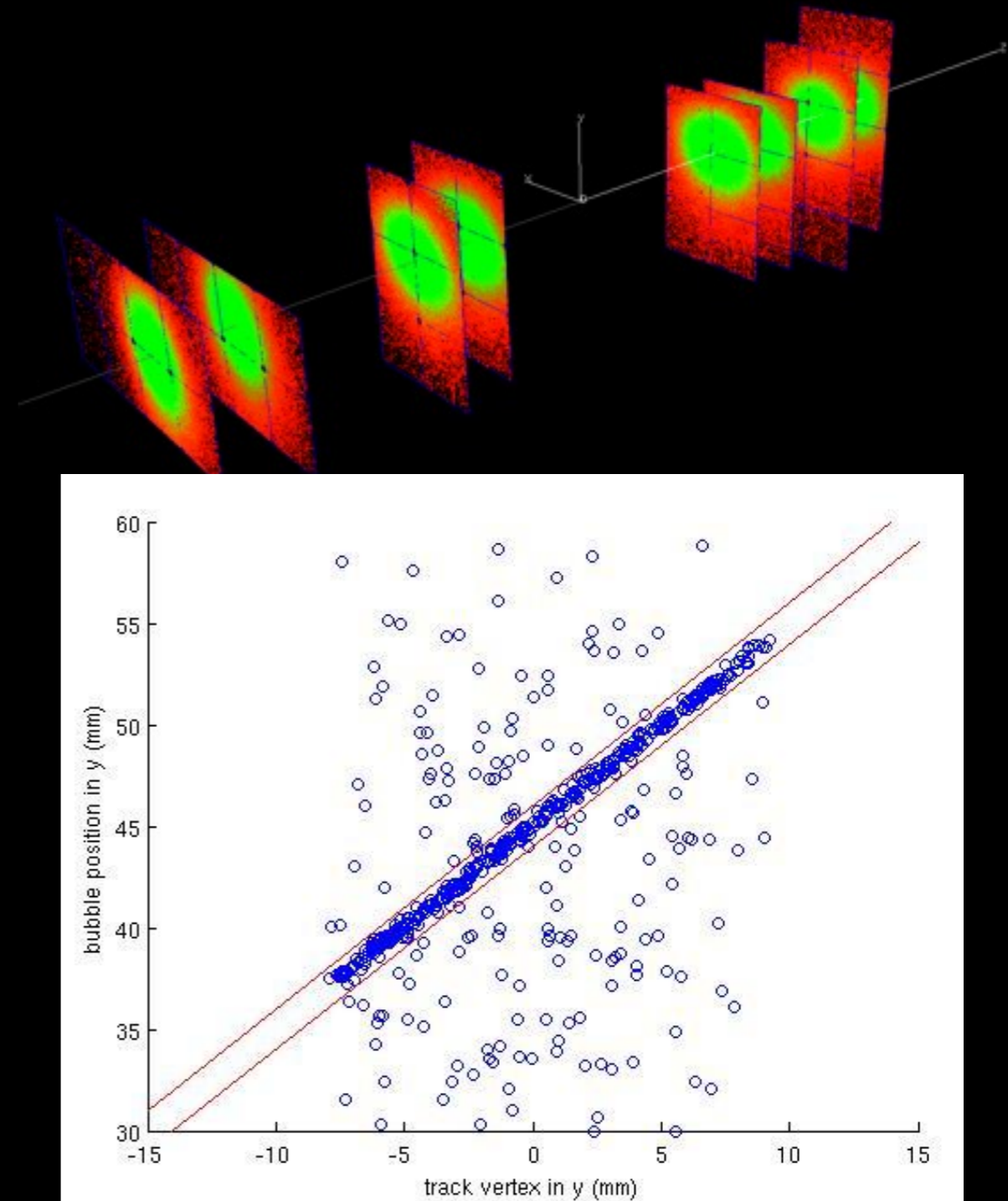
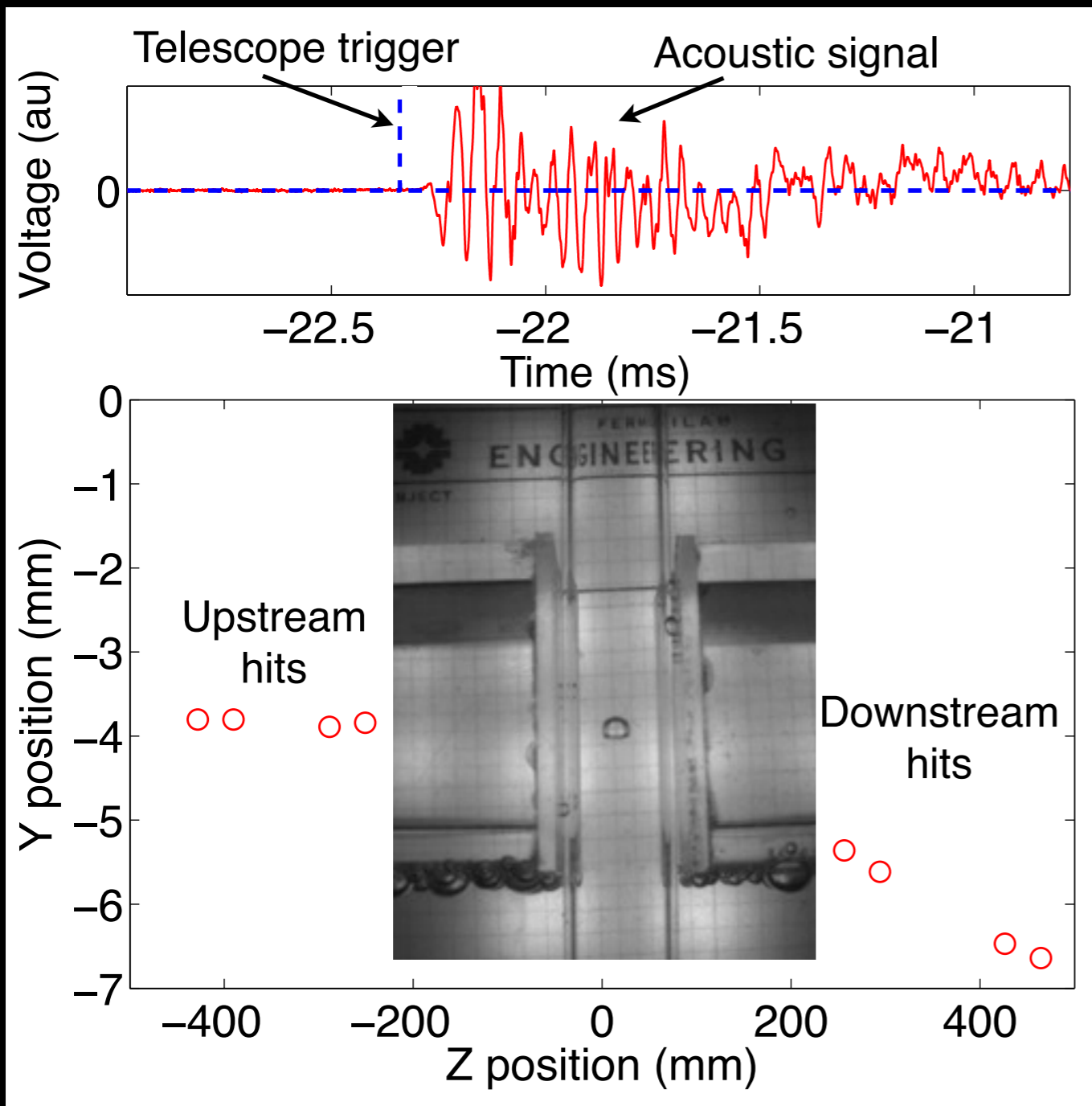


Beam tube
Water bath



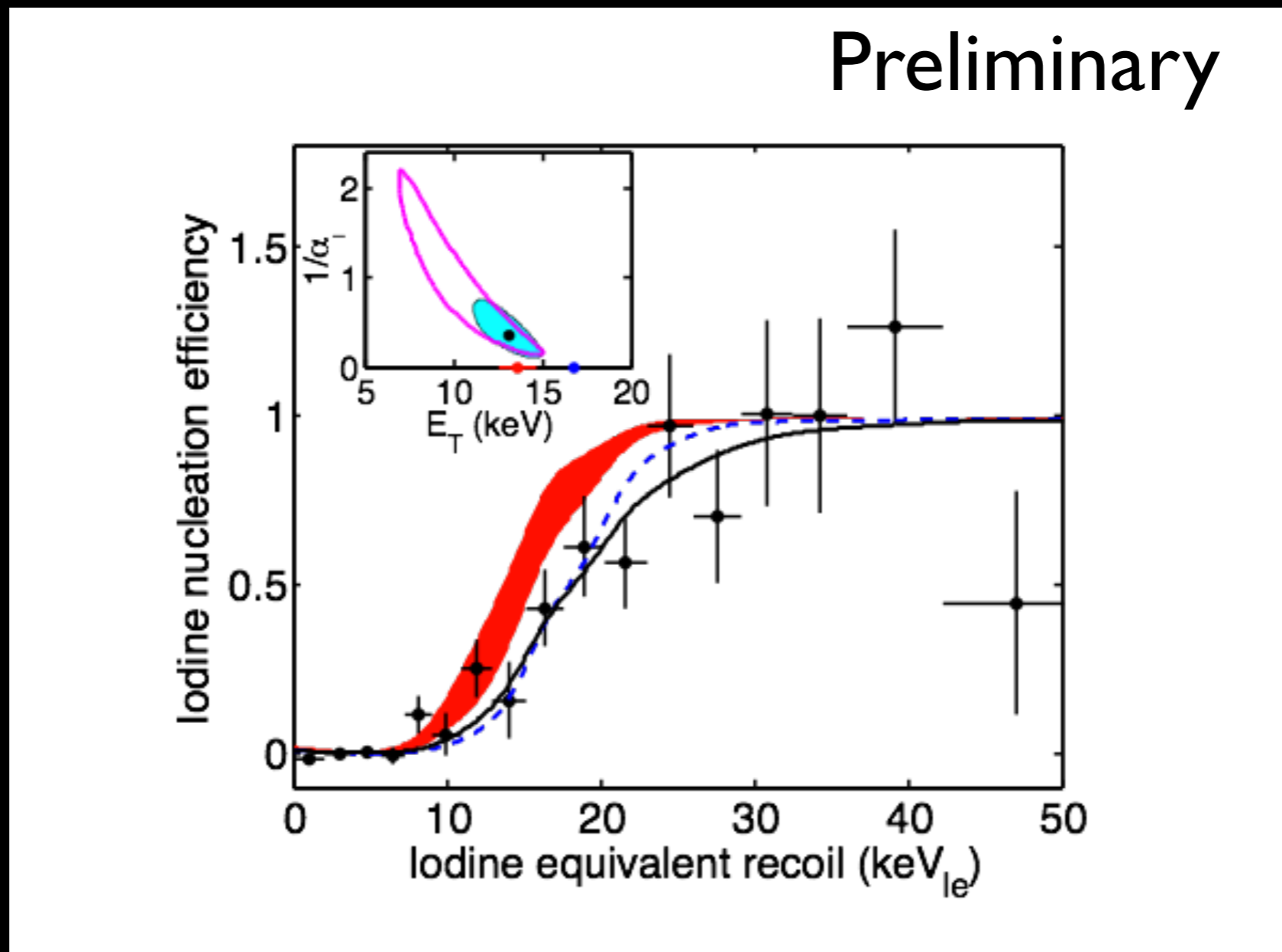
COUPP Iodine Recoil Threshold Experiment

- Beam run at Fermilab in March, 2012



COUPP Iodine Recoil Threshold Experiment

- Analysis shows that iodine threshold is very close to a step function at the predicted energy
 - Limited by resolution (MCS) and statistics



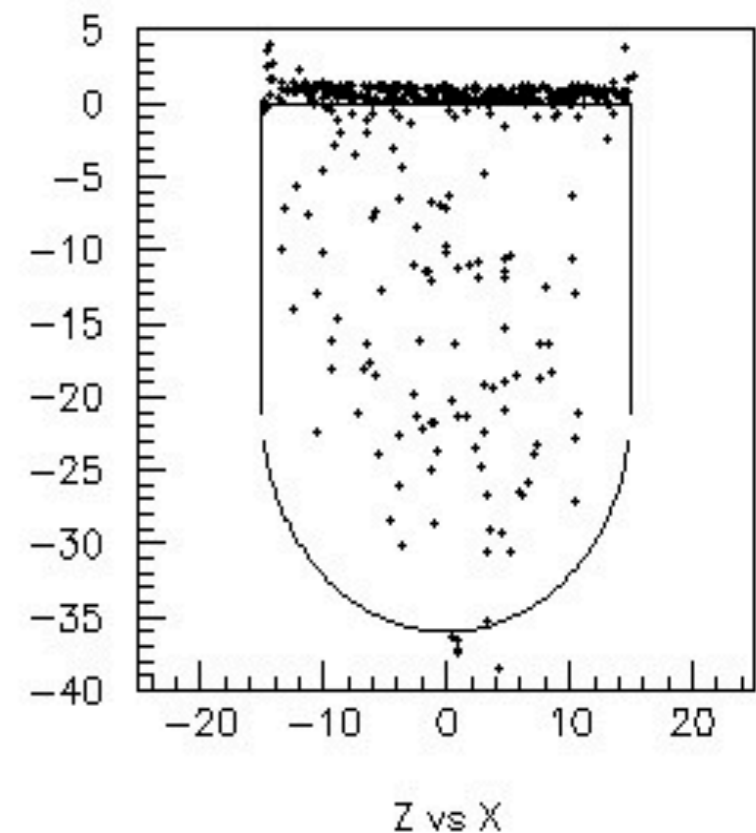
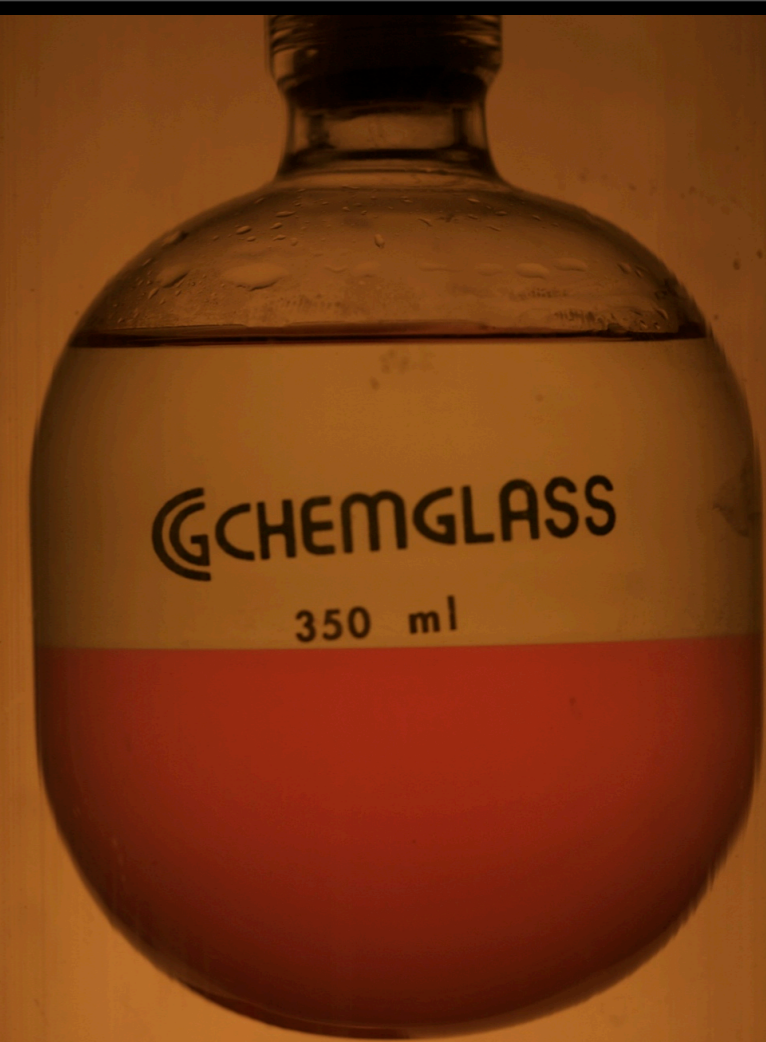
COUPP60

- Engineering run at shallow site in 2010
- Low backgrounds and acoustic discrimination



COUPP60

- Engineering run at shallow site in 2010
- Low backgrounds and acoustic discrimination
- Fluid darkening due to photodissociation of iodine
- Excessive surface rate



COUPP60

- Engineering run at shallow site in 2010
- Low backgrounds and acoustic discrimination
- Fluid darkening due to photodissociation of iodine
- Excessive surface rate
- Solutions tested in second run November, 2011
- Moving to SNOLAB since last summer

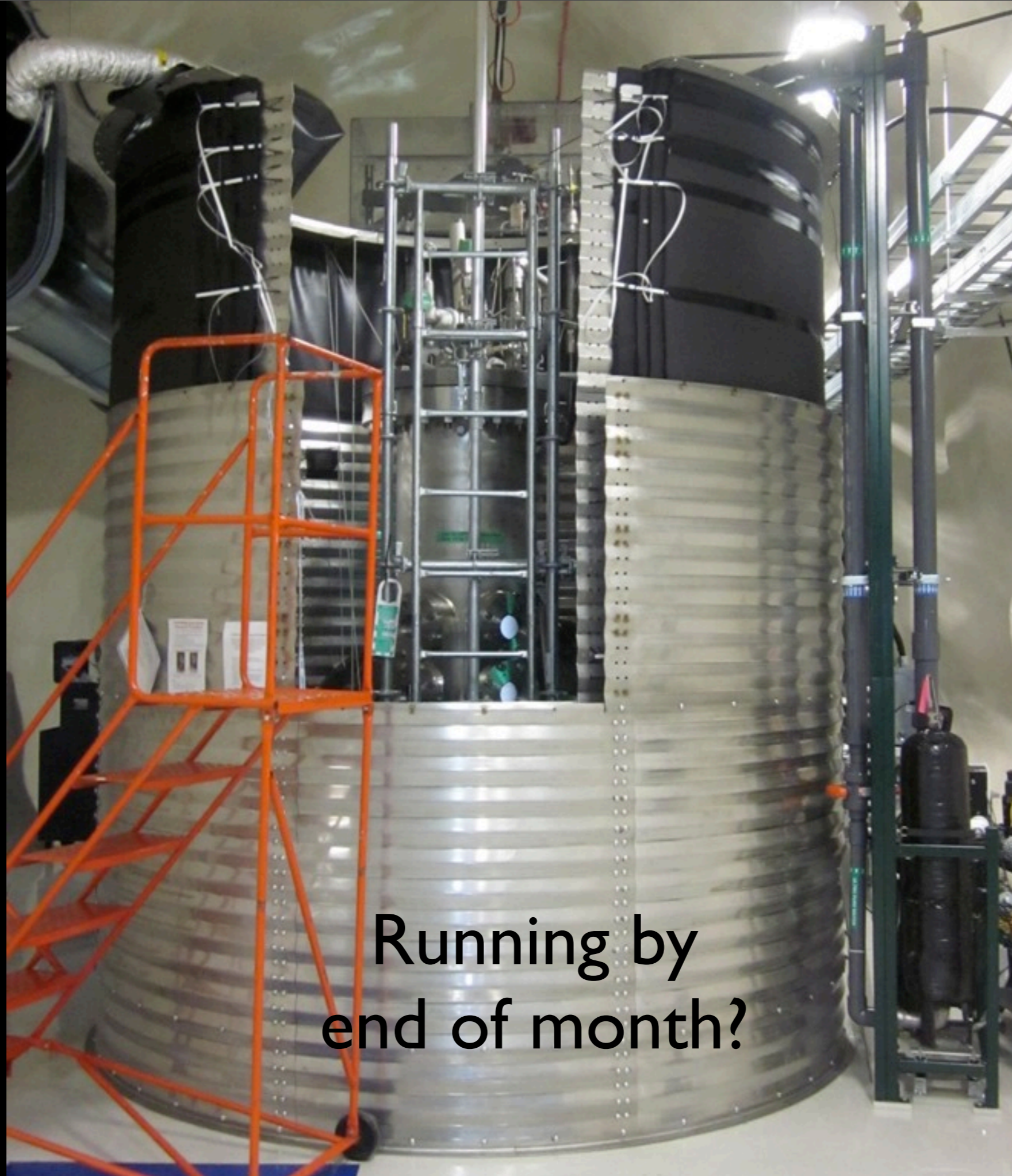




Tuesday, April 16, 2013

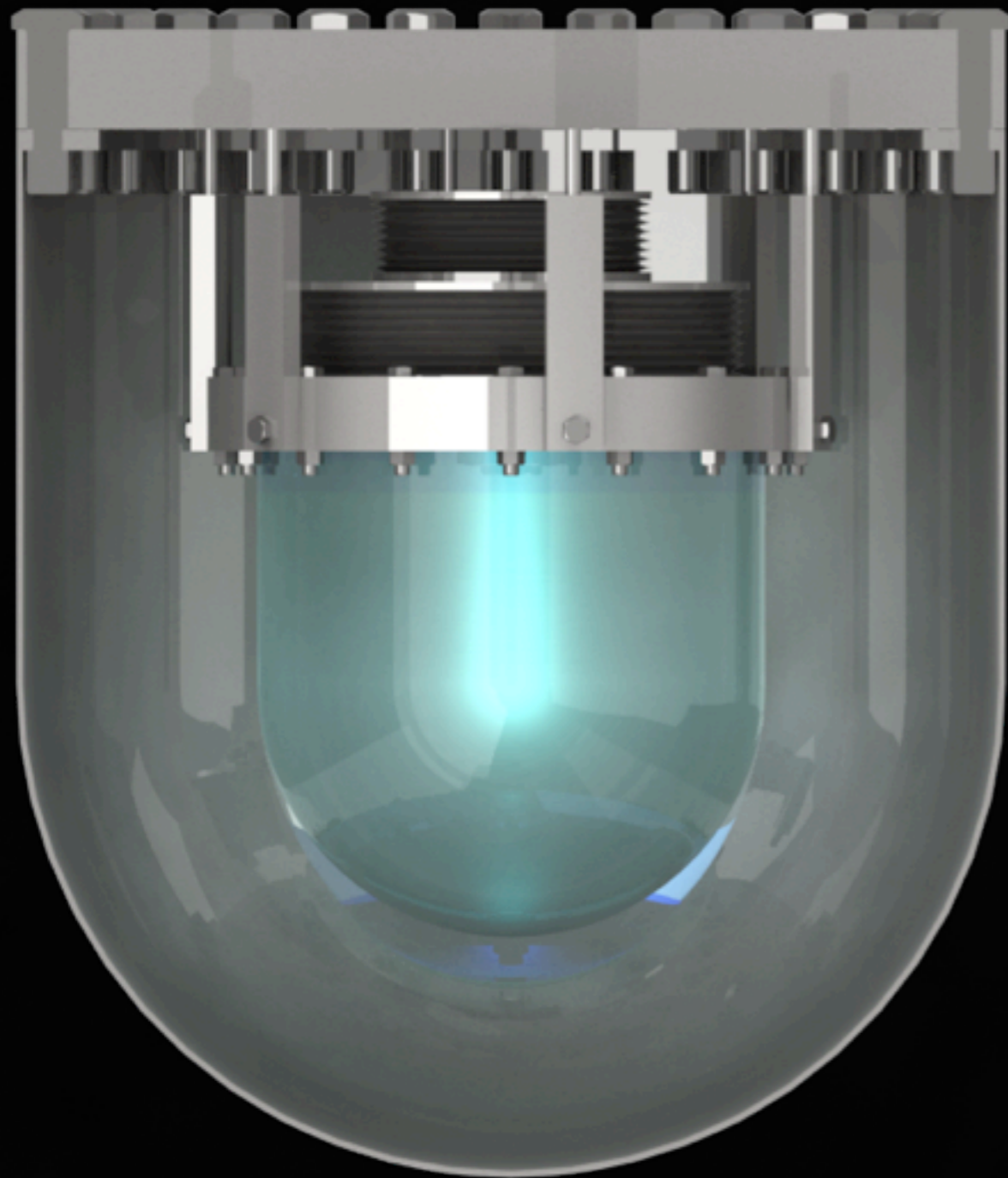


Tuesday, April 16, 2013



Running by
end of month?

COUPP500 (or EREBOS-500)



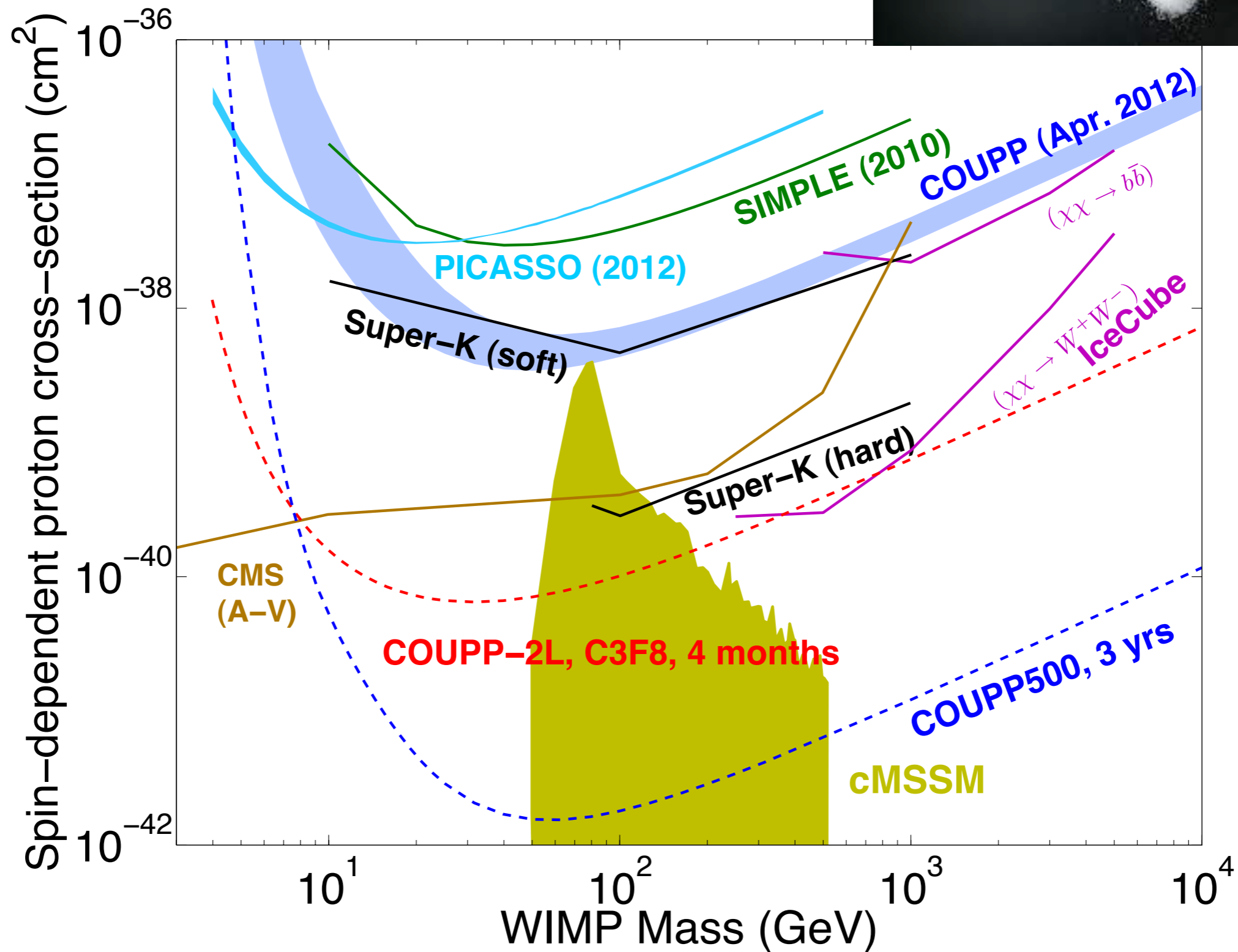
- New merger with Canadian PICASSO collaboration (recent vote chose EREBOS as the new name)
- Funded by NSF and DOE as part of G2 (big showdown in October?)
- Engineering well underway
- Construction 2014-2015?



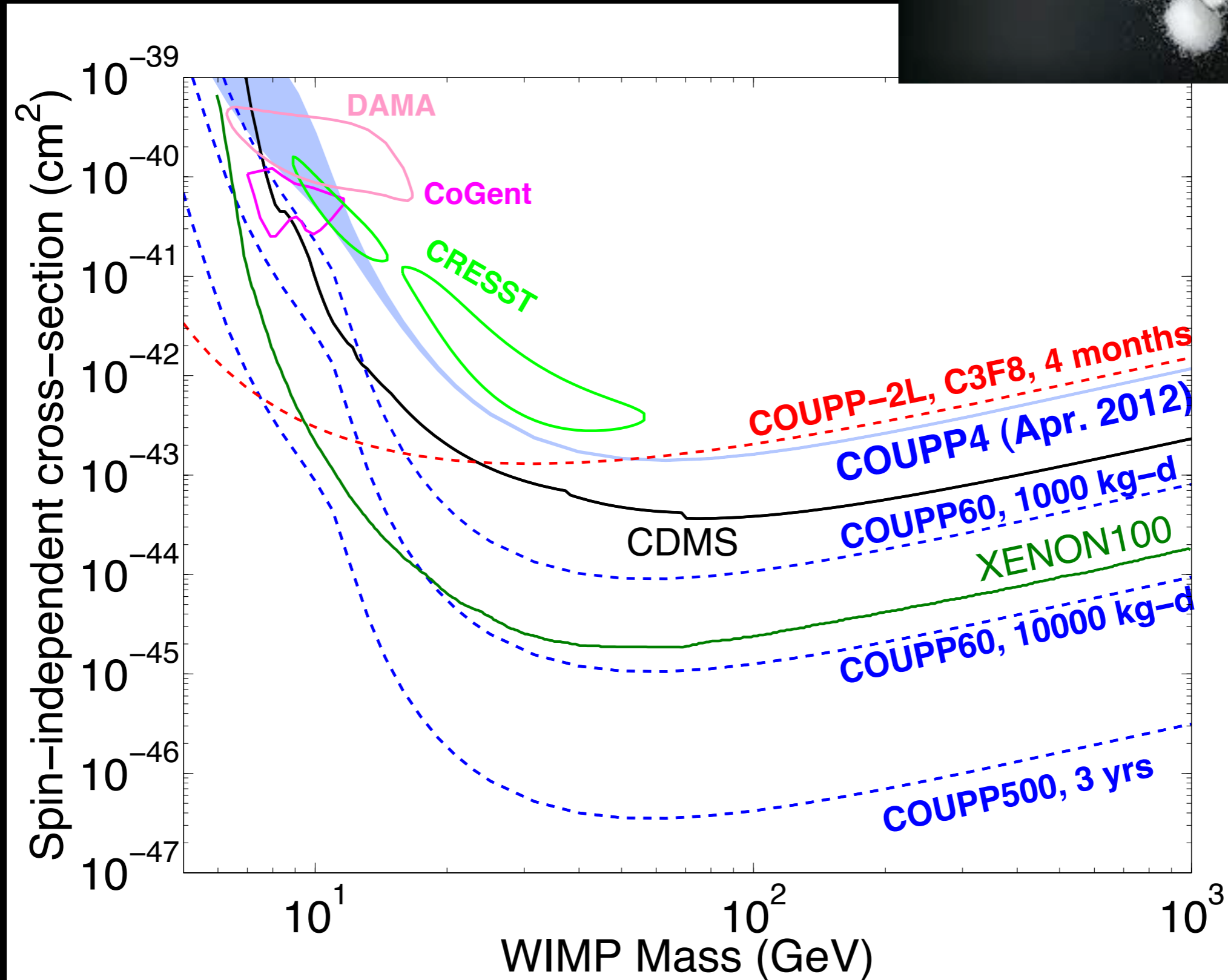
COUPP4 redux

- Alternate fluid - remove the iodine - C_3F_8
 - Lower threshold (down to 3 keV in test stand)
 - Improved sensitivity at low WIMP mass
 - Improved SD sensitivity
- First effort in concert with the PICASSO collaboration
- Possible use in COUPP500 chamber

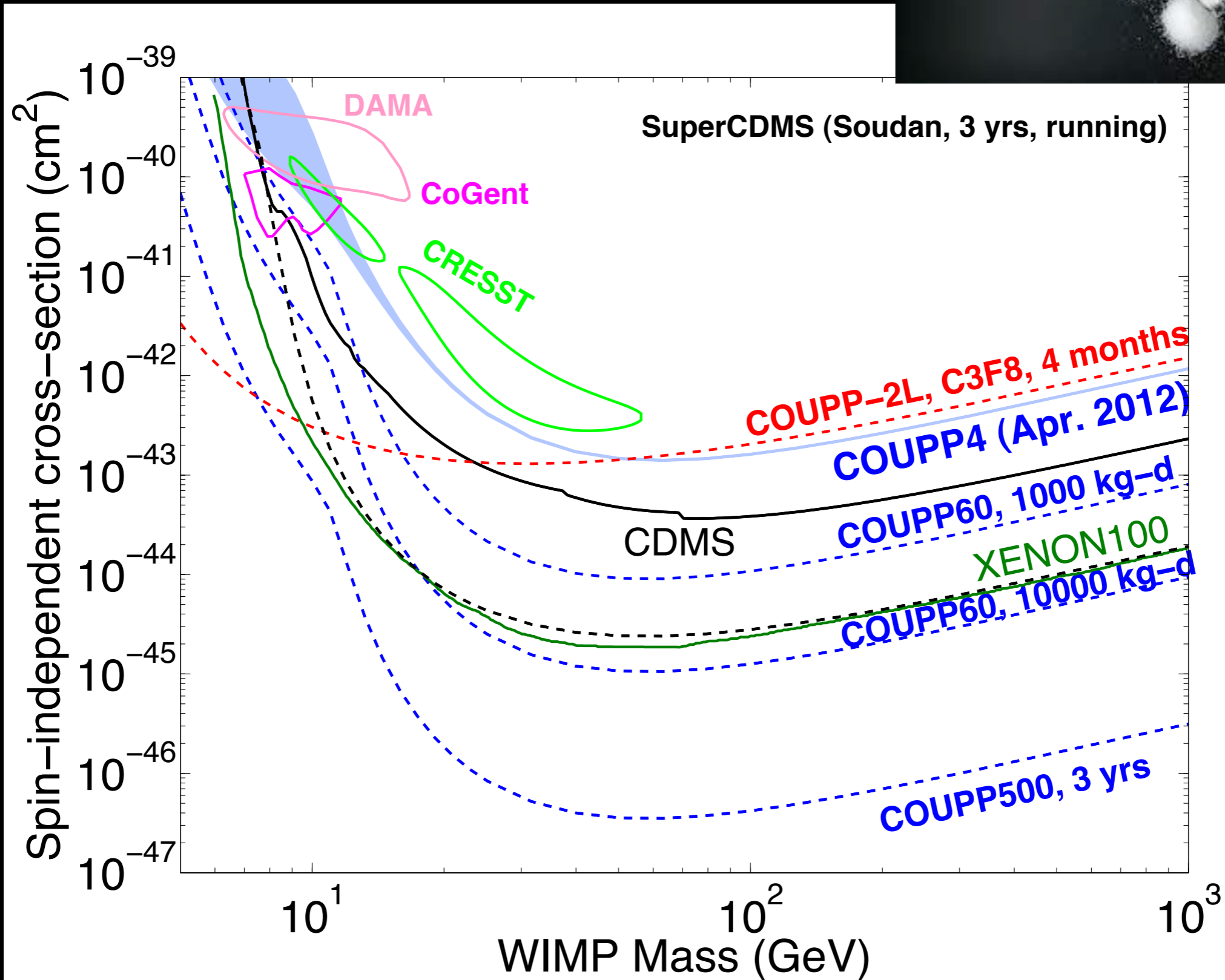
Projections



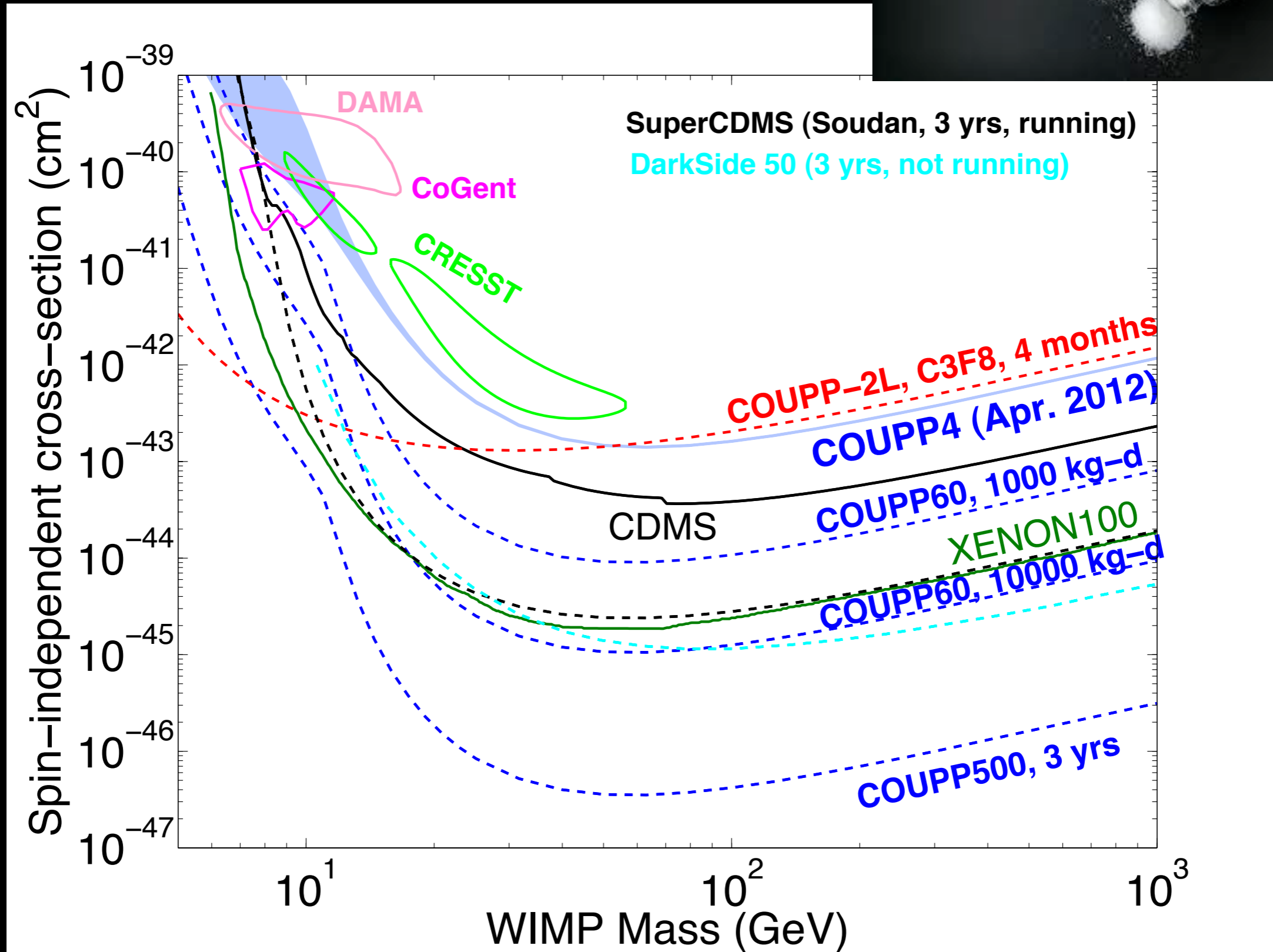
Projections



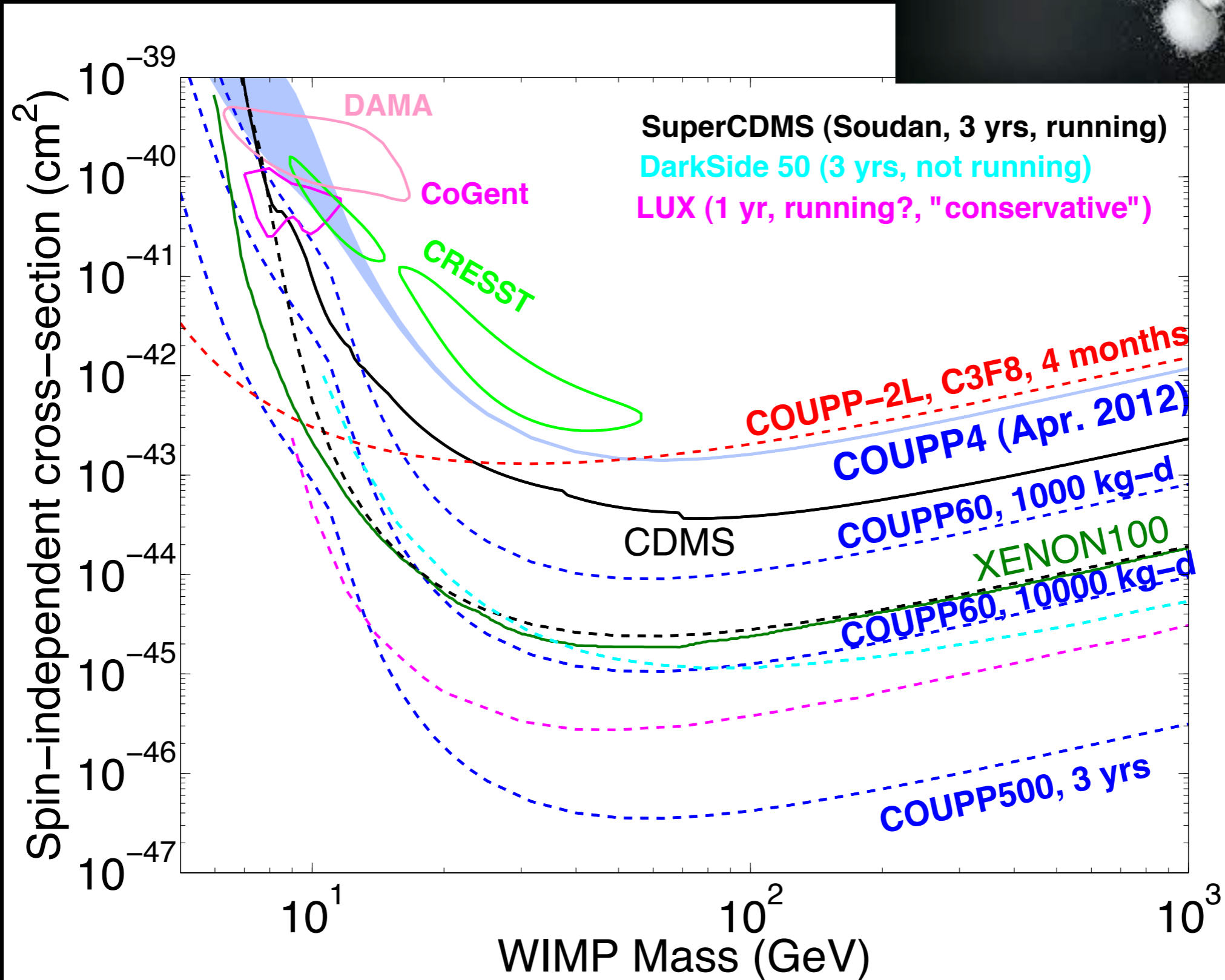
Projections



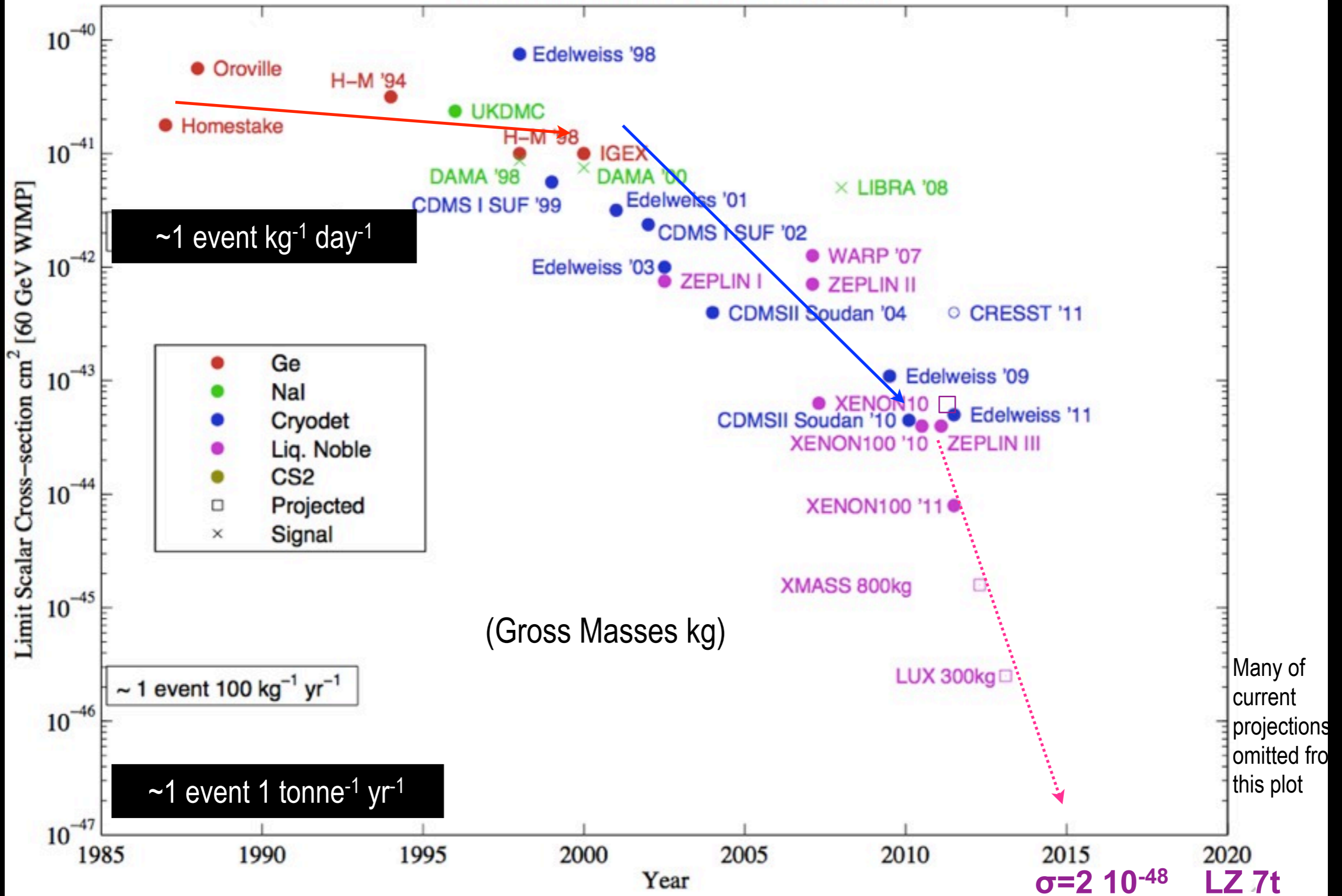
Projections



Projections

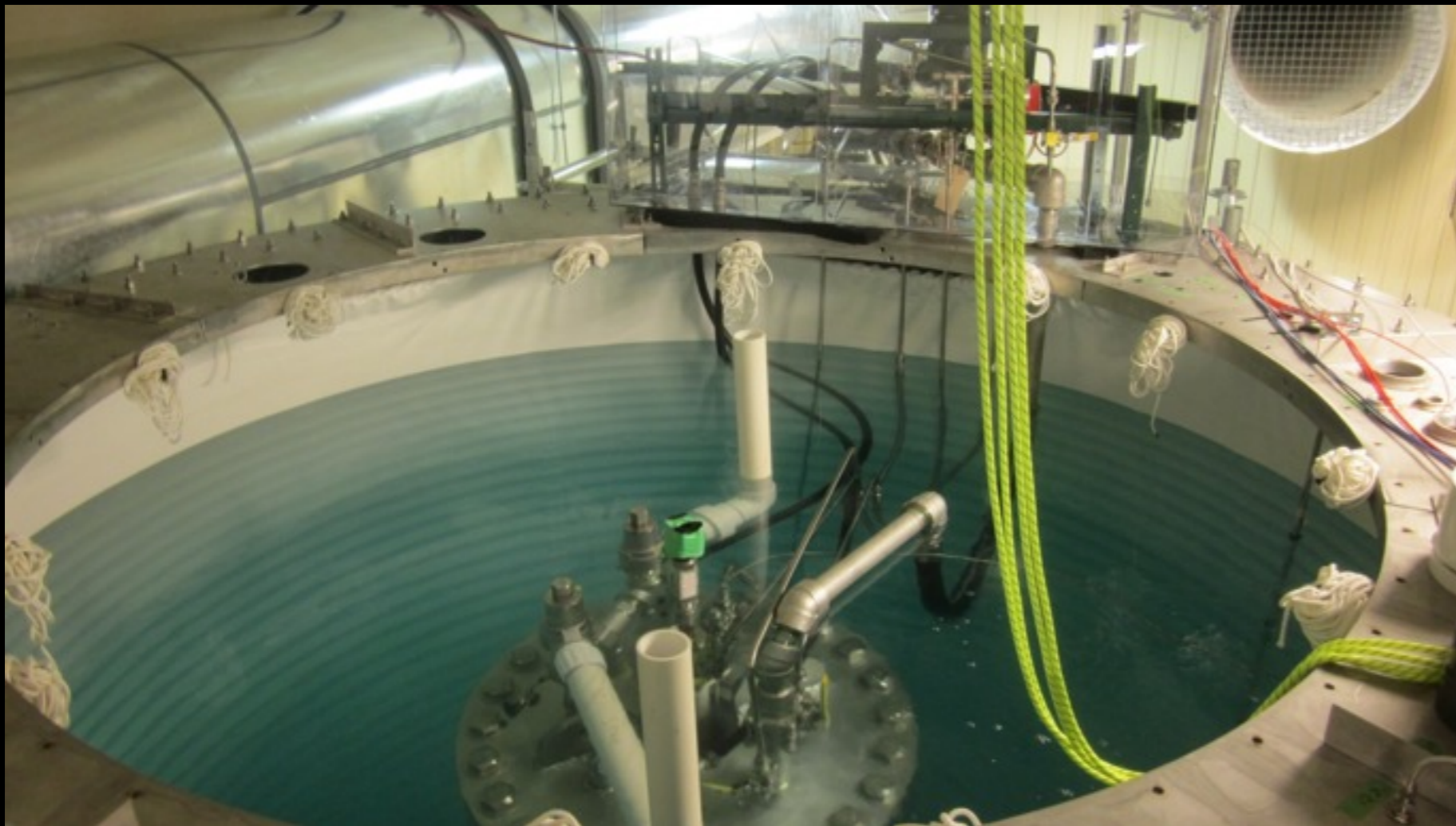


Dark Matter Searches: Past, Present & Future



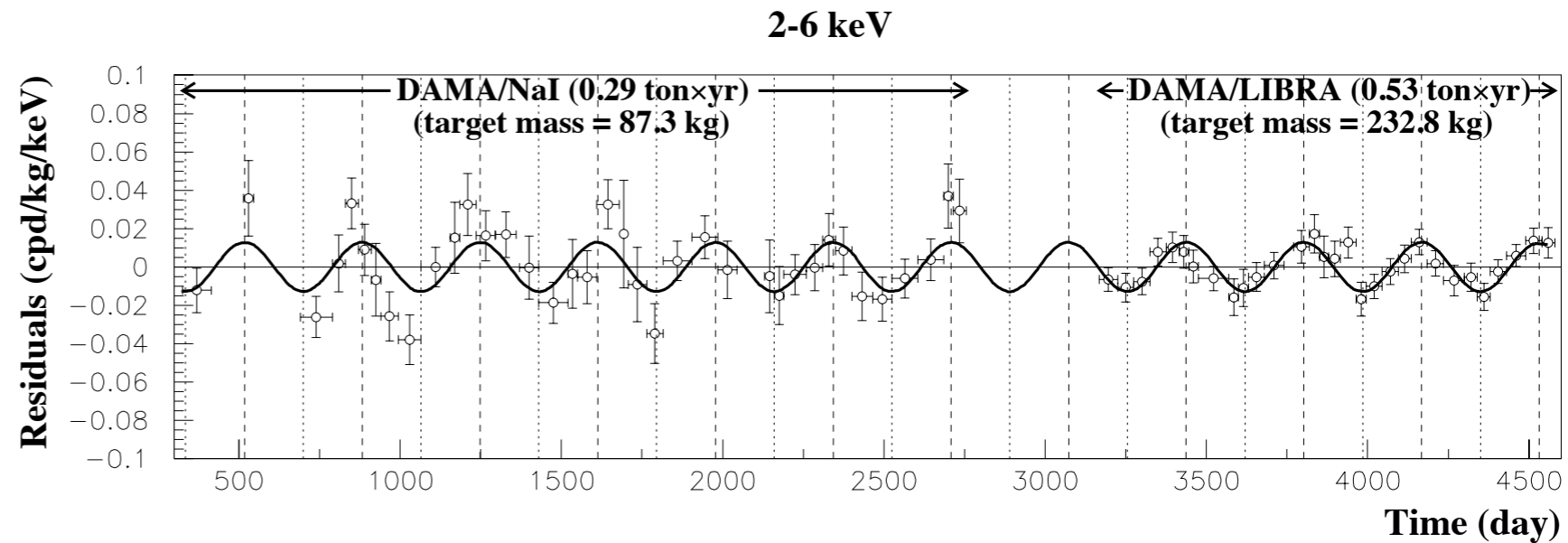
Conclusion

- Dark matter searches are making fast progress (indirect, accelerator and direct)
- COUPP is producing the best direct detection limits on spin-dependent dark matter
- COUPP bubble chambers are also competitive for spin-independent searches



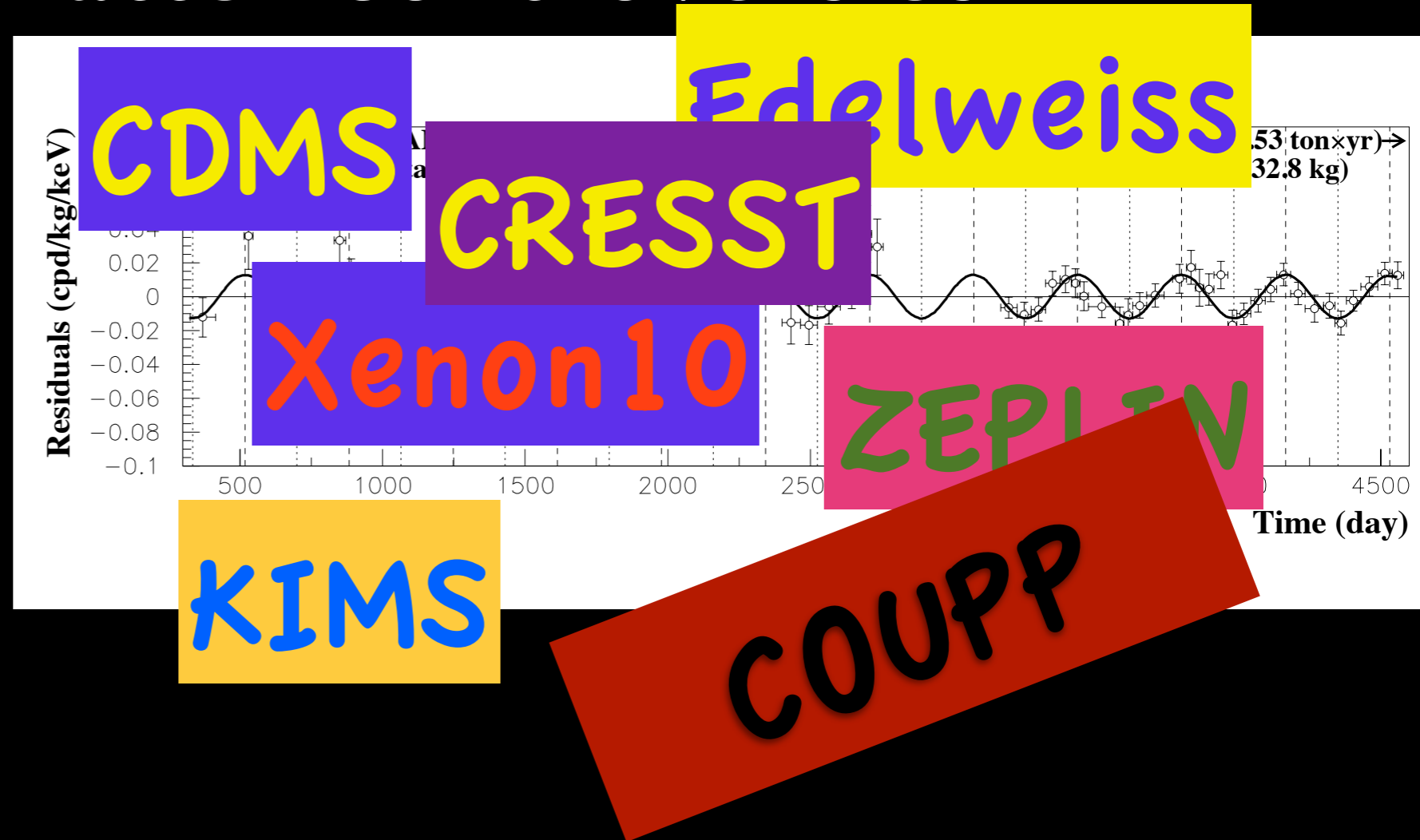
Dark matter controversies

**DAMA -
positive claim
for 10 years!**



Dark matter controversies

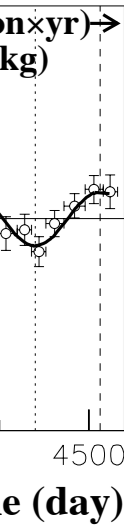
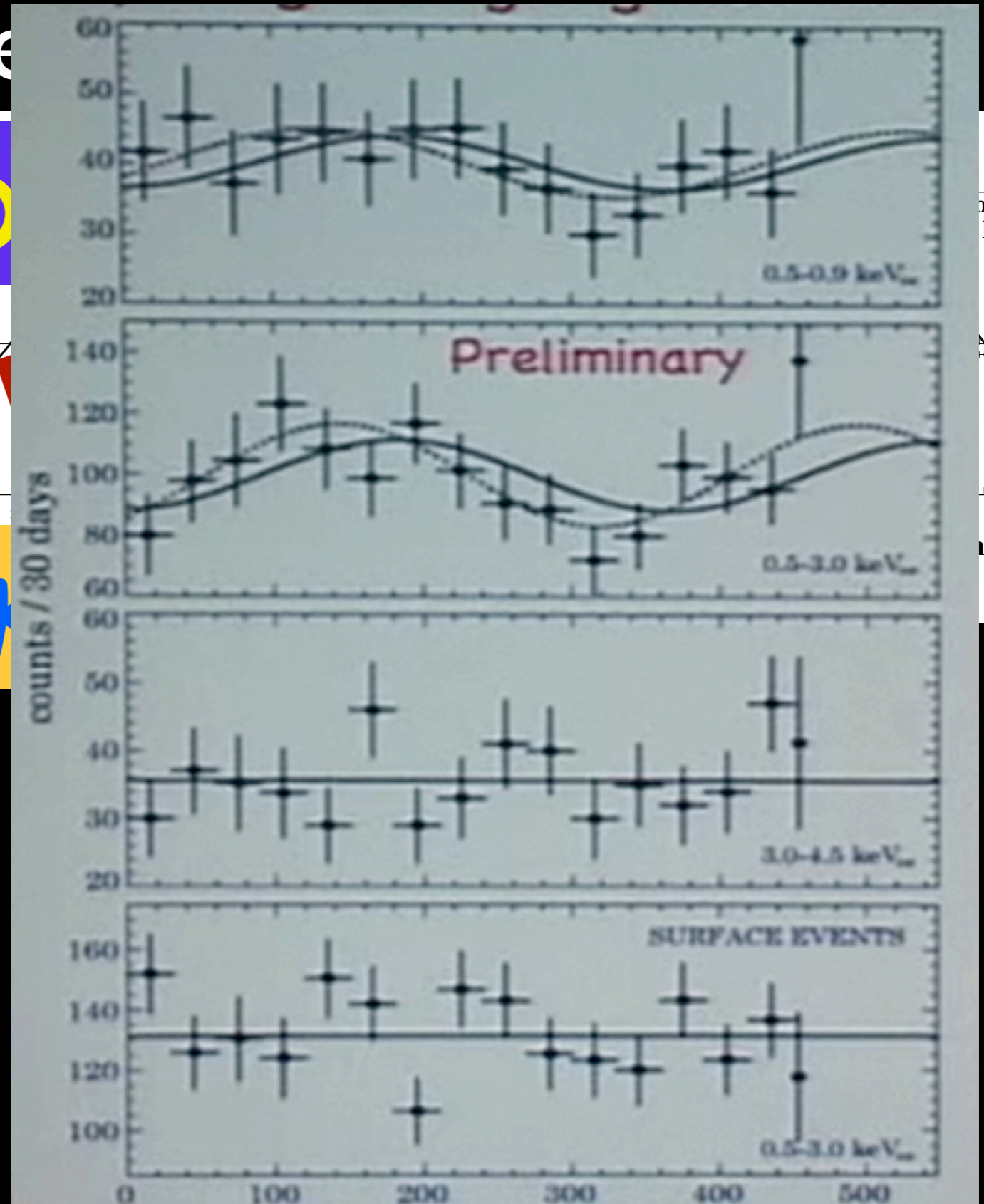
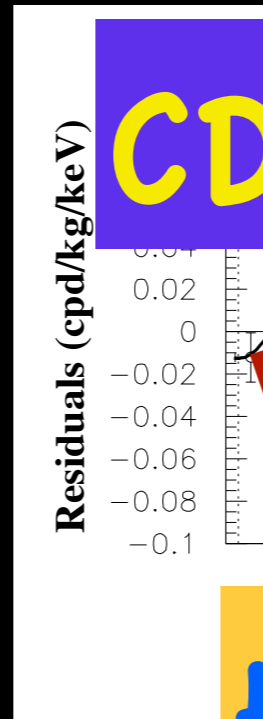
DAMA -
positive claim
for 10 years!



Dark matter

DAMA -
positive claim
for 10 years!

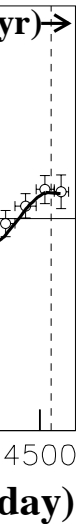
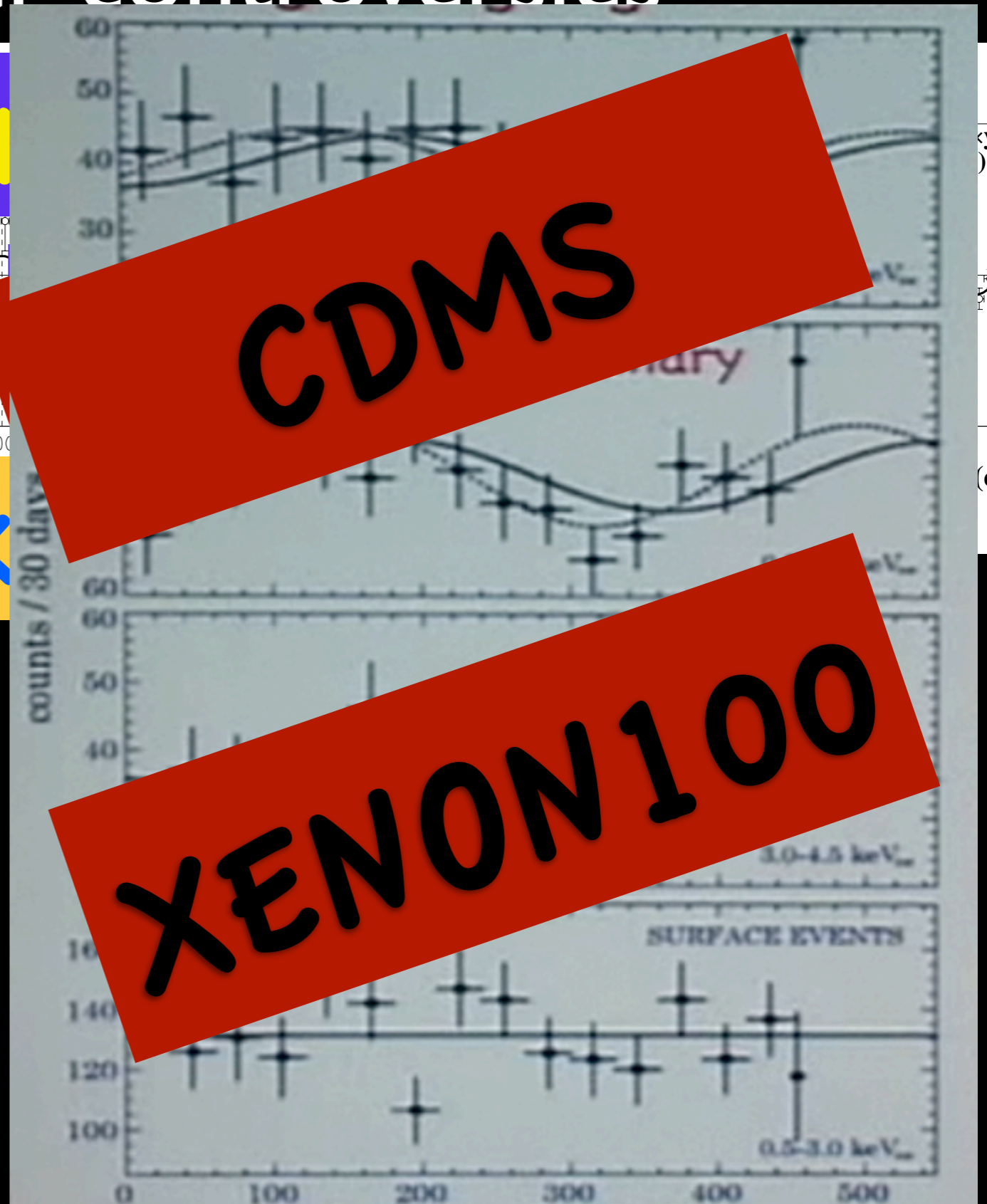
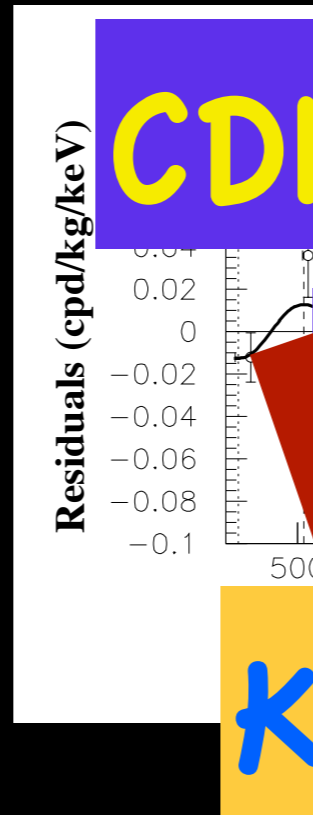
A few years ago,
CoGeNT saw an
excess and then a
possible annual
modulation



Dark matter controversies

DAMA -
positive claim
for 10 years!

CoGeNT reanalyzed
their own data and
found a new
background, decreasing
the sensitivity



DAMA
positive
for 10 ye

Recently
(run by m
an excess
possibl
n

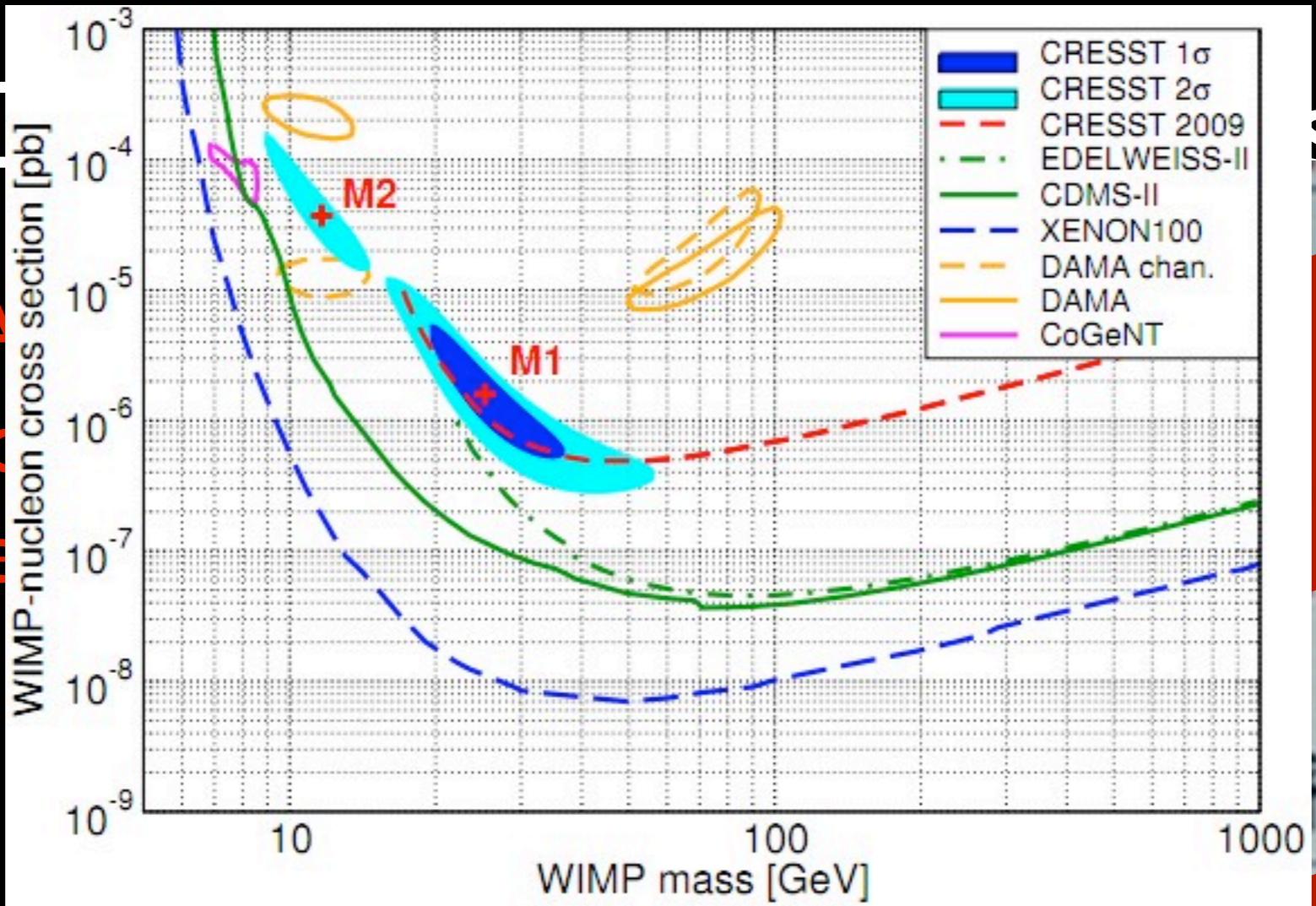
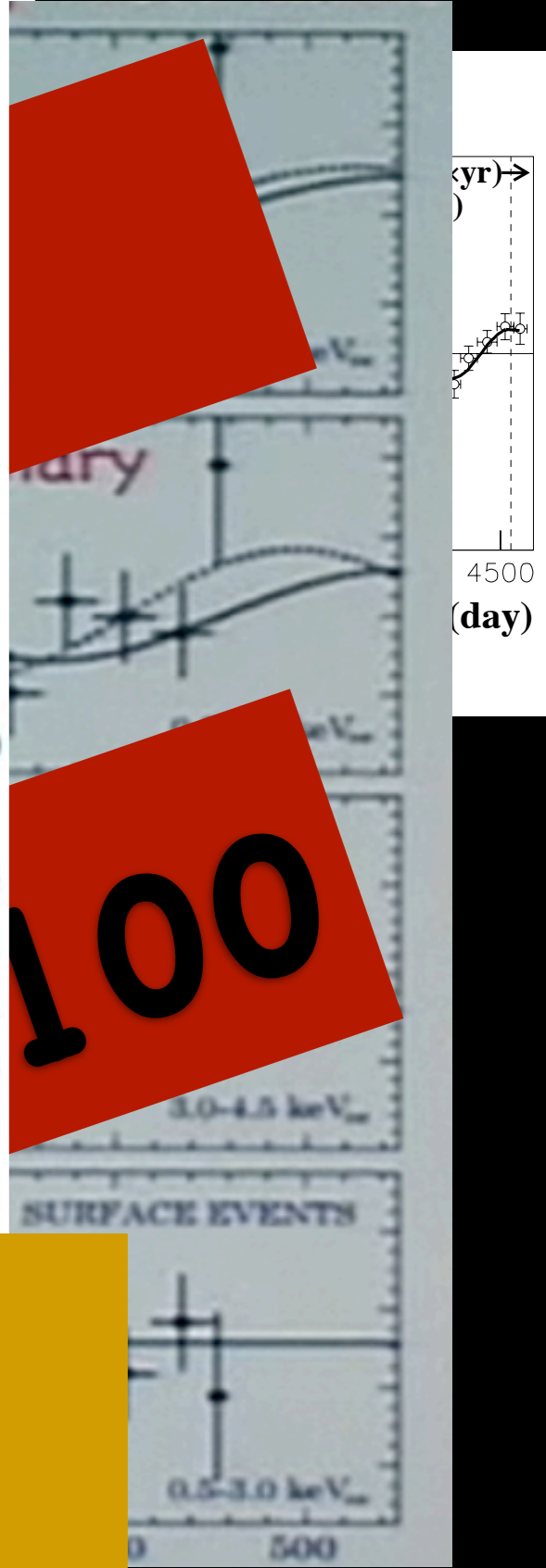


Fig. 13. The WIMP parameter space compatible with the CRESST results discussed here, using the background model described in the text, together with the exclusion limits from CDMS-II [12], XENON100 [13], and EDELWEISS-II [14], as well as the CRESST limit obtained in an earlier run [1]. Additionally, we show the 90% confidence regions favored by CoGeNT [15] and DAMA (LIPDA) [16] (with $\chi^2_{\text{min}} = 1.1$ and 1.2 , respectively).

CRESST



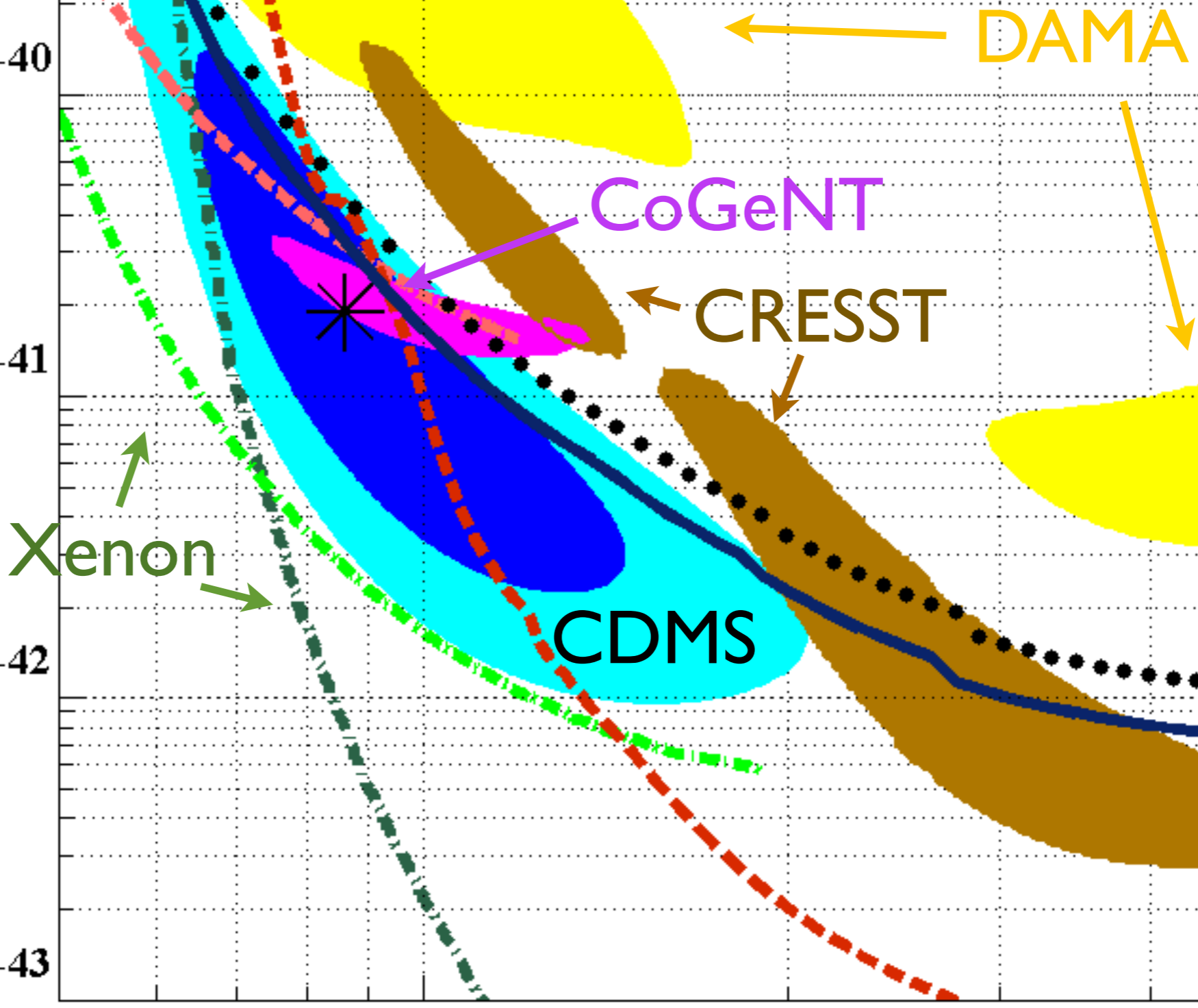
- At APS on Saturday, CDMS (which has historically been the most conservative experiment, culturally speaking) announced a result that is consistent with a light WIMP hypothesis (best fit at 8.6 GeV, $1.9 \times 10^{-41} \text{ cm}^2$)
- 3 candidate events over an estimated background of ~ 0.7 . WIMP hypothesis fits with p-value of 68%, background only at 4.5%
- Nuclear recoil events (CDMS has discrimination, unlike DAMA or CoGeNT)

As of Saturday...

WIMP-nucleon cross-section [cm^2]

10^{-40}
 10^{-41}
 10^{-42}
 10^{-43}

10^1
WIMP Mass [GeV/c^2]



DAMA

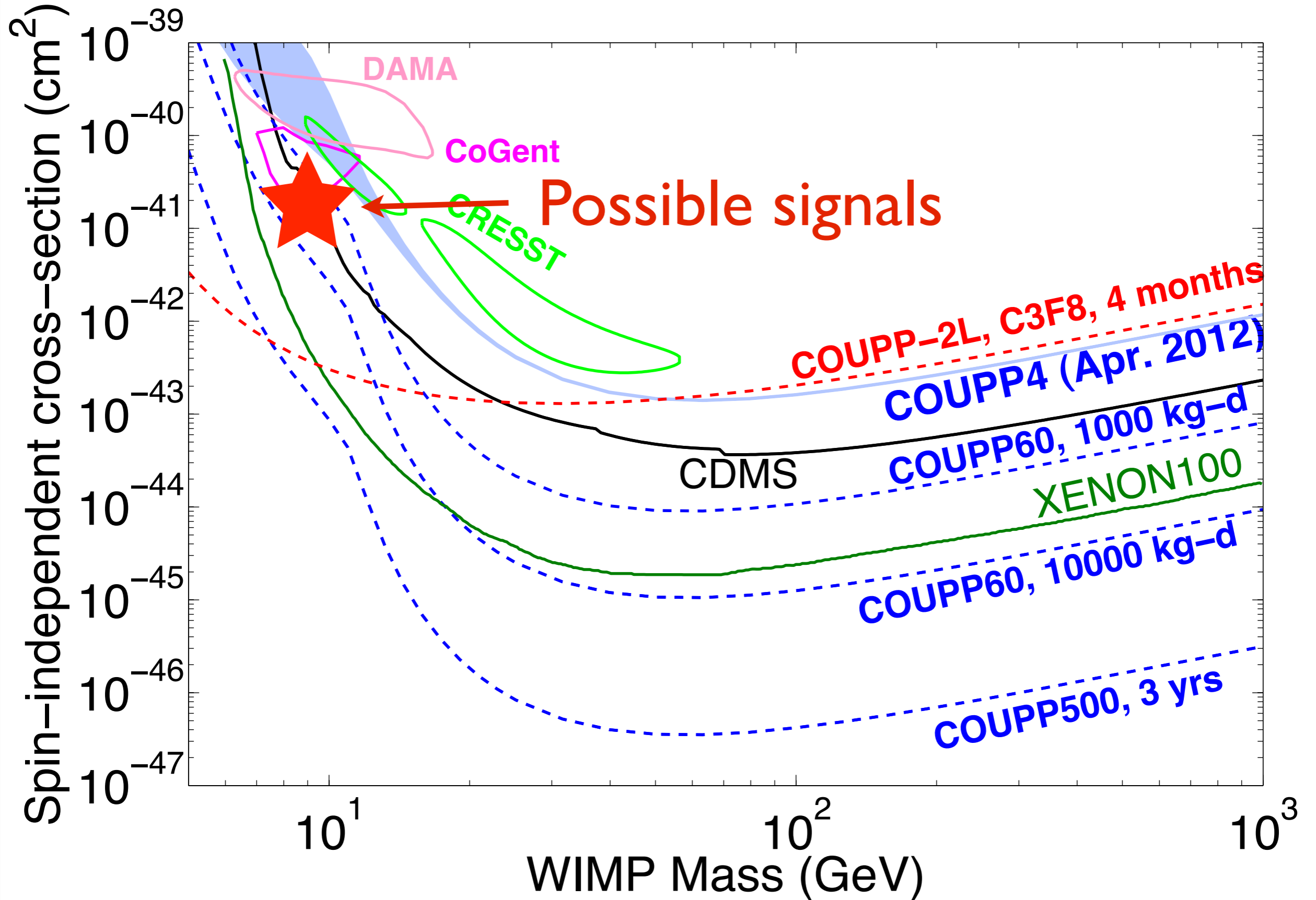
CoGeNT

CRESST

CDMS

Xenon

We'll see?



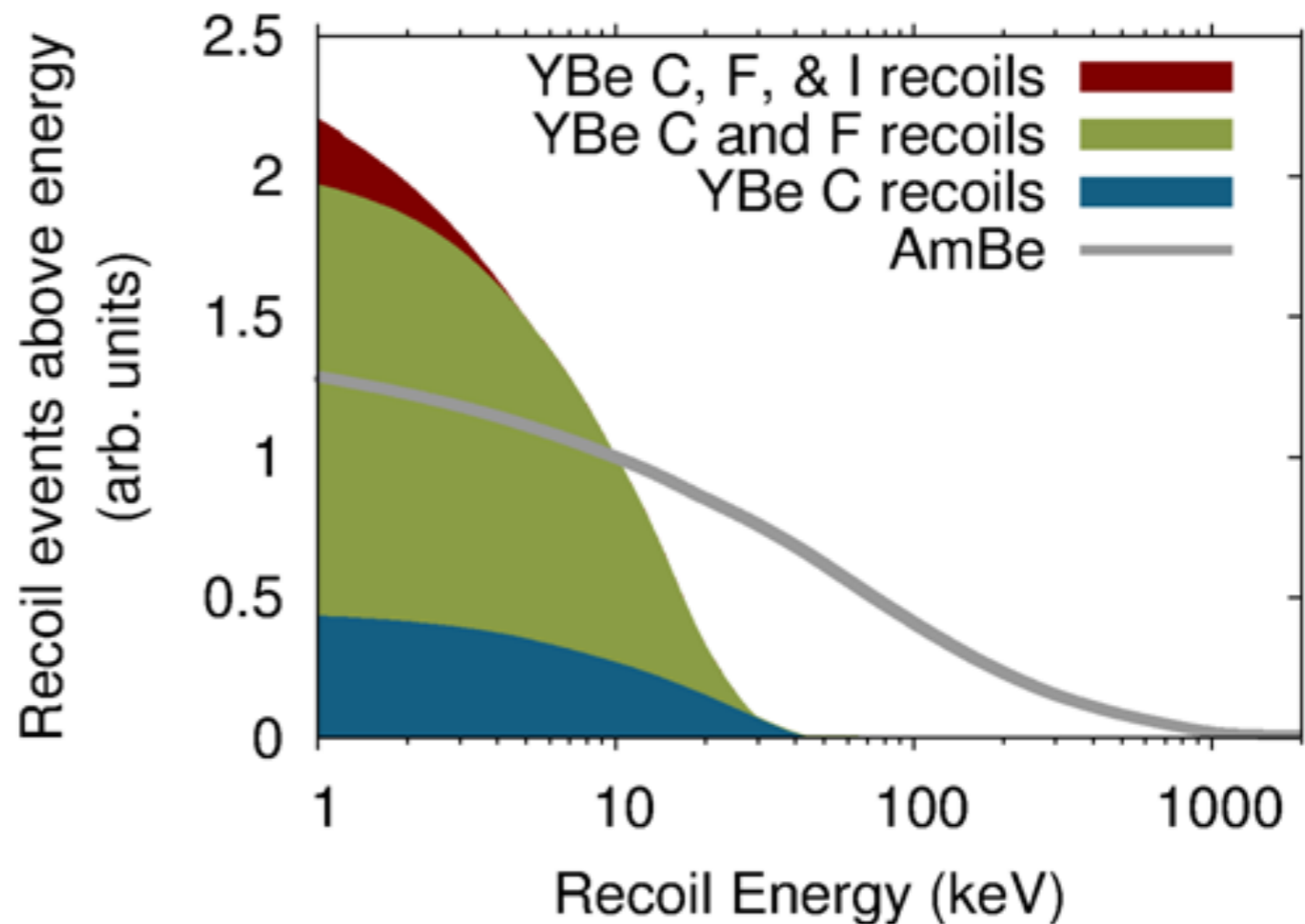
^{88}YBe (γ, n) neutron source

Mono-energetic 152 keV
neutron source.

TABLE 3. Results of present measurements

E_γ (keV)	$\sigma(E_\gamma)$ (mb)
1674.7	0.88 ± 0.16
1705.2	1.33 ± 0.24
1724.9	1.10 ± 0.20
1778.9	0.73 ± 0.13
1836.0	0.47 ± 0.09
2167.6	0.18 ± 0.04

M. Fujishiro et al., Can. J. Phys. **60**,
1672 (1982).



WIMP-nucleon scattering

Spin-independent

Spin-dependent

$$\sigma_0 = \frac{4\mu^2}{\pi} [f_p N_p + f_n N_n]^2 + \frac{32G_F^2 \mu^2}{\pi} \frac{J+1}{J} [a_p \langle S_p \rangle + a_n \langle S_n \rangle]^2$$

Nucleus	Z	Odd Nucleon	J	$\langle S_p \rangle$	$\langle S_n \rangle$	C_A^p/C_p	C_A^n/C_n
^{19}F	9	p	1/2	0.477	-0.004	9.10×10^{-1}	6.40×10^{-5}
^{23}Na	11	p	3/2	0.248	0.020	1.37×10^{-1}	8.89×10^{-4}
^{27}Al	13	p	5/2	-0.343	0.030	2.20×10^{-1}	1.68×10^{-3}
^{29}Si	14	n	1/2	-0.002	0.130	1.60×10^{-5}	6.76×10^{-2}
^{35}Cl	17	p	3/2	-0.083	0.004	1.53×10^{-2}	3.56×10^{-5}
^{39}K	19	p	3/2	-0.180	0.050	7.20×10^{-2}	5.56×10^{-3}
^{73}Ge	32	n	9/2	0.030	0.378	1.47×10^{-3}	2.33×10^{-1}
^{93}Nb	41	p	9/2	0.460	0.080	3.45×10^{-1}	1.04×10^{-2}
^{125}Te	52	n	1/2	0.001	0.287	4.00×10^{-6}	3.29×10^{-1}
^{127}I	53	p	5/2	0.309	0.075	1.78×10^{-1}	1.05×10^{-2}
^{129}Xe	54	n	1/2	0.028	0.359	3.14×10^{-3}	5.16×10^{-1}
^{131}Xe	54	n	3/2	-0.009	-0.227	1.80×10^{-4}	1.15×10^{-1}