

The White Elephant: Upsilon Physics at the BaBar B-factory

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Programme

- The bottomonium system: prospects for discovery
- The BaBar/PEP-II ~~B~~-factory

$$\begin{aligned}B^0 &= (\bar{b} d) \\B^+ &= (\bar{b} u)\end{aligned}\qquad Y = (\bar{b} b)$$

- A matter of QCD – the η_b
- A matter of new physics – the light Higgs
- Prospects for further discovery

The Bottomonium System: Prospects for Discovery



1977

16 authors

Observation of a Dimuon Resonance at 9.5 GeV in 400-GeV Proton-Nucleus Collisions

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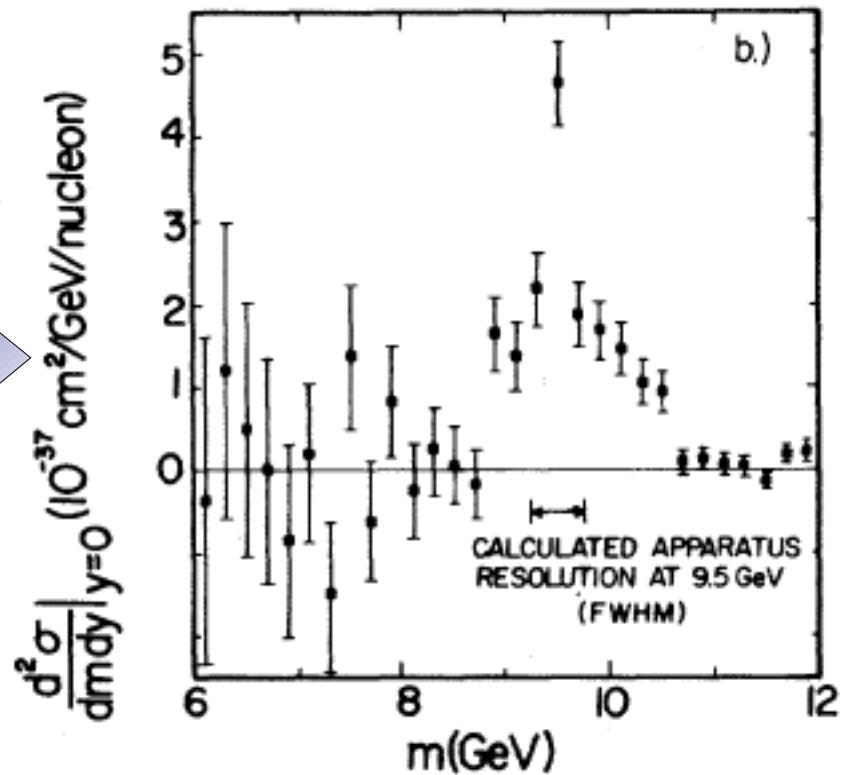
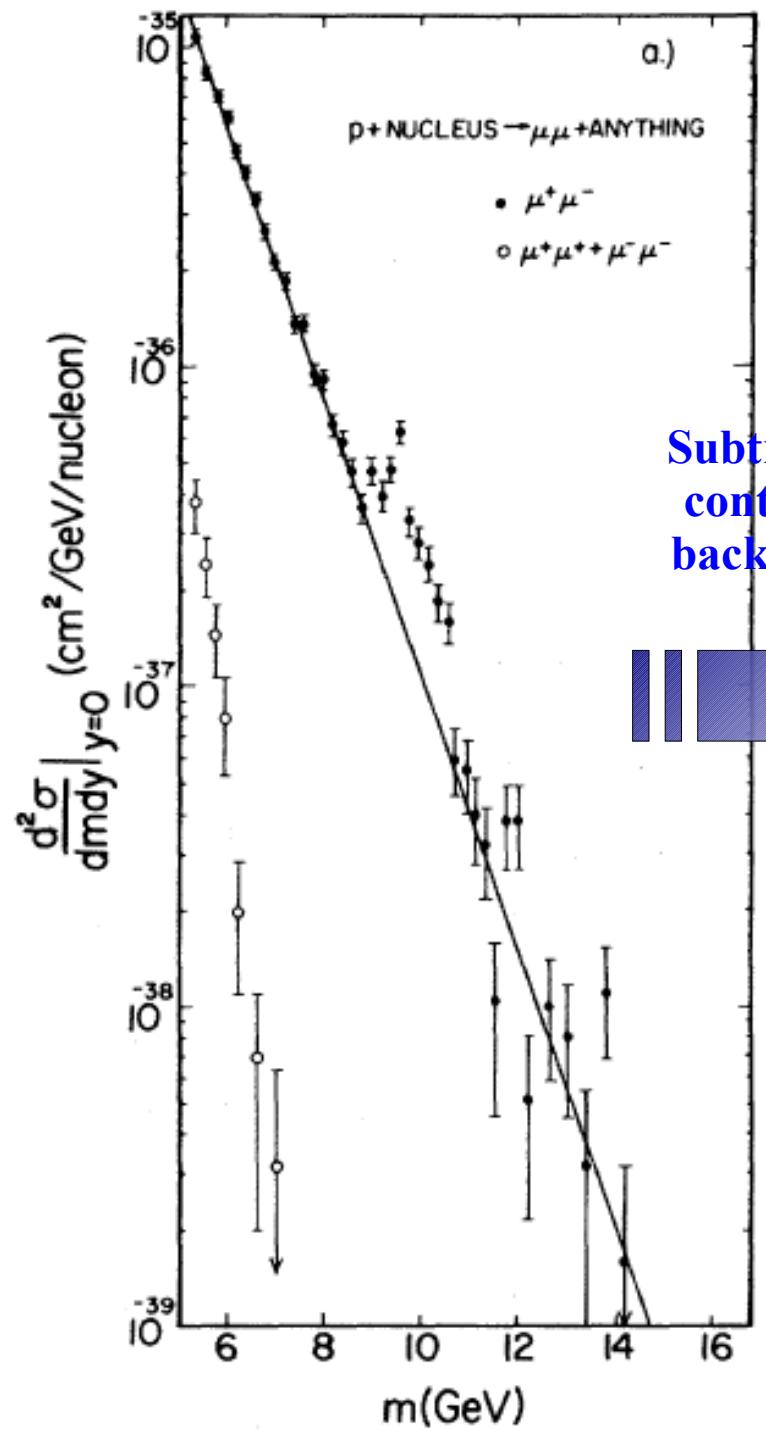
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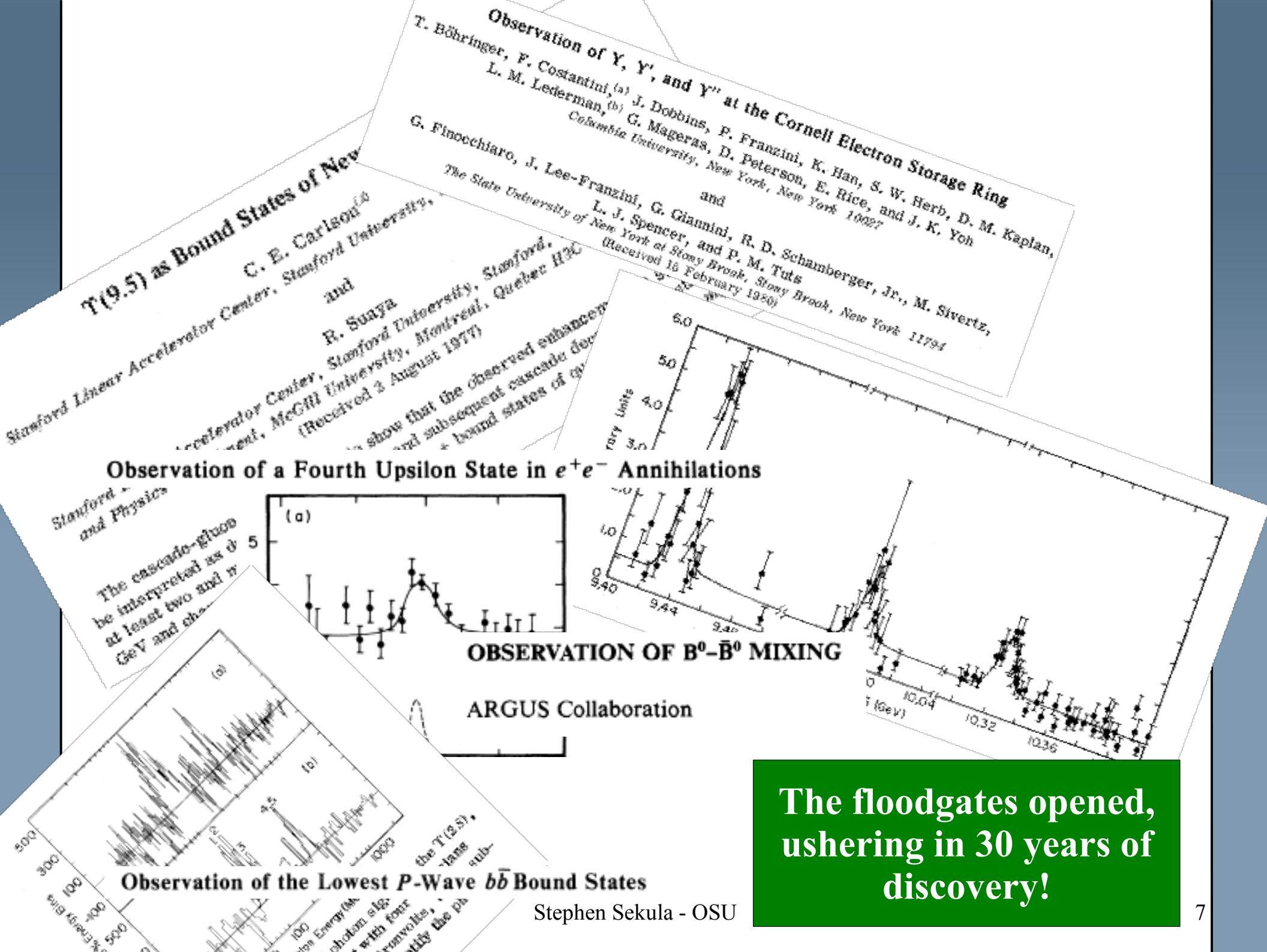
(Received 1 July 1977)

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Dimuon production is studied in 400-GeV proton-nucleus collisions. A strong enhancement is observed at 9.5 GeV mass in a sample of 9000 dimuon events with a mass $m_{\mu^+\mu^-} > 5$ GeV.



The Upsilon is discovered, and identified as the first resonance of a new quark – the bottom quark





2007

$\Upsilon(3S)$

$I^G(J^{PC}) = 0^-(1^- -)$

 $\Upsilon(3S)$ MASS

VALUE (GeV)	DOCUMENT ID	TECN	COMMENT
10.3553 ± 0.0005	¹ ARTAMONOV 00	MD1	$e^+e^- \rightarrow$ hadrons
• • • We do not use the following data for averages, fits, limits, etc. • • •			
10.3553 ± 0.0005	^{2,3} BARU	86B	REDE $e^+e^- \rightarrow$ hadrons
1 Reanalysis of BARU 86B using new electron mass (COHEN 87).			
2 Reanalysis of ARTAMONOV 84.			
3 Superseded by ARTAMONOV 00.			

 $\Upsilon(3S)$ WIDTH

VALUE (keV)	DOCUMENT ID
20.32 ± 1.05 OUR EVALUATION	See the Note on "Width Determinations of the Υ States"

 $\Upsilon(3S)$ DECAY MODES

Mode	Fraction (Γ_i/Γ)	Scale factor/ Confidence level
Γ_1 $\Upsilon(2S)$ anything	(10.6 ± 0.8) %	
Γ_2 $\Upsilon(2S)\pi^+\pi^-$	(2.8 ± 0.6) %	S=2.2
Γ_3 $\Upsilon(2S)\pi^0\pi^0$	(2.00 ± 0.32) %	
Γ_4 $\Upsilon(2S)\gamma\gamma$	(5.0 ± 0.7) %	
Γ_5 $\Upsilon(1S)\pi^+\pi^-$	(4.48 ± 0.21) %	
Γ_6 $\Upsilon(1S)\pi^0\pi^0$	(2.06 ± 0.28) %	
Γ_7 $\Upsilon(1S)\eta$	< 2.2 × 10 ⁻³	CL=90%
Γ_8 $\tau^+\tau^-$	(2.29 ± 0.30) %	
Γ_9 $\mu^+\mu^-$	(2.18 ± 0.21) %	S=2.1
Γ_{10} e^+e^-	seen	
Radiative decays		
Γ_{11} $\gamma\chi_{b2}(2P)$	(13.1 ± 1.6) %	S=3.4
Γ_{12} $\gamma\chi_{b1}(2P)$	(12.6 ± 1.2) %	S=2.4
Γ_{13} $\gamma\chi_{b0}(2P)$	(5.9 ± 0.6) %	S=1.4
Γ_{14} $\gamma\chi_{b0}(1P)$	(3.0 ± 1.1) × 10 ⁻³	
Γ_{15} $\gamma\eta_b(2S)$	< 6.2 × 10 ⁻⁴	CL=90%
Γ_{16} $\gamma\eta_b(1S)$	< 4.3 × 10 ⁻⁴	CL=90%
Γ_{17} $\gamma X \rightarrow \gamma + \geq 4$ prongs	[a] < 2.2 × 10 ⁻⁴	CL=95%
[a] 1.5 GeV < m_X < 5.0 GeV		

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 $\Upsilon(2S)$

$I^G(J^{PC}) = 0^-(1^- -)$

 $\Upsilon(2S)$ MASS

VALUE (GeV)	DOCUMENT ID	TECN	COMMENT
10.02326 ± 0.00031 OUR AVERAGE	¹ ARTAMONOV 00	MD1	$e^+e^- \rightarrow$ hadrons
10.0235 ± 0.0005	² BARBER	84	REDE $e^+e^- \rightarrow$ hadrons
• • • We do not use the following data for averages, fits, limits, etc. • • •			
10.0231 ± 0.0004	^{2,3} BARU	86B	REDE $e^+e^- \rightarrow$ hadrons
1 Reanalysis of BARU 86B using new electron mass (COHEN 87).			
2 Reanalysis of ARTAMONOV 84.			
3 Superseded by ARTAMONOV 00.			

 $\Upsilon(2S)$ WIDTH

VALUE (keV)	DOCUMENT ID
31.98 ± 2.63 OUR EVALUATION	See the Note on "Width Determinations of the Υ States"

 $\Upsilon(2S)$ DECAY MODES

Mode	Fraction (Γ_i/Γ)	Scale factor/ Confidence level
Γ_1 $\Upsilon(1S)\pi^+\pi^-$	(18.8 ± 0.6) %	
Γ_2 $\Upsilon(1S)\pi^0\pi^0$	(9.0 ± 0.8) %	
Γ_3 $\tau^+\tau^-$	(2.00 ± 0.21) %	
Γ_4 $\mu^+\mu^-$	(1.93 ± 0.17) %	S=2.2
Γ_5 e^+e^-	(1.91 ± 0.16) %	
Γ_6 $\Upsilon(1S)\pi^0$	< 1.1 × 10 ⁻³	CL=90%
Γ_7 $\Upsilon(1S)\eta$	< 2 × 10 ⁻³	CL=90%
Γ_8 $J/\psi(1S)$ anything	< 6 × 10 ⁻³	CL=90%
Γ_9 d anything	(3.4 ± 0.6) × 10 ⁻⁵	
Γ_{10} hadrons	(94 ± 11) %	
Radiative decays		
Γ_{11} $\gamma\chi_{b2}(2P)$	(6.9 ± 0.4) %	
Γ_{12} $\gamma\chi_{b1}(2P)$	(7.15 ± 0.35) %	
Γ_{13} $\gamma\chi_{b0}(2P)$	(3.8 ± 0.4) %	
Γ_{14} $\gamma\chi_{b0}(1P)$	< 10 ⁻⁴	CL=90%
Γ_{15} $\gamma\eta_b(2S)$	< 5.9 × 10 ⁻⁴	CL=90%
Γ_{16} $\gamma\eta_b(1S)$	< 5.3 × 10 ⁻⁴	CL=90%
Γ_{17} $\gamma f_J(2200)$	< 2.41 × 10 ⁻⁴	CL=90%
Γ_{18} $\gamma\eta_b(1S)$	< 5.1 × 10 ⁻⁴	CL=90%
Γ_{19} $\gamma X \rightarrow \gamma + \geq 4$ prongs	[a] < 1.95 × 10 ⁻⁴	CL=95%

 $\Upsilon(2S)$ DECAY MODES

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 $\Upsilon(1S)$

$I^G(J^{PC}) = 0^-(1^- -)$

 $\Upsilon(1S)$ MASS

VALUE (MeV)	DOCUMENT ID	TECN	COMMENT
9460.30 ± 0.36 OUR AVERAGE			Error includes scale factor of 3.3
9460.51 ± 0.09 ± 0.05	¹ ARTAMONOV 00	MD1	$e^+e^- \rightarrow$ hadrons
9459.77 ± 0.11 ± 0.07	MACKAY	84	REDE $e^+e^- \rightarrow$ hadrons
• • • We do not use the following data for averages, fits, limits, etc. • • •			
9460.60 ± 0.09 ± 0.05	^{2,3} BARU	92B	REDE $e^+e^- \rightarrow$ hadrons
9460.59 ± 0.12	BARU	86	REDE $e^+e^- \rightarrow$ hadrons
9460.6 ± 0.4	^{3,4} ARTAMONOV 84	REDE	$e^+e^- \rightarrow$ hadrons
1 Reanalysis of BARU 92B and ARTAMONOV 84 using new electron mass (COHEN 87).			
2 Superseding BARU 86.			
3 Value included by ARTAMONOV 84.			
4 Superseded by ARTAMONOV 00.			

 $\Upsilon(1S)$ WIDTH

VALUE (keV)	DOCUMENT ID
54.02 ± 1.25 OUR EVALUATION	See the Note on "Width Determinations of the Υ States"

 $\Upsilon(1S)$ DECAY MODES

Mode	Fraction (Γ_i/Γ)	Confidence level
Γ_1 $\pi^+\pi^-$	(2.60 ± 0.10) %	
Γ_2 e^+e^-	(2.38 ± 0.11) %	
Γ_3 $\mu^+\mu^-$	(2.48 ± 0.05) %	
Hadronic decays		
Γ_4 $\eta'(958)$ anything	(2.94 ± 0.24) %	
Γ_5 $J/\psi(1S)$ anything	(6.5 ± 0.7) × 10 ⁻⁴	
Γ_6 χ_{c0} anything	< 5 × 10 ⁻³	90%
Γ_7 χ_{c1} anything	(2.3 ± 0.7) × 10 ⁻⁴	
Γ_8 χ_{c2} anything	(3.4 ± 1.0) × 10 ⁻⁴	
Γ_9 $\psi(2S)$ anything	(2.7 ± 0.9) × 10 ⁻⁴	
Γ_{10} $\rho\pi$	< 2 × 10 ⁻⁴	90%
Γ_{11} $\pi^+\pi^-$	< 5 × 10 ⁻⁴	90%
Γ_{12} K^+K^-	< 5 × 10 ⁻⁴	90%
Γ_{13} $p\bar{p}$	< 5 × 10 ⁻⁴	90%
Γ_{14} $\pi^0\pi^+ \pi^-$	< 1.84 × 10 ⁻⁵	90%
Γ_{15} $D^+(2010)^+$ anything	(2.86 ± 0.28) × 10 ⁻⁵	
Γ_{16} d anything		
Radiative decays		
Γ_{17} $\gamma\pi^+\pi^-$	(6.3 ± 1.8) × 10 ⁻⁵	
Γ_{18} $\gamma\pi^0\pi^0$	(1.7 ± 0.7) × 10 ⁻⁵	
Γ_{19} $\gamma\pi^0\eta$	< 2.4 × 10 ⁻⁶	90%
Γ_{20} K^+K^- with $2 < m_{K^+K^-} < 3$ GeV	(1.14 ± 0.13) × 10 ⁻⁵	
Γ_{21} $\gamma p\bar{p}$ with $2 < m_{p\bar{p}} < 3$ GeV	< 6 × 10 ⁻⁶	90%
Γ_{22} $\gamma 2\pi^+ 2\pi^-$	(7.0 ± 1.5) × 10 ⁻⁴	
Γ_{23} $\gamma 3\pi^+ 3\pi^-$	(5.4 ± 2.0) × 10 ⁻⁴	
Γ_{24} $\gamma 4\pi^+ 4\pi^-$	(7.4 ± 3.5) × 10 ⁻⁴	
Γ_{25} $\gamma\pi^+\pi^- K^+K^-$	(2.9 ± 0.9) × 10 ⁻⁴	
Γ_{26} $\gamma 2\pi^+ 2\pi^-$	(2.5 ± 0.9) × 10 ⁻⁴	
Γ_{27} $\gamma 3\pi^+ 3\pi^-$	(2.5 ± 1.2) × 10 ⁻⁴	
Γ_{28} $\gamma 2\pi^+ 2\pi^- K^+K^-$	(2.4 ± 1.2) × 10 ⁻⁴	
Γ_{29} $\gamma\pi^+\pi^- \rho\bar{p}$	(1.5 ± 0.6) × 10 ⁻⁴	
Γ_{30} $\gamma 2\pi^+ 2\pi^- \rho\bar{p}$	(4 ± 6) × 10 ⁻⁵	
Γ_{31} $\gamma 2K^+ 2K^-$	(2.0 ± 2.0) × 10 ⁻⁵	
Γ_{32} $\gamma\eta'(958)$	< 1.9 × 10 ⁻⁶	90%
Γ_{33} $\gamma\eta$	< 1.0 × 10 ⁻⁶	90%
Γ_{34} $\gamma\eta(980)$	< 3 × 10 ⁻⁵	90%
Γ_{35} $\gamma f_0(1525)$	(3.7 ± 1.2) × 10 ⁻⁵	
Γ_{36} $\gamma f_0(1270)$	(1.01 ± 0.09) × 10 ⁻⁴	
Γ_{37} $\gamma f_0(1500)$	< 8.2 × 10 ⁻⁵	90%
Γ_{38} $\gamma f_0(1710)$	< 1.5 × 10 ⁻⁵	90%
Γ_{39} $\gamma f_0(1710) \rightarrow \gamma K^+K^-$	< 2.6 × 10 ⁻⁴	90%
Γ_{40} $\gamma f_0(1710) \rightarrow \gamma\pi^0\pi^0$	< 7 × 10 ⁻⁶	90%
Γ_{41} $\gamma f_0(1710) \rightarrow \gamma\eta\eta$	< 1.4 × 10 ⁻⁶	90%
Γ_{42} $\gamma f_0(2050)$	< 1.8 × 10 ⁻⁶	90%
Γ_{43} $\gamma f_0(2200) \rightarrow \gamma K^+K^-$	< 2.2 × 10 ⁻⁴	90%
Γ_{44} $\gamma f_0(2200) \rightarrow \gamma\pi^+ \pi^-$	< 8 × 10 ⁻⁷	90%
Γ_{45} $\gamma f_0(2200) \rightarrow \gamma K^+K^-$	< 6 × 10 ⁻⁷	90%
Γ_{46} $\gamma f_0(2200) \rightarrow \gamma\pi^+ \pi^-$	< 1.1 × 10 ⁻⁶	90%
Γ_{47} $\gamma f_0(2200) \rightarrow \gamma p\bar{p}$	< 1.3 × 10 ⁻⁶	90%
Γ_{48} $\gamma\eta(2225) \rightarrow \gamma\phi\phi$	< 3 × 10 ⁻³	90%
Γ_{49} $\gamma X \rightarrow \gamma + \geq 4$ prongs	[a] < 3 × 10 ⁻⁵	90%
Γ_{50} $\gamma X \rightarrow \gamma + \geq 4$ prongs	[b] < 1 × 10 ⁻³	90%
Γ_{51} invisible	[c] < 1.78 × 10 ⁻⁴	90%
Other decays		
Γ_{52} invisible	< 2.5 × 10 ⁻³	90%

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Gunion, C. Amsler et al. (Particle Data Group), PL Short, 1 (2006) (URL: http://pdg.lbl.gov)		Gunion, C. Amsler et al. (Particle Data Group), PL Short, 1 (2006) (URL: http://pdg.lbl.gov)		Gunion, C. Amsler et al. (Particle Data Group), PL Short, 1 (2006) (URL: http://pdg.lbl.gov)		Gunion, C. Amsler et al. (Particle Data Group), PL Short, 1 (2006) (URL: http://pdg.lbl.gov)	
For inclusive branching fractions, e.g., $B \rightarrow D^+ \pi^-$, it usually are multiplications, not branching fractions. They can be shown here.		τX	$\bar{D}_{s0}^*(2460)^+ \pi^+$	ℓ	$D_{s0}(2457)^+ D^-$	$(-2.81 \pm 0.24) \times 10^{-4}$	
		$c X$	$B(D_s(2460)^- \rightarrow D^0 \pi^-)$		$\bar{D}_{s0}^*(2457)^0 \bar{\rho}^0$	$(-2.2 \pm 0.9) \times 10^{-4}$	
		$\tau c X$	$D_s(2460)^- \pi^0$		$D_{s0}(2457)^0 \pi^-$	$(-1.2 \pm 0.4) \times 10^{-4}$	
Mode	Fraction (T_f/T)	$D, D^*, \text{ or } D_s \pi$		$D_{s0}(2457)^0 \pi^0$	$(-2.02 \pm 0.35) \times 10^{-4}$	S=1.6	
$\tau_1 \ell^+ \nu_\ell \text{ anything}$	[4] $(10.3 \pm 0.28) \%$	$\tau_{31} D^- \pi^+$	$D_{s0}(2457)^0 \pi^+$	$(-1.26 \pm 0.21) \times 10^{-4}$	S=1.1		
$e^+ \nu_e X_0$	$(10.2 \pm 0.4) \%$	$\tau_{32} D^- \pi^0$	$D_{s0}(2457)^0 \pi^0$	$(-2.59 \pm 0.36) \times 10^{-4}$			
$\tau_3 D^+ \ell^- \nu_\ell \text{ anything}$	$(9.8 \pm 0.9) \%$	$\tau_{33} D^- K^0 \pi^+$	$D_{s0}(2457)^0 K^0$	$< 1.16 \times 10^{-5}$	CL=90%		
$\tau_4 D^- \pi^+ \nu_\ell$	$(2.37 \pm 0.12) \%$	$\tau_{34} D^- K^{*0}$	$D_{s0}(2457)^0 K^{*0}$	$< 1.9 \times 10^{-5}$	CL=90%		
$\tau_5 D^- \pi^+ \nu_\tau$	$(1.0 \pm 0.4) \%$	$\tau_{35} D^- \pi^+ \pi^0$	$D_{s0}(2457)^0 \pi^+ \pi^0$	$< 1.1 \times 10^{-5}$	CL=90%		
$\tau_6 D^*(2010)^- \ell^+ \nu_\ell$	[4] $(8.16 \pm 0.11) \%$	$\tau_{36} D^- K^+ \pi^0$	$D_{s0}(2457)^0 K^+ \pi^0$	$< 2.1 \times 10^{-5}$	CL=90%		
$\tau_7 D^*(2010)^- \ell^+ \nu_\tau$	$(1.8 \pm 0.5) \%$	$\tau_{37} D^- K^+ \pi^0$	$D_{s0}(2457)^0 K^+ \pi^0$	$< 1.8 \pm 0.1) \times 10^{-4}$	S=1.8		
$\tau_8 D^2 \pi^- \ell^+ \nu_\ell$	$(4.3 \pm 0.6) \times 1$	$\tau_{38} D^2 \pi^- \pi^+$	$D_{s0}(2457)^0 D^2 \pi^+$	$< 1.21 \pm 0.18) \times 10^{-4}$	S=1.8		
$\tau_9 D_s(2460)^- \ell^+ \nu_\ell \times$	$(2.0 \pm 0.9) \times 1$	$\tau_{39} D^2 \pi^- \pi^0$	$D_{s0}(2457)^0 D^2 \pi^0$	$< 6.2 \pm 2.2) \times 10^{-4}$			
$B(D_s^+ \rightarrow \overline{D}^0 \pi^+)$		$\tau_{40} D^2 \pi^+ \pi^-$	$D_{s0}(2457)^0 D^2 \pi^-$	$< 3.6 \pm 1.2) \times 10^{-5}$			
$\tau_{10} D_s(2460)^- \ell^+ \nu_\ell \times$	$(-2.2 \pm 0.6) \times 1$	$\tau_{41} D^2 \pi^+ \pi^0$	$D_{s0}(2457)^0 D^2 \pi^0$	$< 6.9 \times 10^{-5}$	CL=90%		
$B(D_s^+ \rightarrow \overline{D}^0 \pi^+)$		$\tau_{42} (D^+ \pi^+ \pi^+ \pi^-) \text{ nonresonant}$	$D_{s0}(2457)^0 (D^+ \pi^+ \pi^-)$	$< 4.0 \times 10^{-5}$	CL=90%		
$\tau_{11} D_s(2460)^- \ell^+ \nu_\ell (\alpha \geq 1)$	$(-2.4 \pm 0.8) \%$	$\tau_{43} D^2 \pi^+ \rho^0$	$D_{s0}(2457)^0 D^2 \rho^0$	$< 2.7 \pm 0.5) \times 10^{-5}$			
$\tau_{12} \overline{D}^0 \pi^- \ell^+ \nu_\ell$	$(4.9 \pm 0.8) \times 1$	$\tau_{44} D^*(2010)^- \pi^+$	$D_{s0}(2457)^0 K^+ \pi^-$	$< 8.2 \pm 0.9) \times 10^{-4}$			
$\tau_{13} D_s(2460)^- \ell^+ \nu_\lambda \times$	$(5.4 \pm 2.1) \times 1$	$\tau_{45} D^*(2010)^- \rho^0$	$D_{s0}(2457)^0 D^* \rho^0$	$< 2.7 \pm 0.8) \times 10^{-4}$	S=1.8		
$\tau_{14} D_s(2460)^- \ell^+ \nu_\lambda \times$	$< 8.0 \times 1$	$\tau_{46} D^*(2010)^- K^+$	$D_{s0}(2457)^0 D^* K^+$	$< 6.1 \pm 1.9) \times 10^{-4}$	S=1.6		
$B(D_s^+ \rightarrow \overline{D}^0 \pi^+)$		$\tau_{47} D^*(2010)^- K^0$	$D_{s0}(2457)^0 D^* K^0$	$< 9 \times 10^{-5}$	CL=90%		
$\tau_{15} D_s(2460)^- \ell^+ \nu_\lambda \times$	$< 3.0 \times 1$	$\tau_{48} D^*(2010)^- K^{*0}$	$D_{s0}(2457)^0 D^* K^{*0}$	$< 1.7 \pm 0.4) \times 10^{-4}$	S=1.8		
$B(D_s^+ \rightarrow \overline{D}^0 \pi^+)$		$\tau_{49} D^*(2010)^- \pi^0$	$D_{s0}(2457)^0 D^* \pi^0$	$< 5.1 \times 10^{-4}$	CL=90%		
$\tau_{50} \rho^- \ell^+ \nu_\ell$	$(-2.47 \pm 0.31) \times 1$	$\tau_{51} D^*(2010)^- K^+ \overline{K}^0$	$D_{s0}(2457)^0 D^* K^+ \overline{K}^0$	$< 1.8 \pm 0.1) \times 10^{-4}$	S=1.8		
$\tau_{51} \rho^- \ell^+ \nu_\lambda$	$(-2.36 \pm 0.08) \times 1$	$\tau_{52} D^*(2010)^- \pi^+ \pi^-$	$D_{s0}(2457)^0 \pi^+ \pi^-$	$< 4.4 \pm 1.0) \times 10^{-4}$			
$\tau_{53} (D^*(2010)^- \pi^+ \pi^- \pi^-)$		$\tau_{54} D^*(2010)^- \pi^0 \pi^0$	$D_{s0}(2457)^0 \pi^0 \pi^0$	$< 3.8 \times 10^{-4}$	CL=90%		
Inclusive modes		$\tau_{55} D^*(2010)^- \pi^+ \pi^+ \pi^-$	$D_{s0}(2457)^0 \pi^+ \pi^+ \pi^-$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{56} D_s(2460)^- \pi^+ \nu_\ell \times$	$< 8.0 \times 1$	$\tau_{57} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$B(D_s^+ \rightarrow \overline{D}^0 \pi^+)$		$\tau_{58} D_s(2460)^- \pi^0 \nu_\ell \times$	$D_{s0}(2457)^0 \pi^0 \nu_\ell \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{59} D_s(2460)^- \pi^0 \nu_\lambda \times$	$< 3.0 \times 1$	$\tau_{60} D_s(2460)^- \pi^- \nu_\ell \times$	$D_{s0}(2457)^0 \pi^- \nu_\ell \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{61} D_s(2460)^- \pi^- \nu_\lambda \times$		$\tau_{62} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{63} D_s(2460)^- \pi^0 \nu_\lambda \times$	$< 2.8 \times 1$	$\tau_{64} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{65} D_s(2460)^- \pi^+ \nu_\lambda \times$	$< 3.3 \times 1$	$\tau_{66} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{67} \overline{D}_s^- \pi^+ \times$	$(5.0 \pm 2.1) \%$	$\tau_{68} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{69} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(10.3 \pm 1.8) \%$	$\tau_{70} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{71} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{72} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{73} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{74} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{75} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{76} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{77} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{78} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{79} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{80} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{81} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{82} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{83} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{84} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{85} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{86} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{87} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{88} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{89} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{90} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{91} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{92} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{93} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{94} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{95} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{96} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{97} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{98} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{99} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{100} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{101} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{102} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{103} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{104} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{105} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{106} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{107} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{108} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{109} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{110} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{111} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{112} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{113} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{114} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{115} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{116} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{117} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{118} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{119} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{120} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{121} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{122} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{123} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{124} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{125} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{126} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{127} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{128} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{129} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{130} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{131} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{132} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{133} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{134} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{135} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{136} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{137} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{138} D_s(2460)^- \pi^0 \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^0 \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{139} D_s(2460)^- \pi^+ \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{140} D_s(2460)^- \pi^- \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^- \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{141} D_s(2460)^- \pi^0 \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	$\tau_{142} D_s(2460)^- \pi^+ \nu_\lambda \times$	$D_{s0}(2457)^0 \pi^+ \nu_\lambda \times$	$< 1.2 \times 10^{-4}$	CL=90%		
$\tau_{143} D_s(2460)^- \pi^- \nu_\lambda \times$	$(-2.4 \pm 0.8) \%$	<math					

The case for BaBar taking data at one of the narrow Upsilon resonances built over time, and involved the whole collaboration.

Here are just a few snapshots . . .

June Collaboration Meeting, 2007

Ideas for searching for a low-mass Higgs [\(pdf\)](#) [\(ppt\)](#) [\(video\)](#)

October, 2007



Higgs and Exotics 2007 Workshop
Workshop on Higgs and other Exotic particles

Babar Home

Higgs Workshop, Monday October 29, 2007

December Collaboration Meeting, 2007

Run Strategy

17:30-17:40

Upsilon (3S) SM Physics [\(pdf\)](#) [\(ppt\)](#) [\(video\)](#)

17:40-17:50

Upsilon (3S) non-SM Physics [\(pdf\)](#) [\(ppt\)](#) [\(video\)](#)

17:50-18:10

Upsilon (5S) Physics [\(pdf\)](#) [\(ppt\)](#) [\(video\)](#)

18:10-18:25

Off-resonance data [\(pdf\)](#) [\(ppt\)](#) [\(video\)](#)

After December 17, 2007:

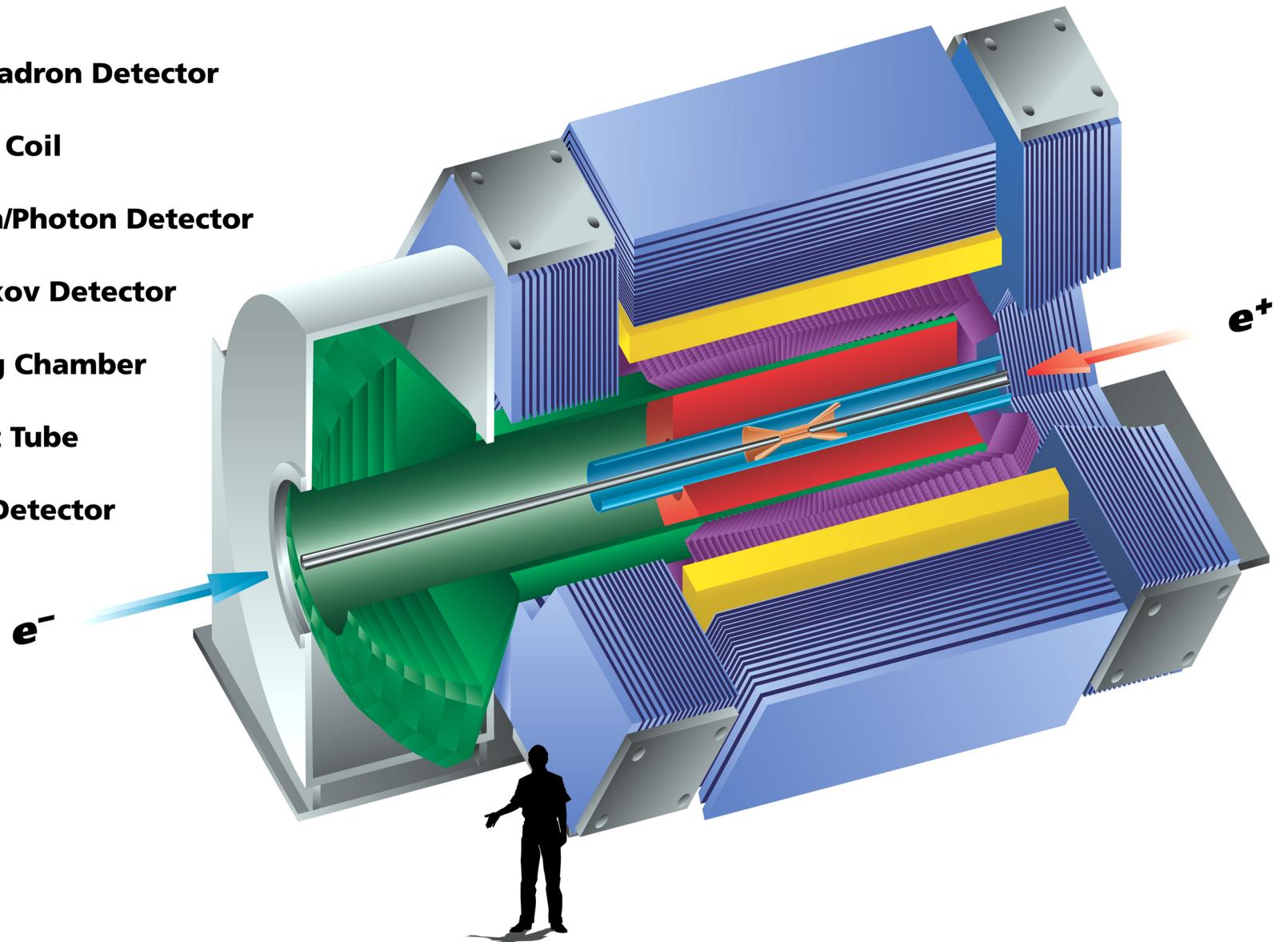
The Physics Case for Running the B-factory at
the $\Upsilon(3S)$ Resonance

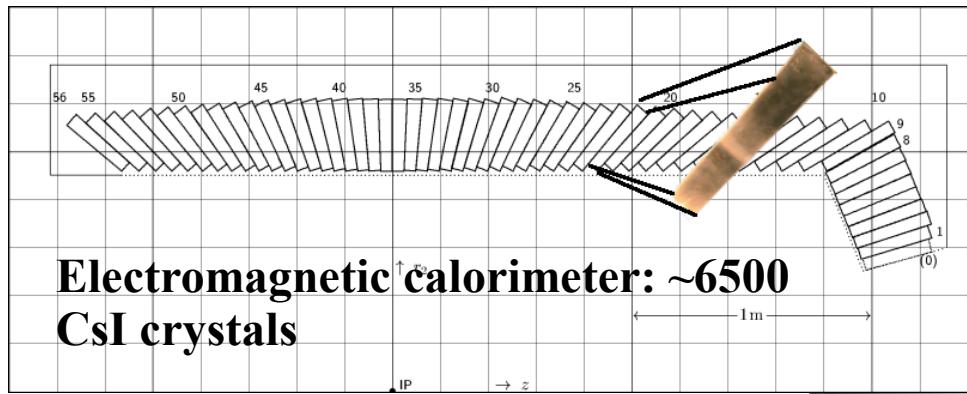
The BaBar/PEP-II b-Factory



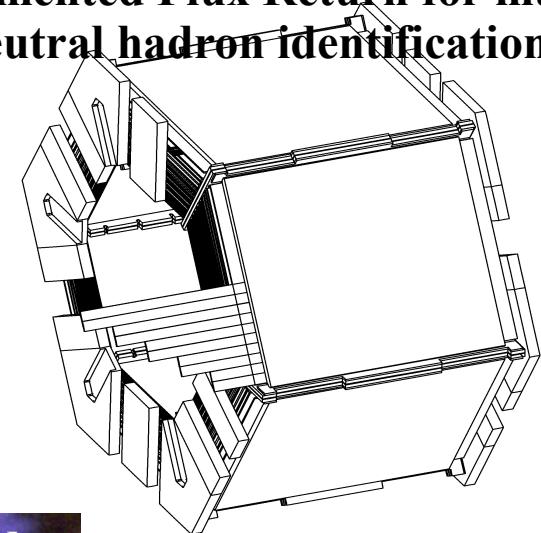
BABAR Detector

- █ Muon/Hadron Detector
- █ Magnet Coil
- █ Electron/Photon Detector
- █ Cherenkov Detector
- █ Tracking Chamber
- █ Support Tube
- █ Vertex Detector

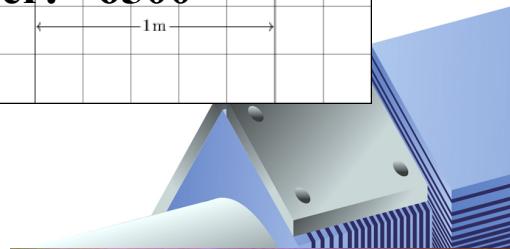




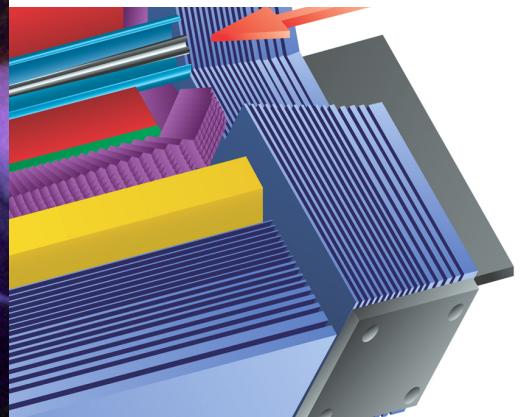
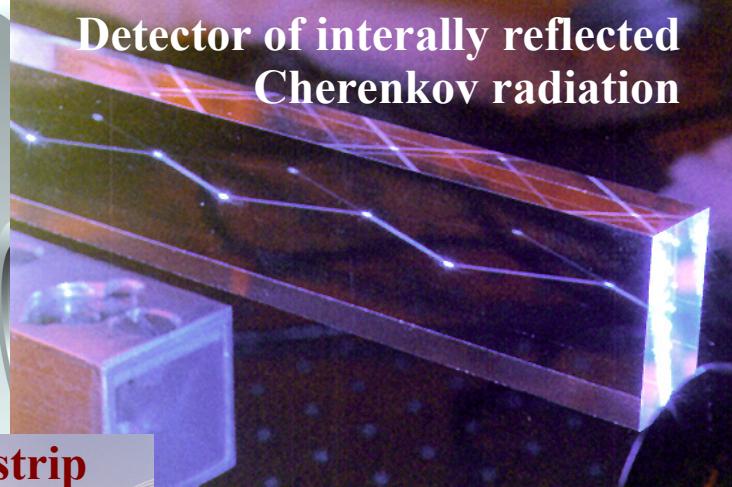
Instrumented Flux Return for muon and neutral hadron identification



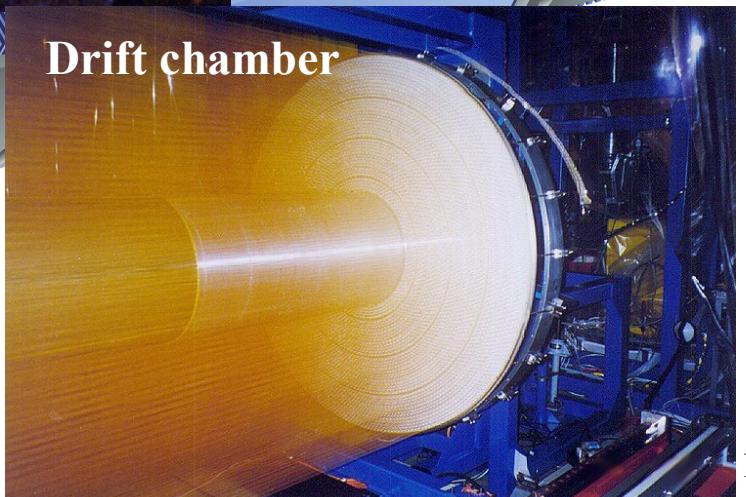
- █ Electron/Photon Detector
- █ Cherenkov Detector
- █ Tracking Chamber
- █ Support Tube
- █ Vertex Detector

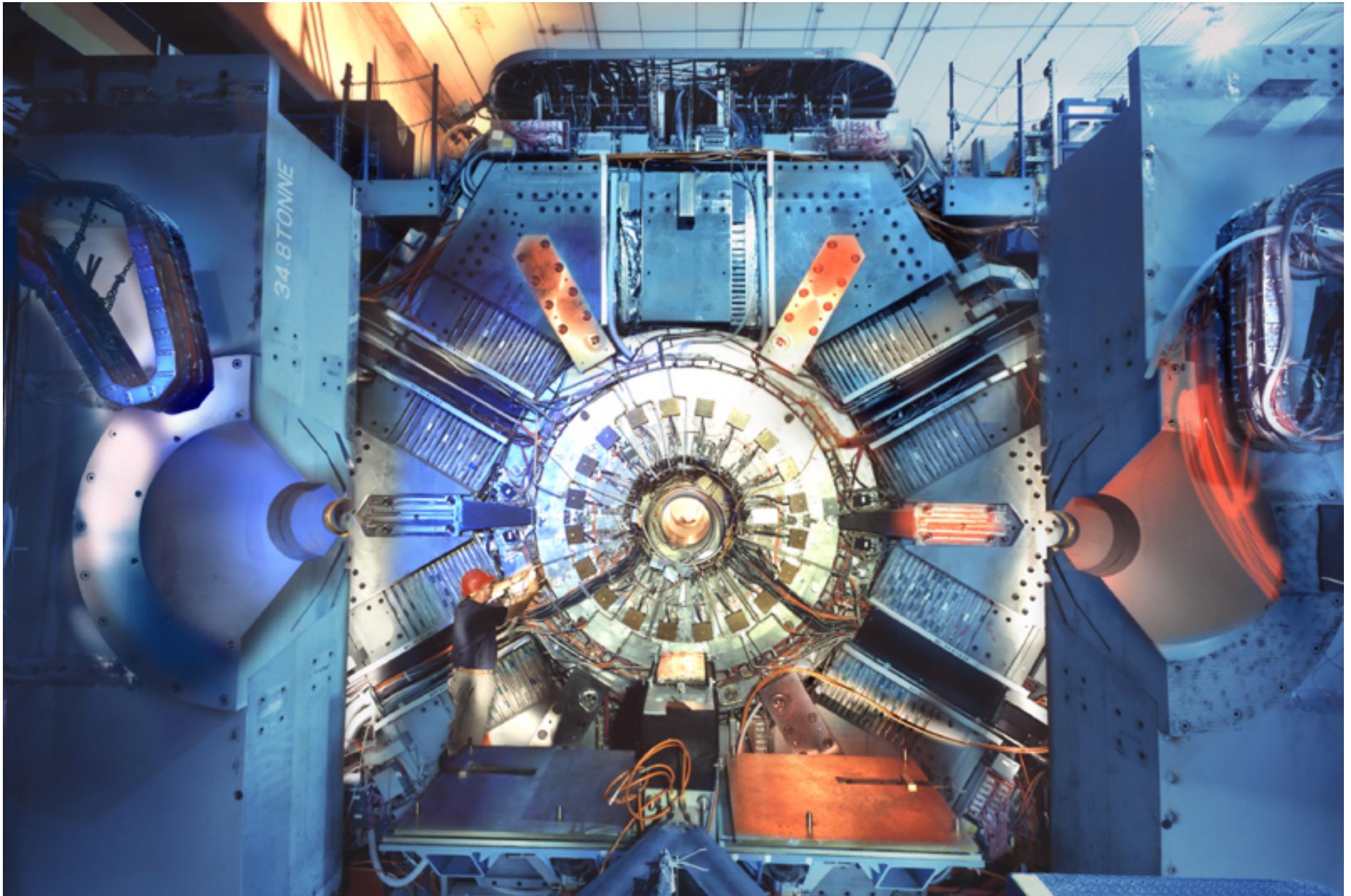


Detector of internally reflected Cherenkov radiation

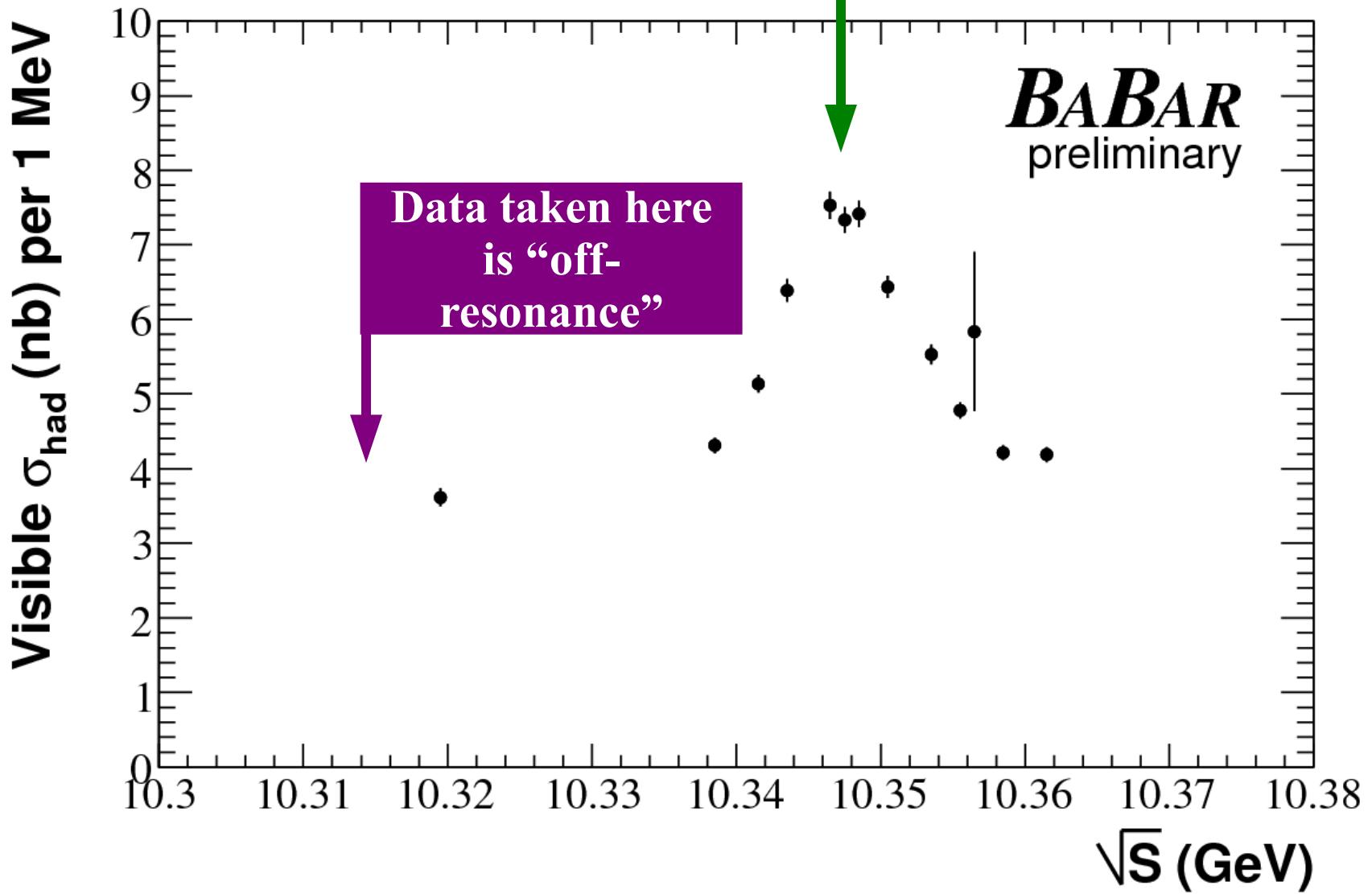


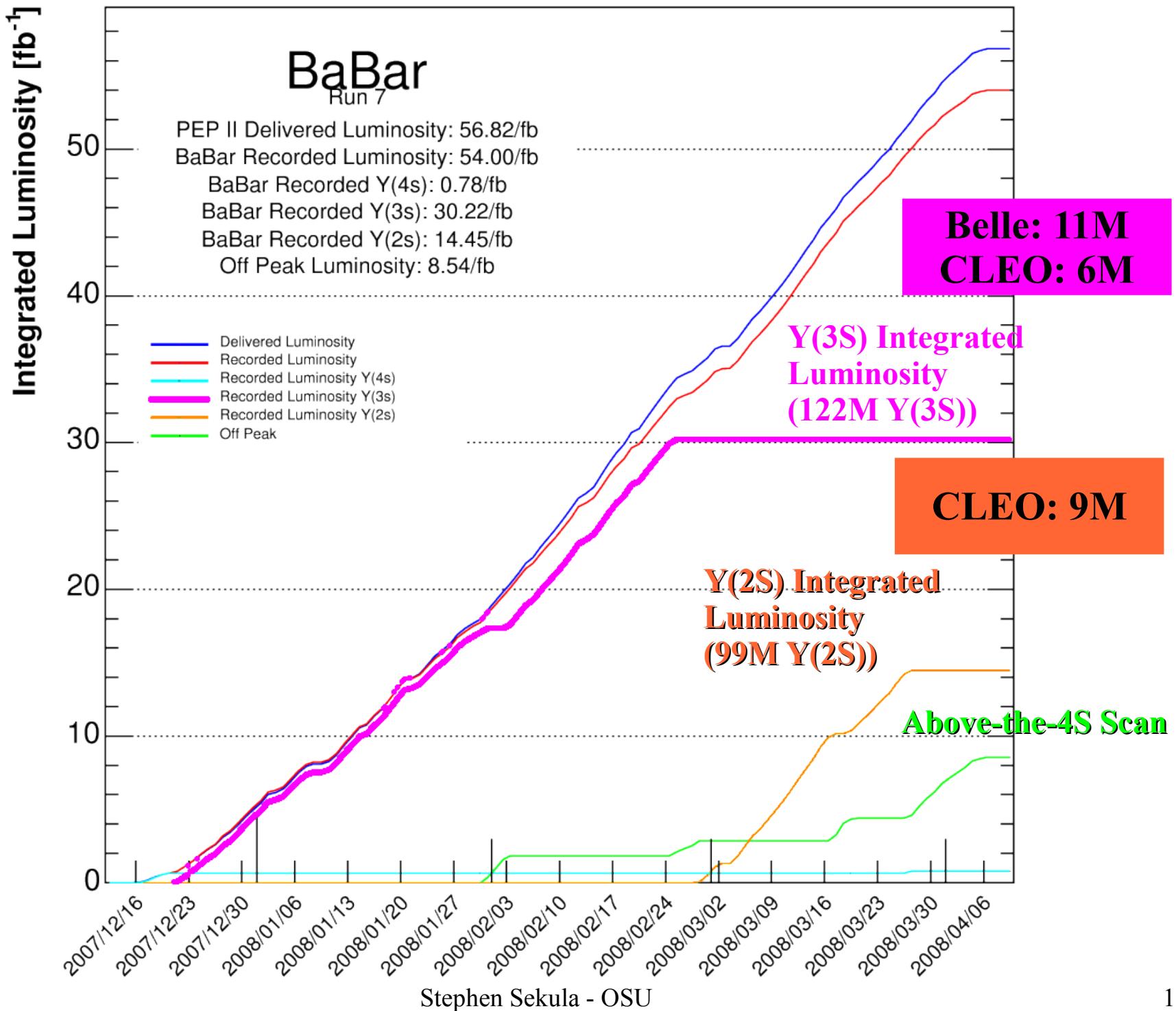
Stephen Sekula - OSU





Scan Data from Dec. 22, 2007





A matter of QCD: The search for the η_b

Quarks

<i>u</i>	<i>c</i>	<i>t</i>
up	charm	top

<i>d</i>	<i>s</i>	<i>b</i>
down	strange	bottom

<i>e</i>	<i>μ</i>	<i>τ</i>
electron	muon	tau

<i>ν_e</i>	<i>ν_μ</i>	<i>ν_τ</i>
electron neutrino	muon neutrino	tau neutrino

Leptons

Forces

<i>Z</i>	<i>γ</i>
Z boson	photon

<i>W</i>	<i>g</i>
W boson	gluon

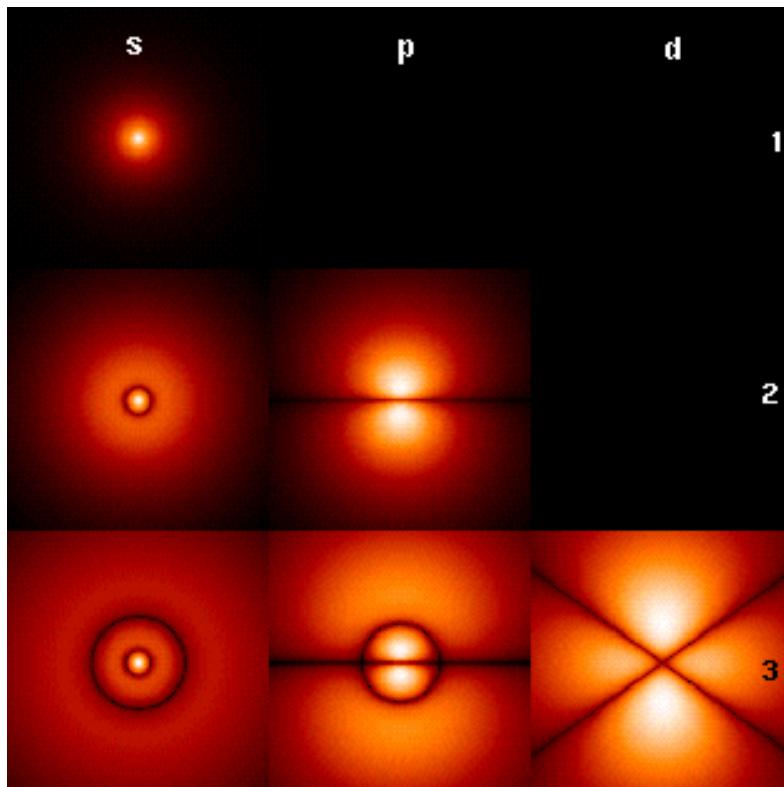
Visible Matter

Remember your Quantum Mechanics

What are the allowed states of a pair of spin-1/2 particles?

SPIN:   $S_{b\bar{b}}=0, 1$

ORBITAL: L=0, 1, 2, ... (S, P, D, ...)

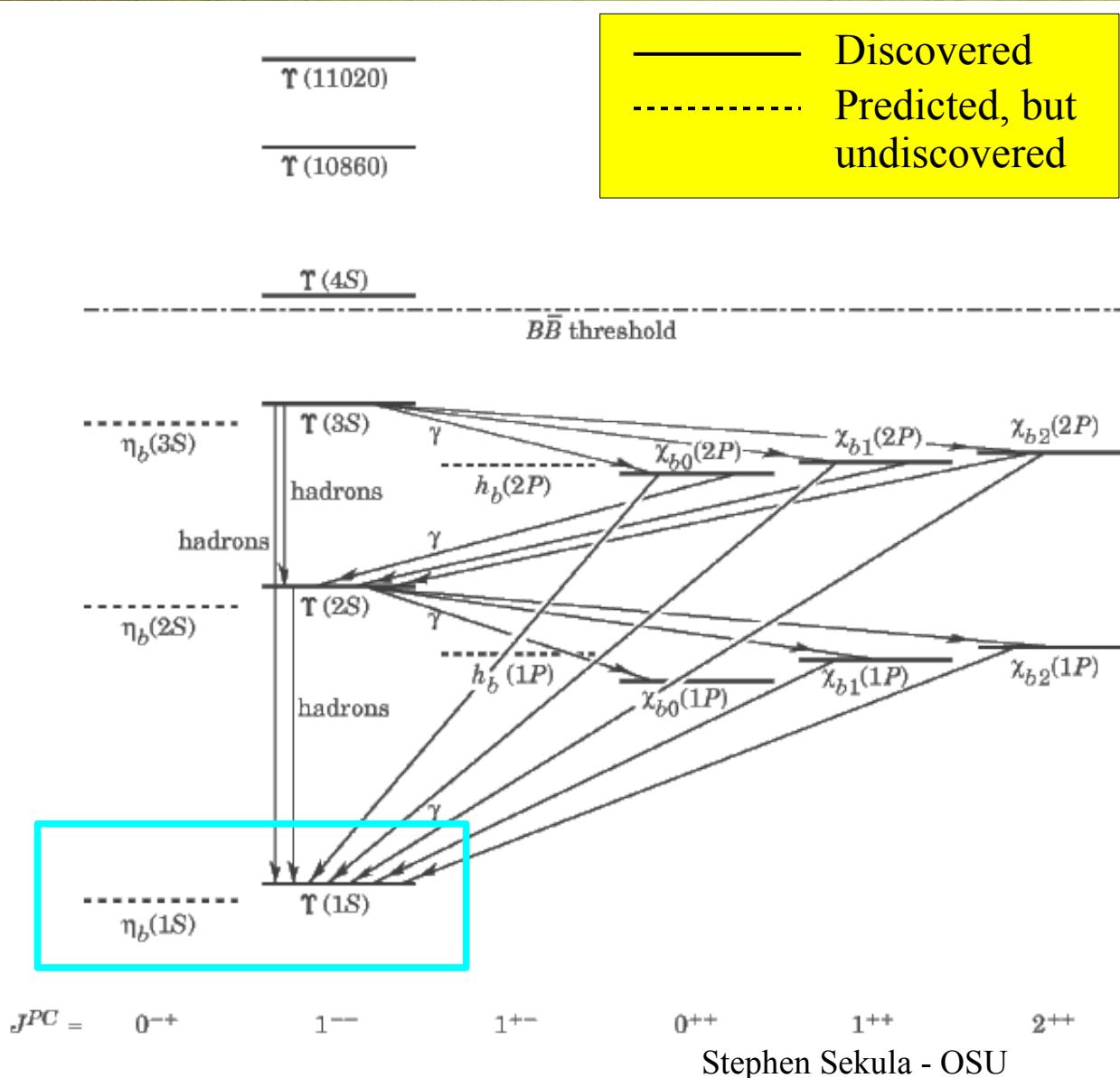


TOTAL ANGULAR MOMENTUM (J):
 $|L-S| < J < L+S$

THE FIRST FEW STATES:

L	S	J	State
0	0	0	$\eta_b(1S, 2S, \dots)$
0	1	1	$\Upsilon(1S, 2S, \dots)$
1	0	1	$h_b(1P, 2P, \dots)$
1	1	0,1,2	$\chi_{bJ}(1P, 2P, \dots)$

Spectroscopy: Find the bottomonium ground state



QCD is assumed to be the dominant factor in defining the spectrum of states. Predictions proceed from this . . .

Hyperfine splitting predictions ($1^3S_1 - 1^1S_0$)

- pNRQCD: **(39-44) MeV** ($\sim 25\%$ uncertainty)
- Potential models: **(46-87) MeV**
- Lattice QCD: **(40-71) MeV (10-25% uncertainty)**

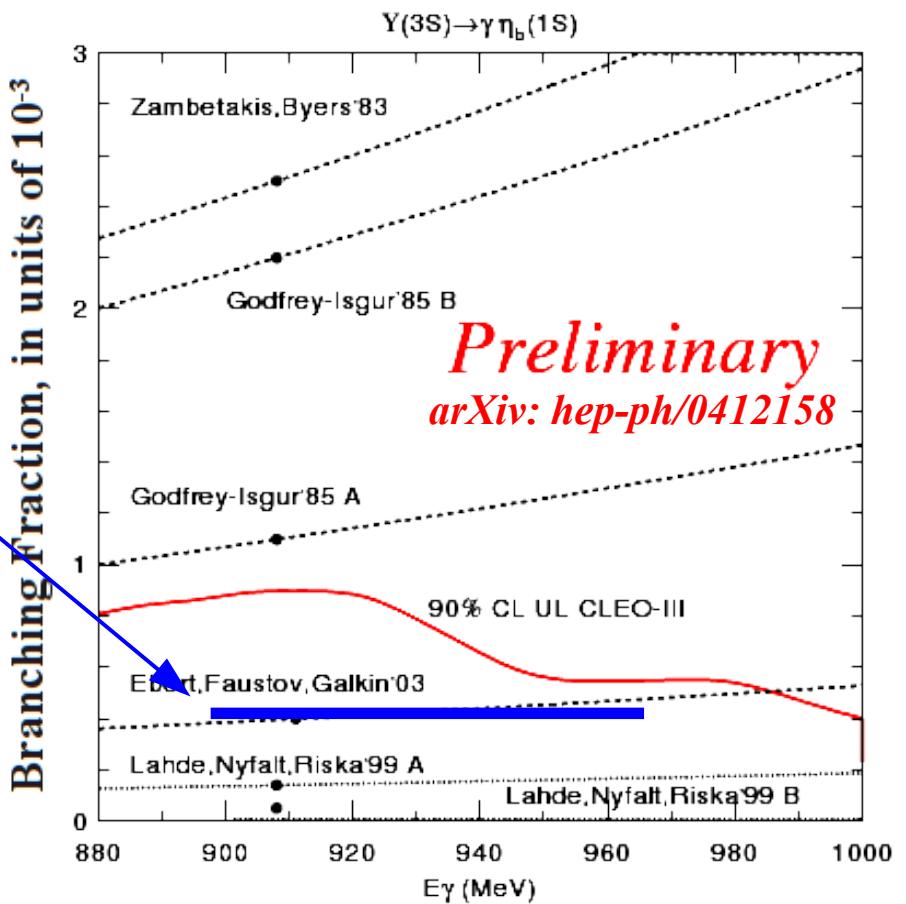
What were the best existing experimental constraints?

$$Y(nS) \rightarrow \gamma \eta_b$$

Published CLEO limits
PRL 94 032001 (2005)

$$e^+ e^- \rightarrow e^+ e^- \gamma^* \gamma^* (\rightarrow \eta_b)$$

Expt	final state	$\Gamma_{\gamma\gamma} \times \mathcal{B}$ (keV)
ALEPH	4 charged	< 0.048
	6 charged	< 0.132
L3	$K^+ K^- \pi^0$	< 2.83
	4 charged	< 0.21
	4 charged π^0	< 0.50
	6 charged	< 0.33
	6 charged π^0	< 5.50
	$\pi^+ \pi^- \eta'$	< 3.00
DELPHI	4 charged	< 0.093
	6 charged	< 0.270
	8 charged	< 0.780



*30 years after the discovery
of the Upsilon, the ground
state of bottomonium had
eluded detection*

Analysis Strategy

Blind Analysis

We never look at the signal region in the final data set until the analysis method is finalized.

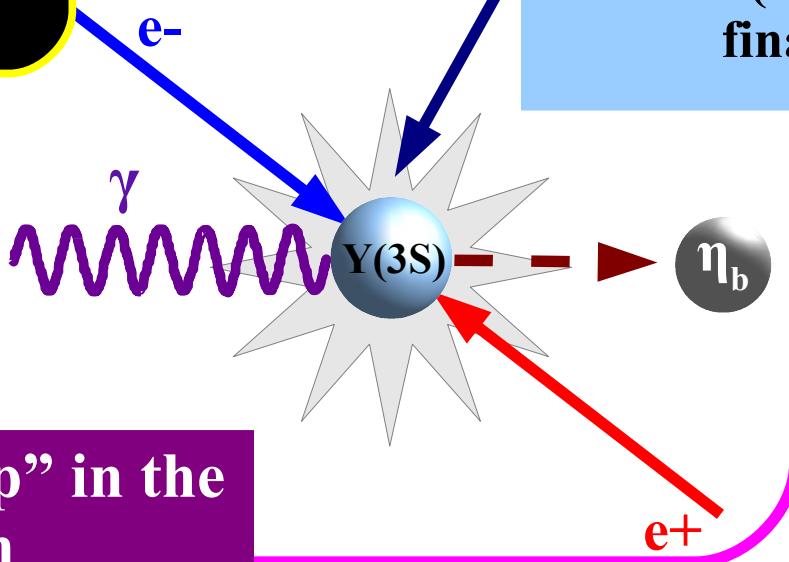
$$E_{\gamma}^* = \frac{m_{Y(3S)}^2 - m_{\eta_b}^2}{2 m_{Y(3S)}^2}$$

Search for a “bump” in the photon spectrum

use maximum likelihood fit,
including backgrounds
and a possible signal

Full dataset: 122M Y(3S) mesons

use a small sample (9%) for tuning the selection
use $(109 \pm 1) \times 10^6$ Y(3S) for final result



Monte Carlo Simulations
used for modeling signal and specific backgrounds
tune selection criteria

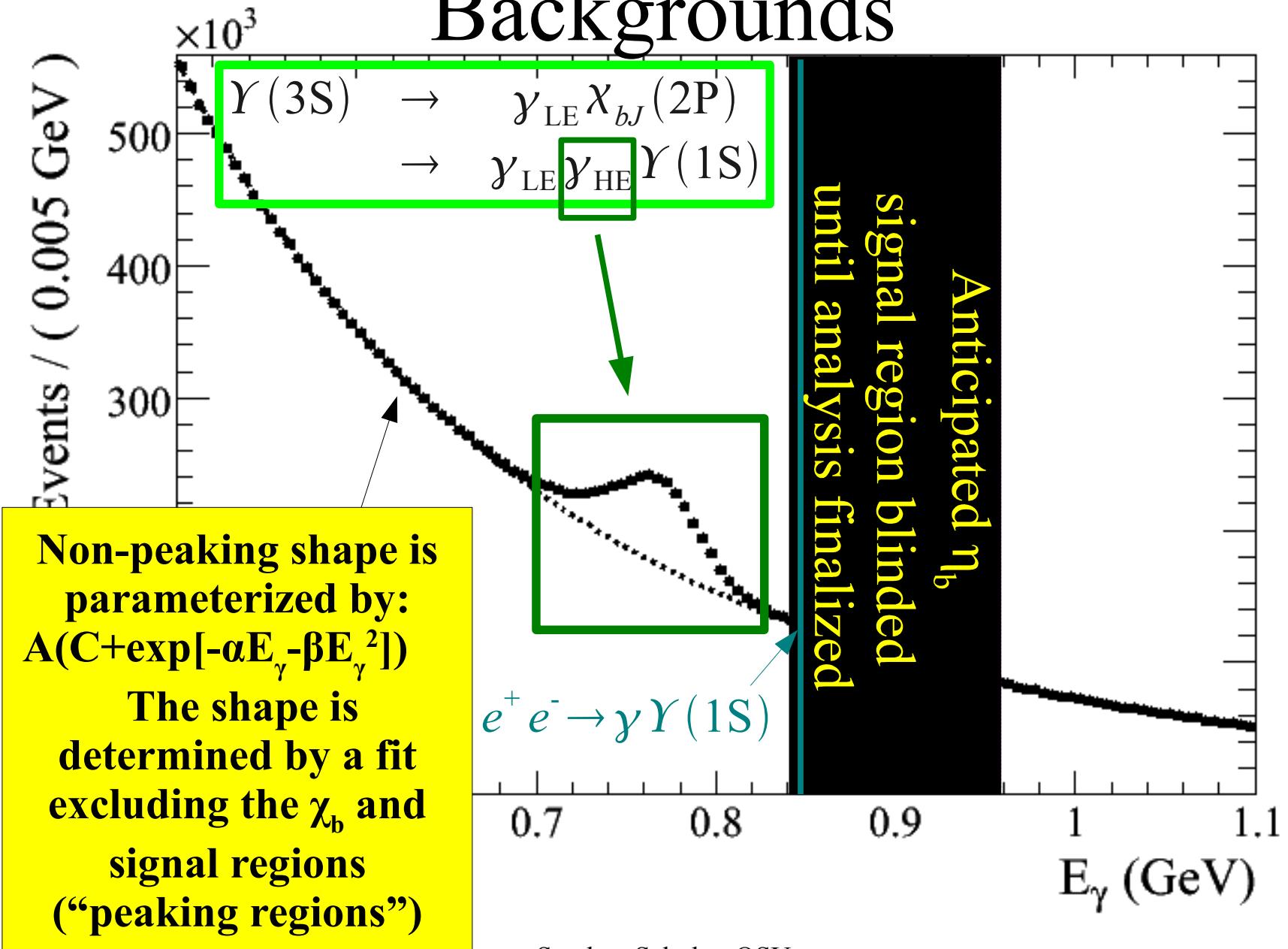
An illustrative signal simulation event . . .

Signal photon required to be reconstructed with high quality, be well within the calorimeter acceptance, and be inconsistent with originating from a π^0

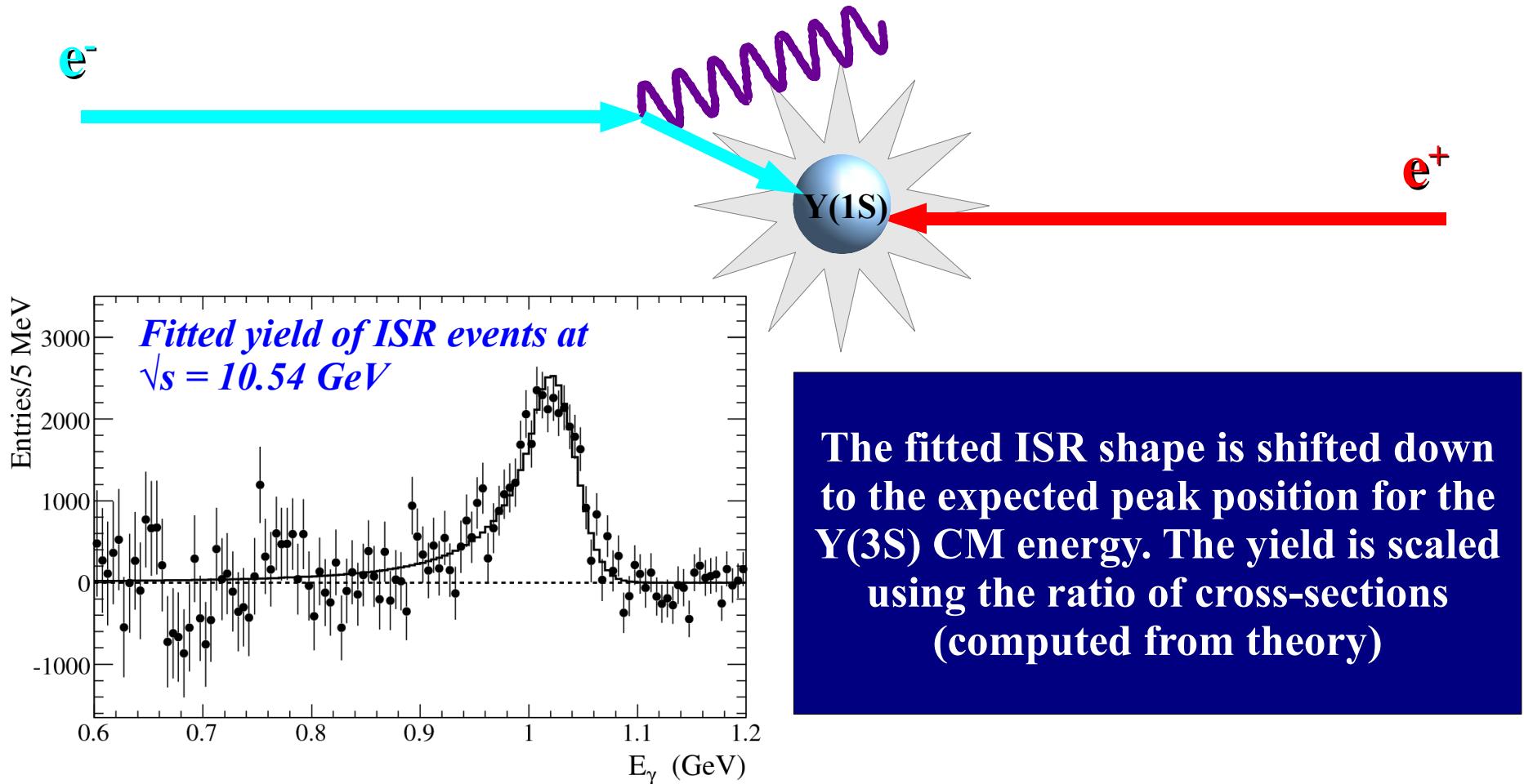
η_b expected to decay into many hadrons (through two gluons), and have uniform distribution of final state particles

Signal Efficiency:
37%

The Single Photon Challenge: Backgrounds



$e^+e^- \rightarrow \gamma_{\text{ISR}} Y(1S)$: Expectation

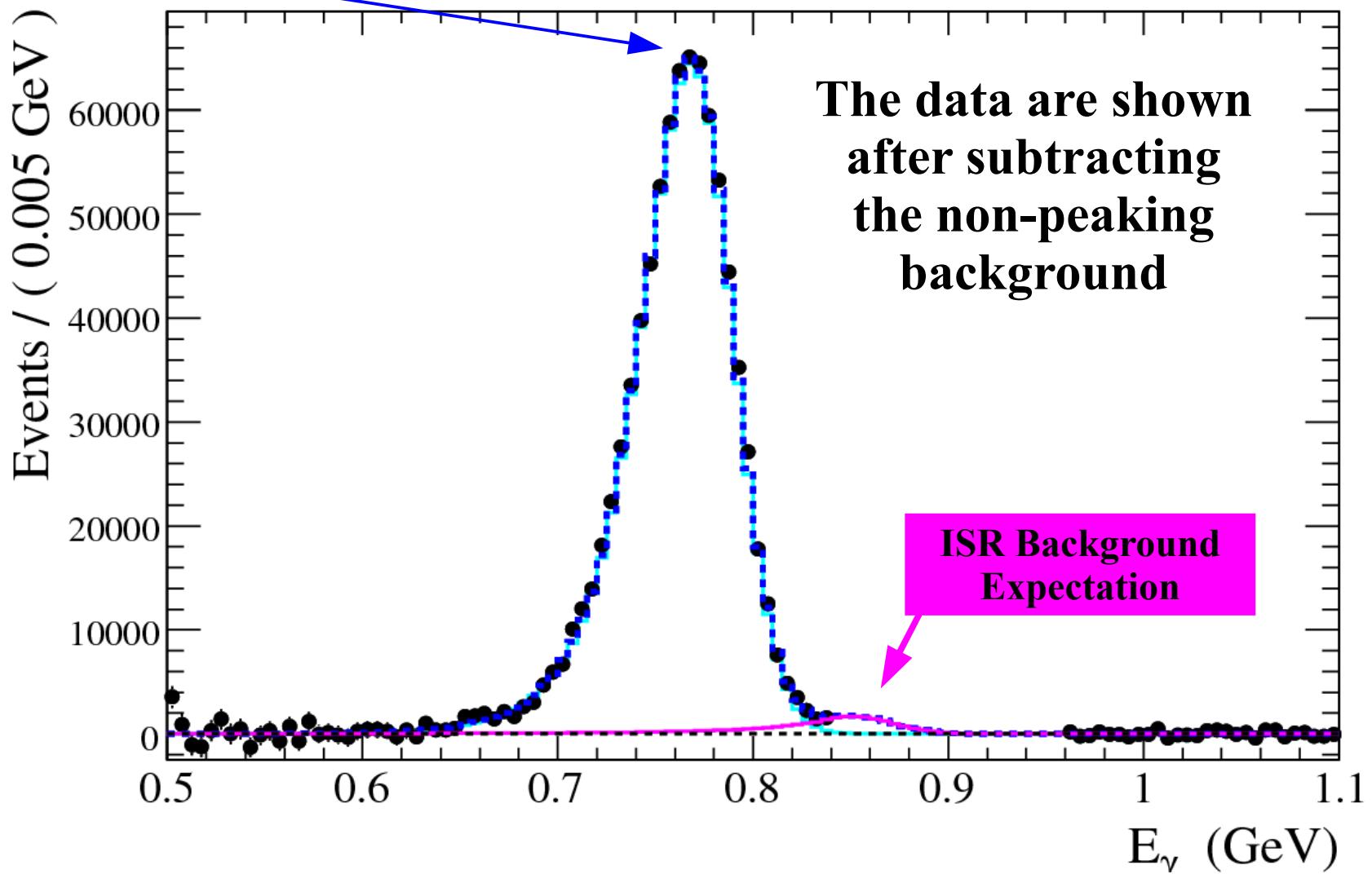


$\sqrt{s} = 10.54 \text{ GeV} \rightarrow \sqrt{s} = 10.3552 \text{ GeV}: 25153 \pm 1677$

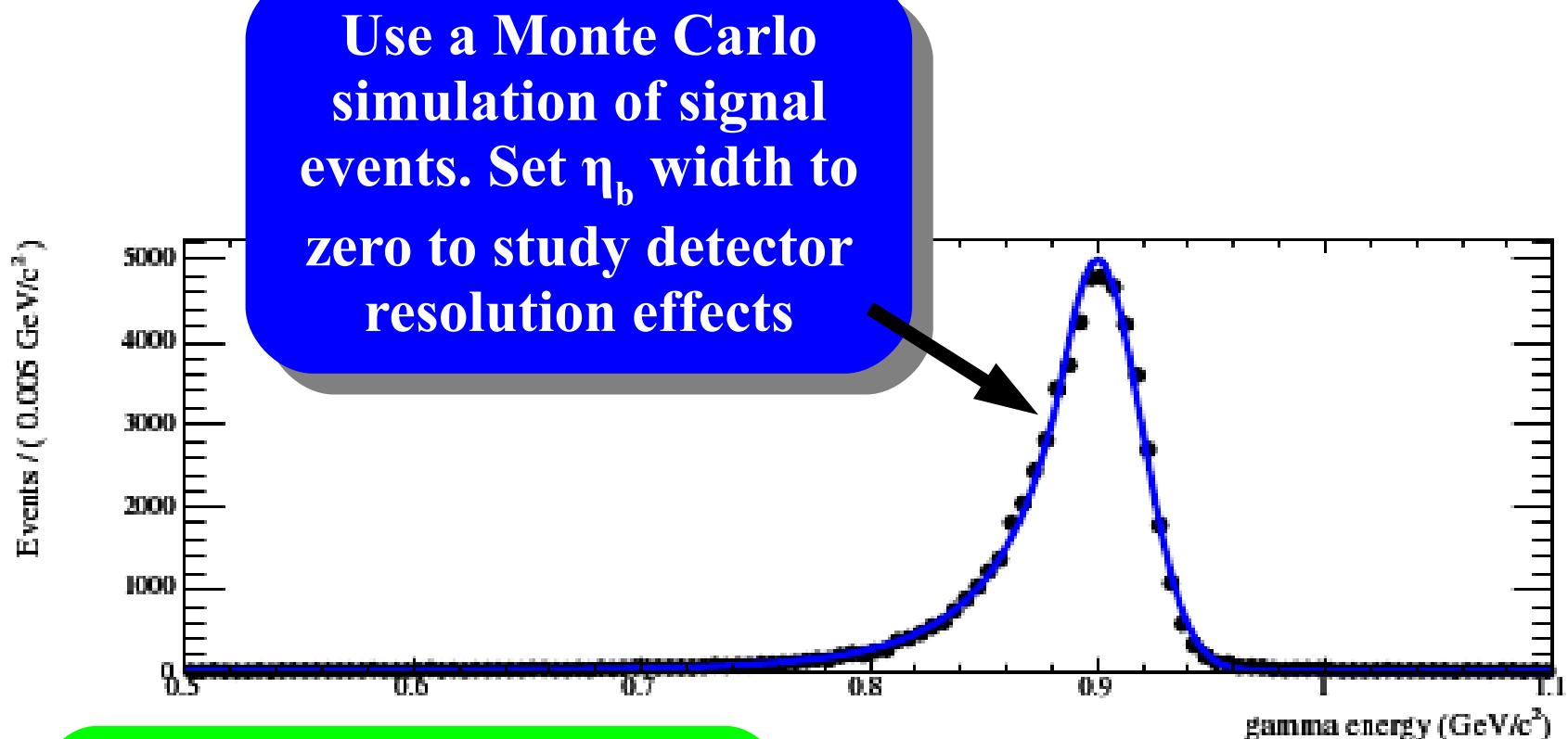
$\sqrt{s} = 10.31 \text{ GeV} \rightarrow \sqrt{s} = 10.3552 \text{ GeV}: 29393 \pm 5014$

The $\chi_{bJ}(2P)$ – background, calibration

*The peak position is shifted by 3.8 MeV below the expectation –
this is used to calibrate the photon energy*



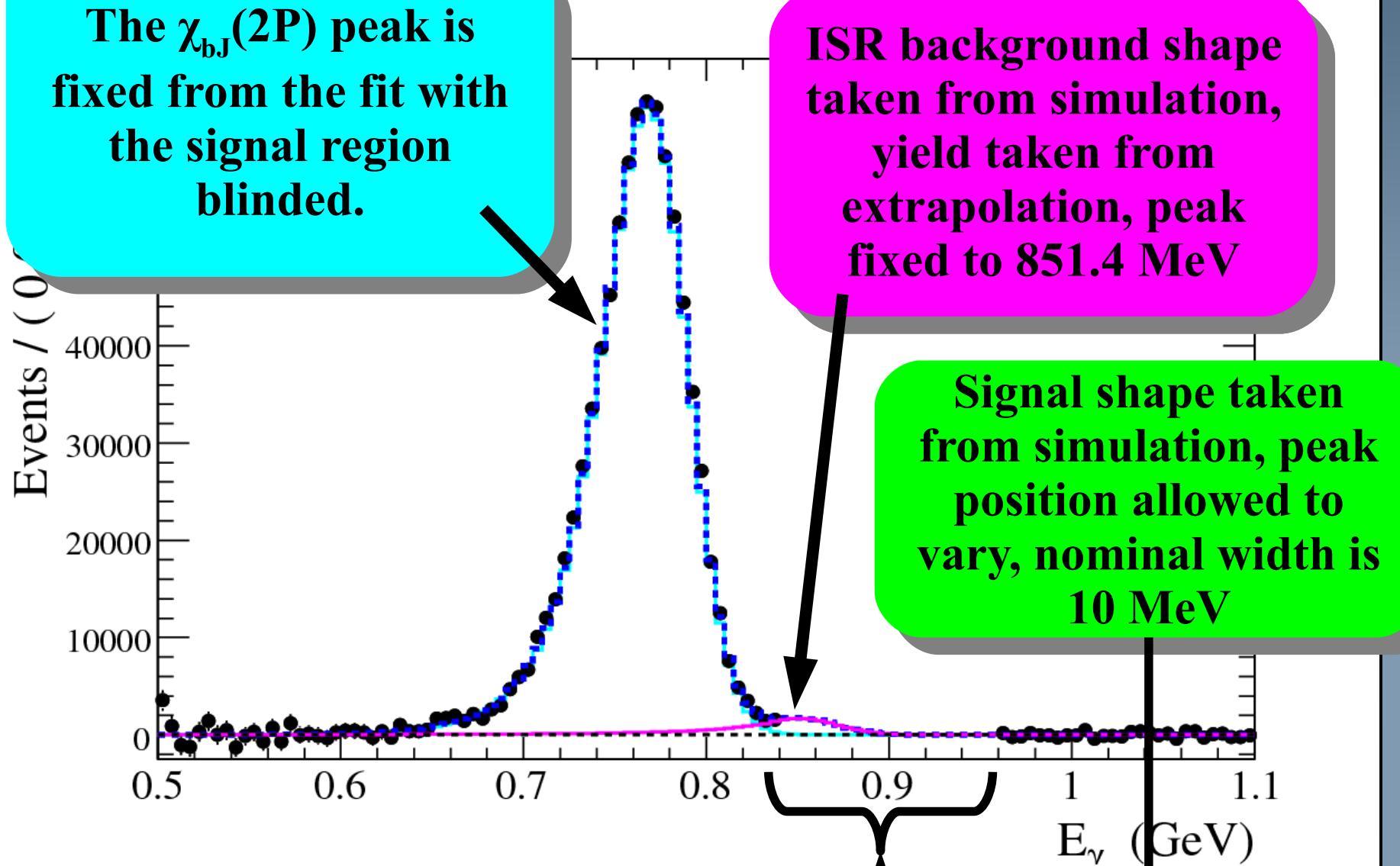
The η_b Signal Model



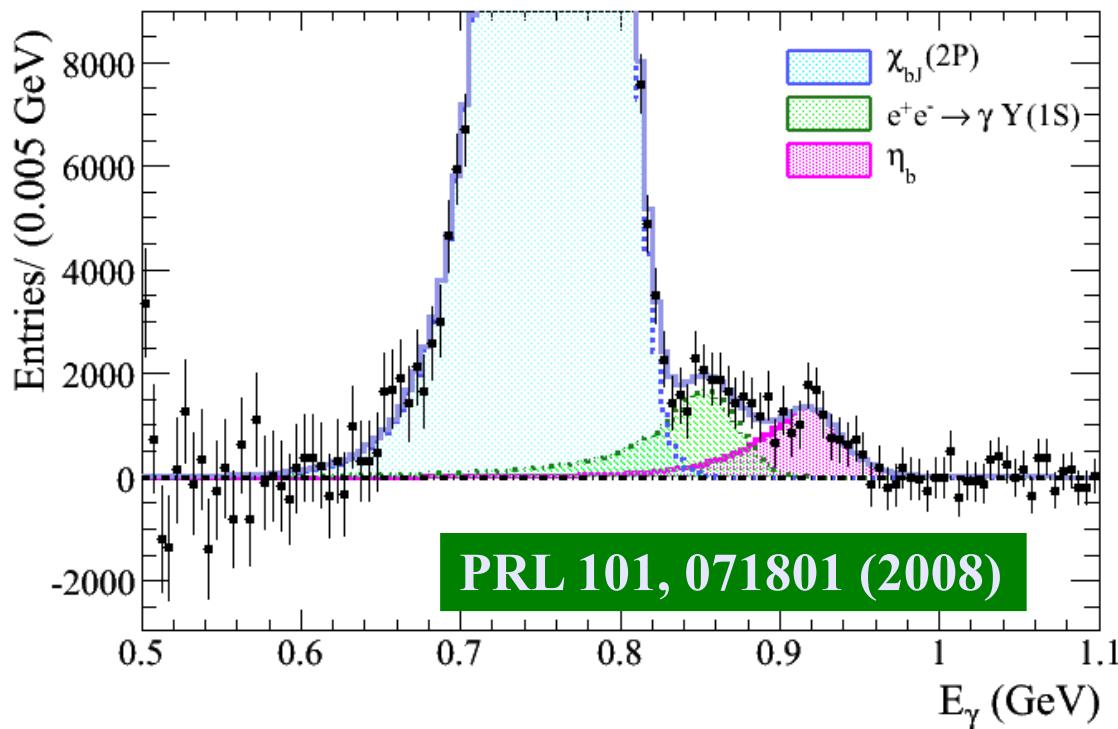
Convolute resolution model with a Breit-Wigner, which represents the resonance

Floating the BW width in the final fit failed to converge. Fix to 10 MeV and vary from 5-20 MeV.

Strategy to search for a “bump”



Results

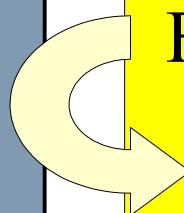


Fitted signal yield:

**$19200 \pm 2000 \text{ (stat.)}$
 $\pm 2100 \text{ (syst.)}$**

Branching Fraction:

$(\xi.8 \pm 1.0 \pm 1.2) \times 10^{-5}$



Fitted Mean: $E_\gamma = 921.2^{+2.1}_{-2.8} \pm 2.4 \text{ MeV}$

$9388.9^{+3.1}_{-2.3} \pm 2.7 \text{ MeV}/c^2$

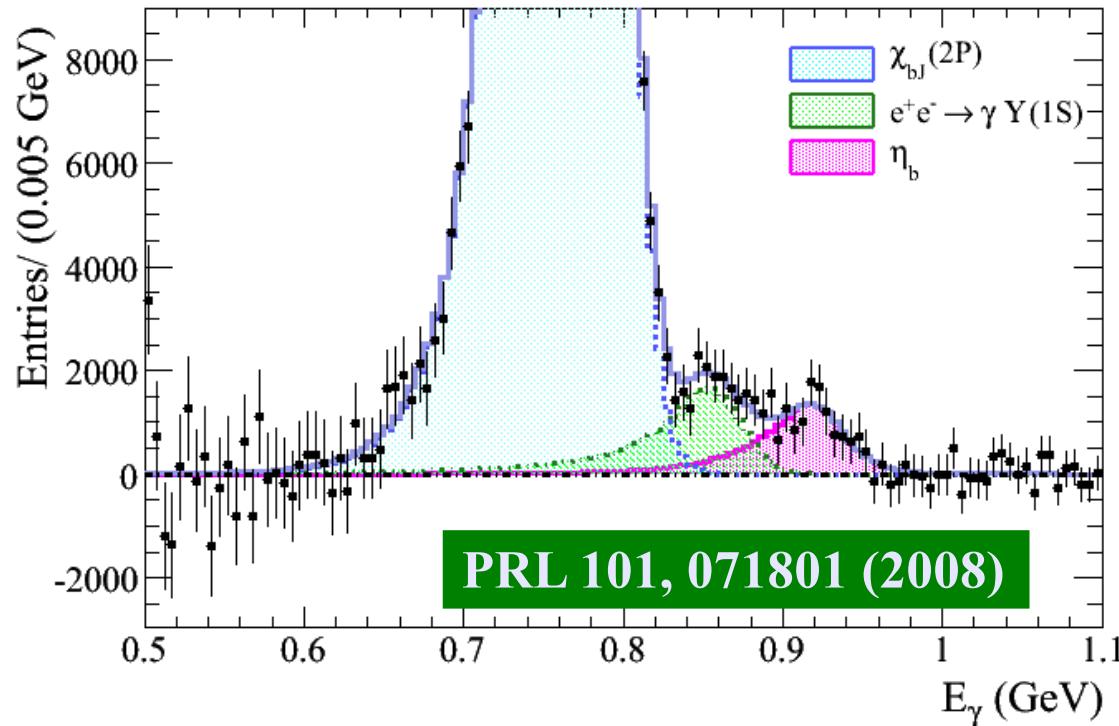
$71.4^{+2.3}_{-3.1} \pm 2.7 \text{ MeV}/c^2$

Mass:
Hyperfine
Splitting:

Consistent with
predictions of the
 η_b properties

Is this really the ground state?

- photon angular distribution can tell us the spin
- are the dominant decay modes to hadrons?
- do we see the “same” state in $Y(2S) \rightarrow \gamma \eta_b$?



A Matter of New Physics: Search for a Light CP-Odd Higgs

Quarks

<i>u</i>	<i>c</i>	<i>t</i>
up	charm	top

<i>d</i>	<i>s</i>	<i>b</i>
down	strange	bottom

<i>e</i>	<i>μ</i>	<i>τ</i>
electron	muon	tau

<i>ν_e</i>	<i>ν_μ</i>	<i>ν_τ</i>
electron neutrino	muon neutrino	tau neutrino

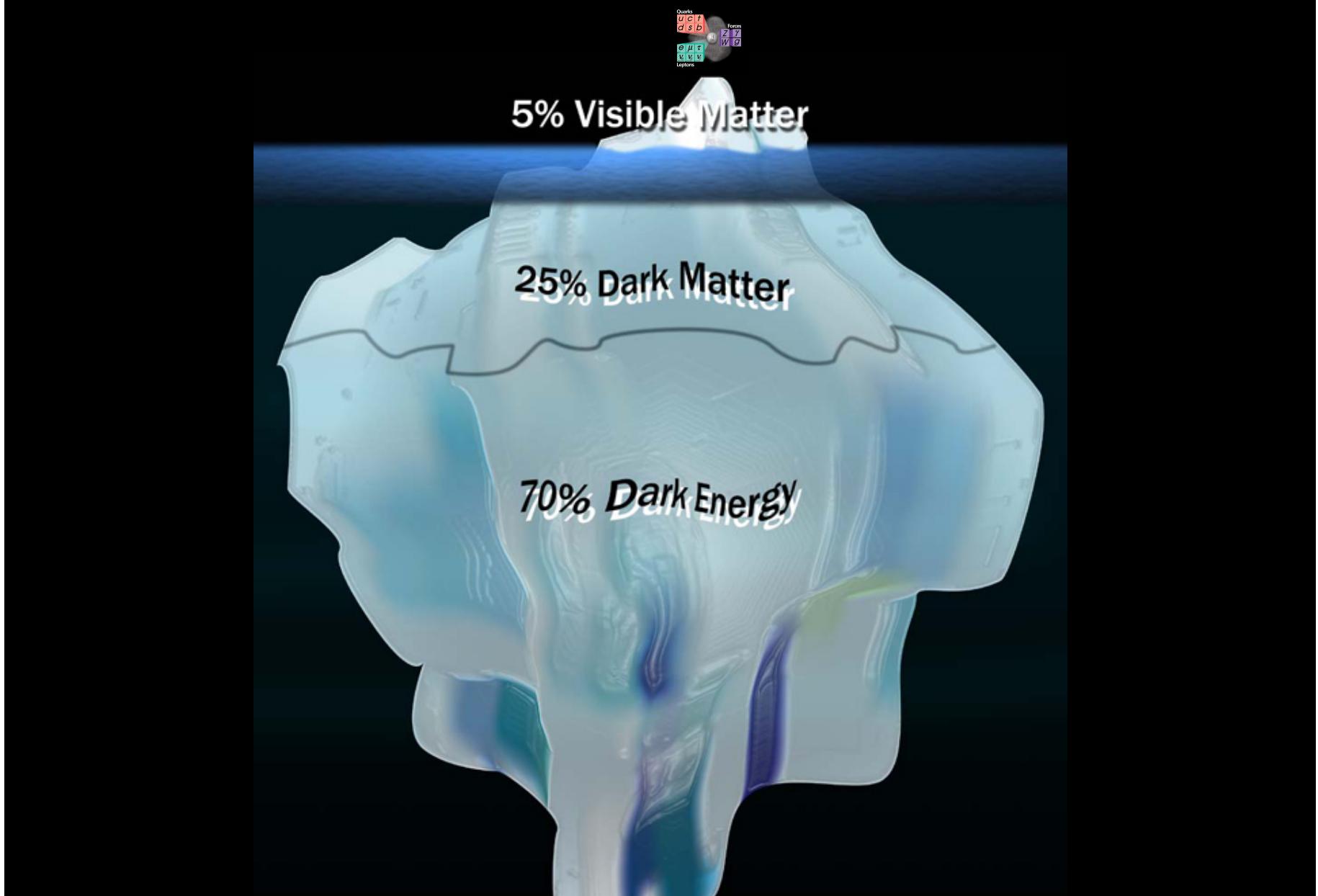
Leptons

Forces

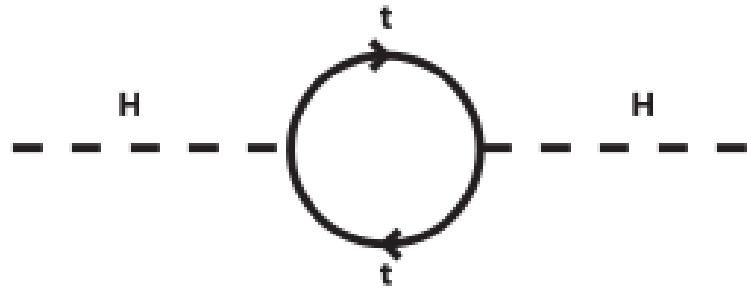
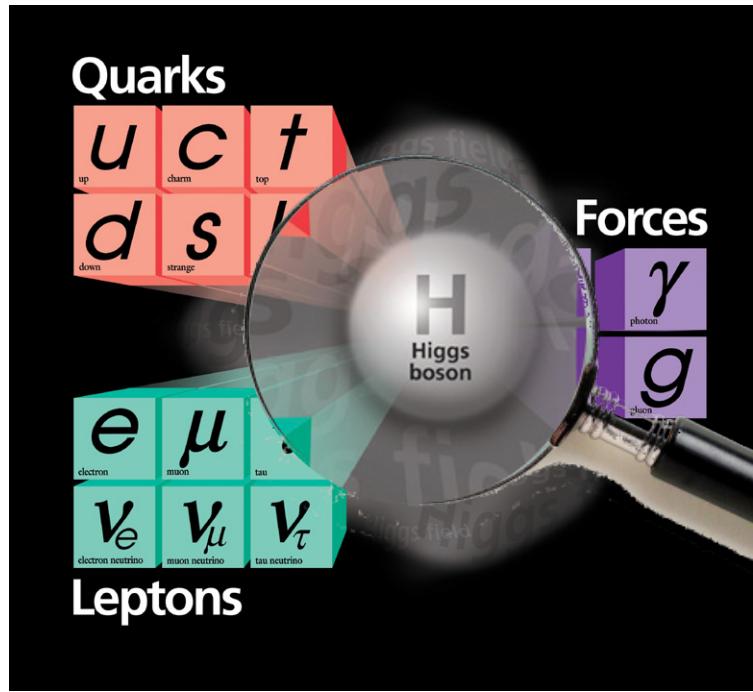
<i>Z</i>	<i>γ</i>
Z boson	photon

<i>W</i>	<i>g</i>
W boson	gluon

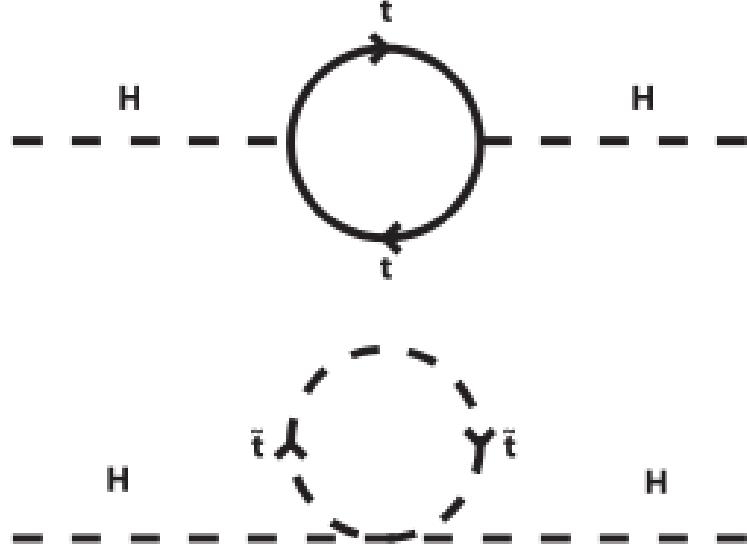
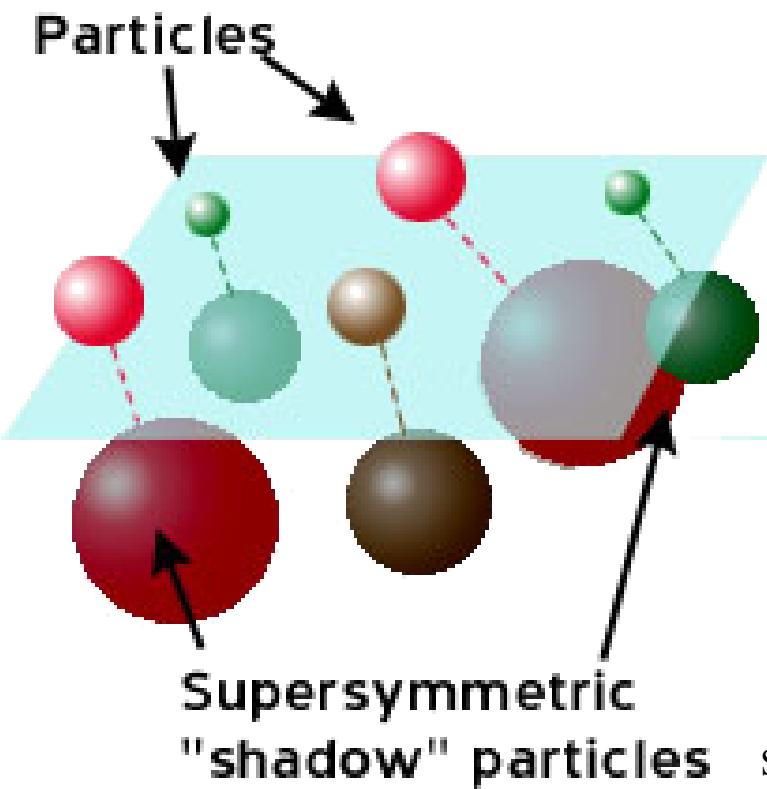
Visible Matter



Can we solve the dark matter puzzle and illuminate the Higgs sector at the same time?



**Higgs self-coupling
diverges in the Standard
Model at high energies**



**Loops involving superpartners
cancel divergences!**

THE “mu” PROBLEM

$$\mu H_u H_d$$

The above term in the superpotential gives the two Higgs doublets non-zero vacuum-expectation values, so that the Higgses can then give mass to the matter particles

μ is then expected to have a value of order the weak scale, far from the next natural scale: the Planck scale. Why is μ so small?

**One Solution:
The Next-to-Minimal Supersymmetric Standard Model (NMSSM)**

$$\mu H_u H_d \longrightarrow \lambda N H_u H_d$$

Add an additional gauge singlet Higgs superfield, effectively promoting μ to a gauge singlet, chiral superfield

This adds a CP-odd Higgs, which I will denote the A^0 , that can radically change the phenomenology of the Higgs sector

New Physics: A Light Higgs Boson

Parameter Scan

blue points: $m_{A^0} < 2m_\tau$

red points: $2m_\tau < m_{A^0} < 7.5 \text{ GeV}$

green points: $7.5 \text{ GeV} < m_{A^0} < 8.8 \text{ GeV}$

black points: $8.8 \text{ GeV} < m_{A^0} < 9.2 \text{ GeV}$

PRL 95:041801,2005 and
PRD 76:051105,2007

For a light A^0 , the dominant SM Higgs decay will be $h \rightarrow A^0 A^0$

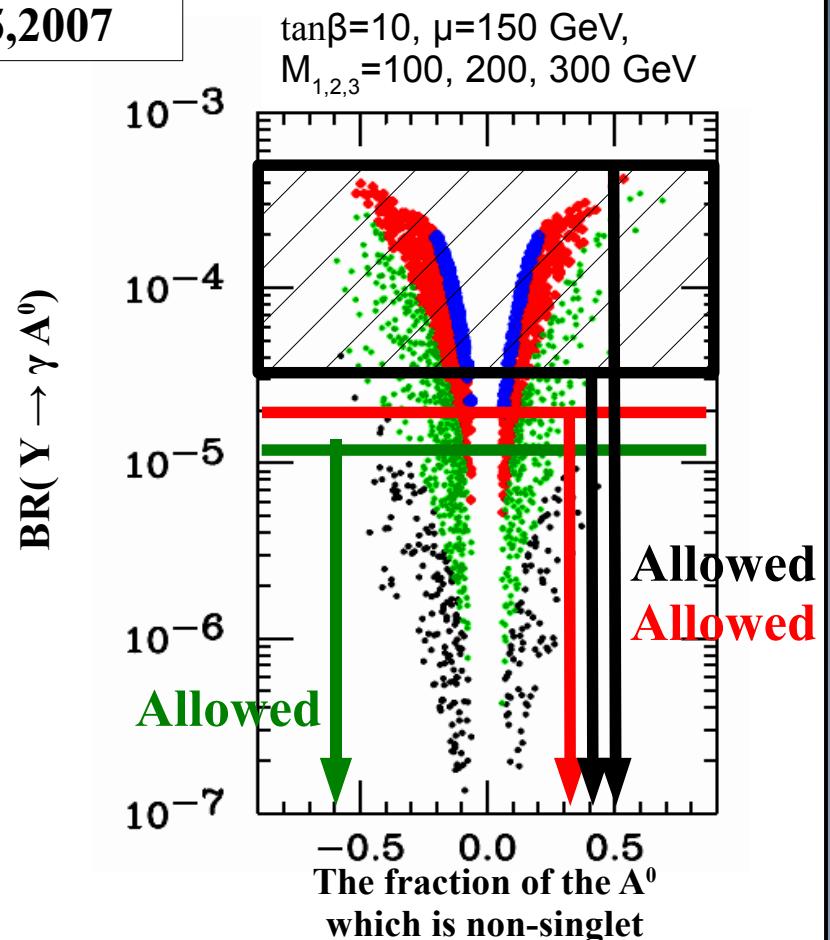
Depending on the A^0 mass, the dominant decay could be:

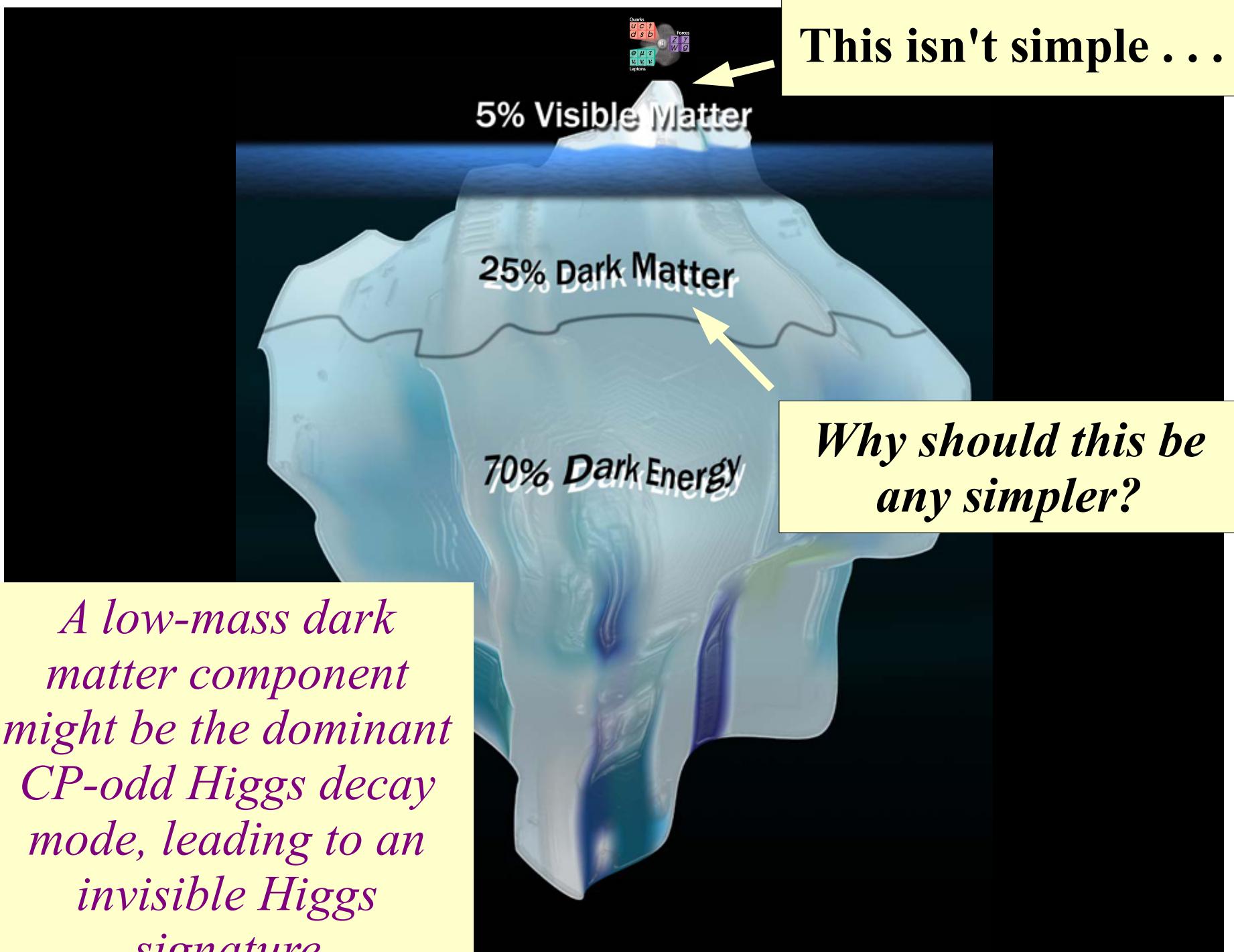
$$A^0 \rightarrow \tau^+ \tau^-$$

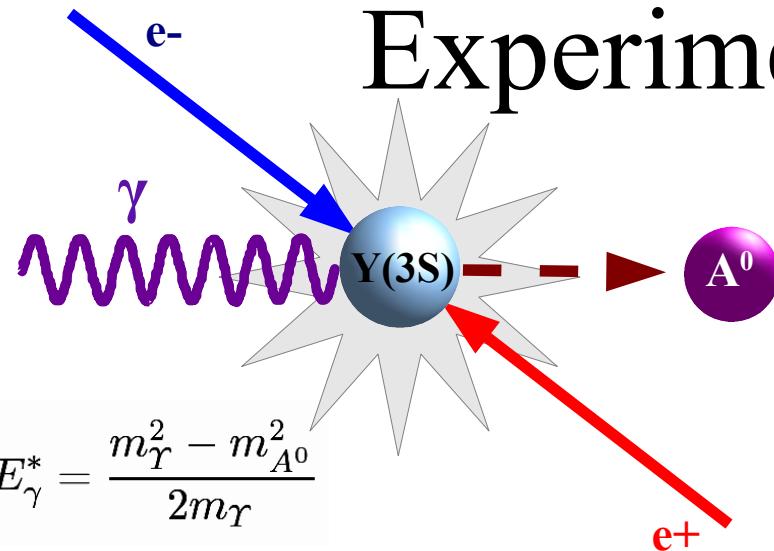
Leading to 4- τ final states at LEP, which were never explored.

Best limits come from recent CLEO search for $A^0 \rightarrow \mu\mu, \tau\tau$

hep/ex arXiv:0807.1427

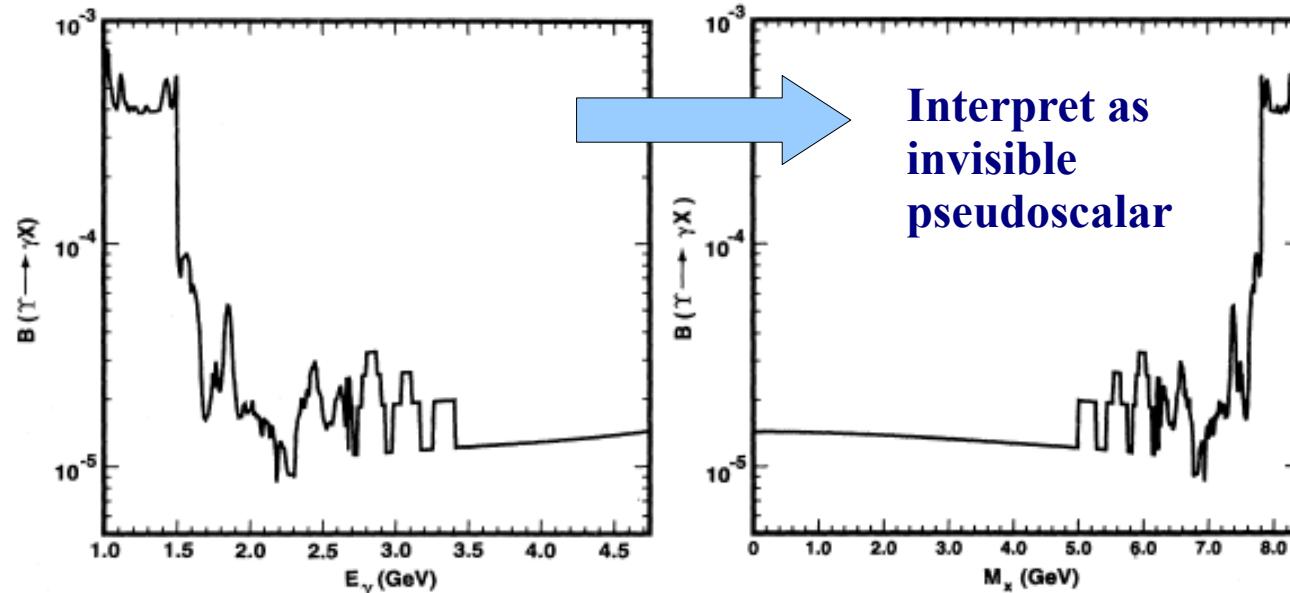
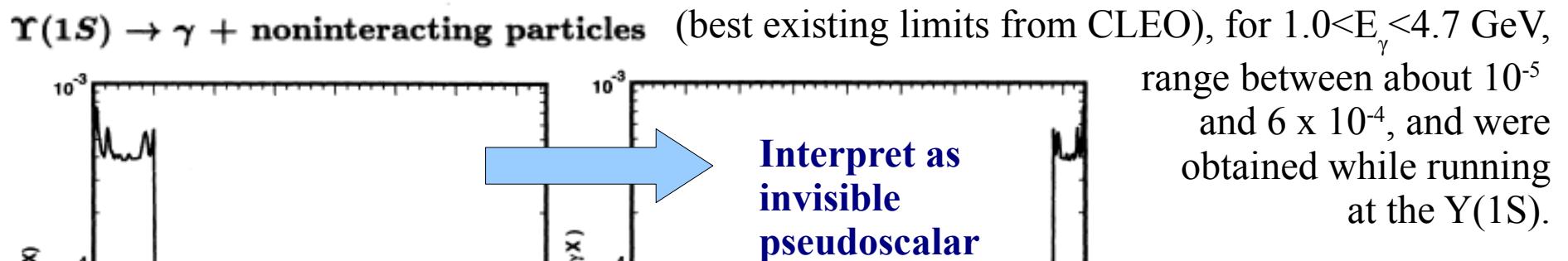






Experimental Signature

Search for an invisibly-decaying particle recoiling against a single photon



There are no limits from the $\Upsilon(3S)$ or the $\Upsilon(2S)$.

**Selection of high-quality photons,
with tighter
criteria for lower
photon energies
(increasing
backgrounds)**

An illustrative signal candidate event . . .

Require very
little additional
detector activity
either in tracking
or in the
calorimeter

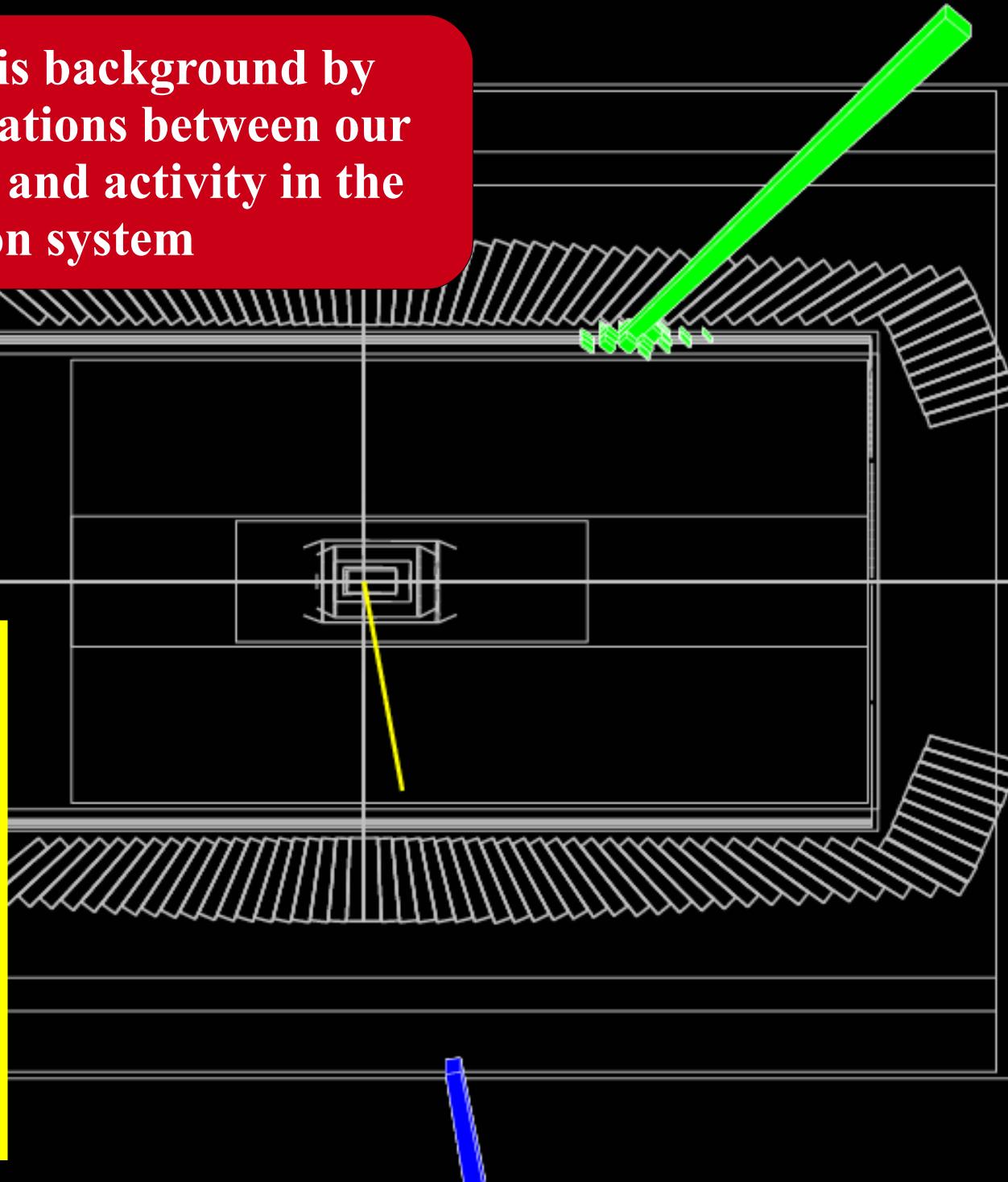
**One catch: this event is a data event from a
problematic background: $e^+ e^- \rightarrow \gamma \gamma$**

We reject this background by
vetoing correlations between our
signal photon and activity in the
muon system

Total Signal Efficiency:

High Energy
Region: 10-11%

Low Energy
Region: 20%



Maximum Likelihood Fit

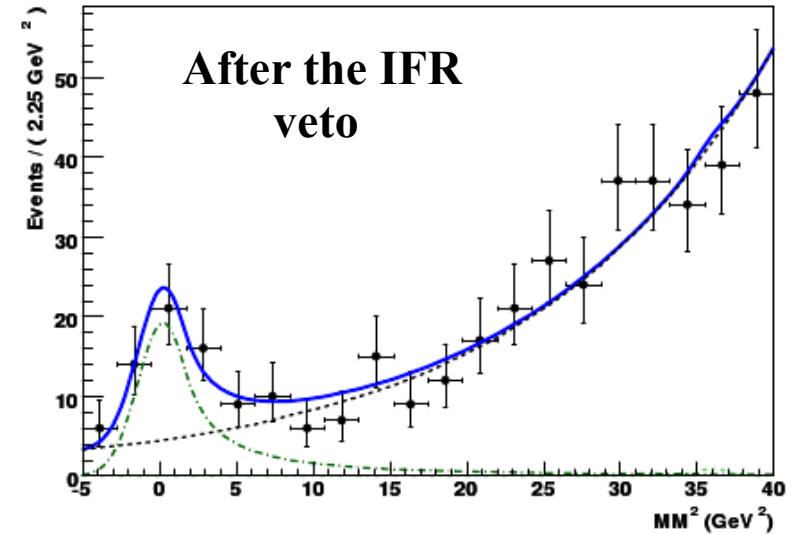
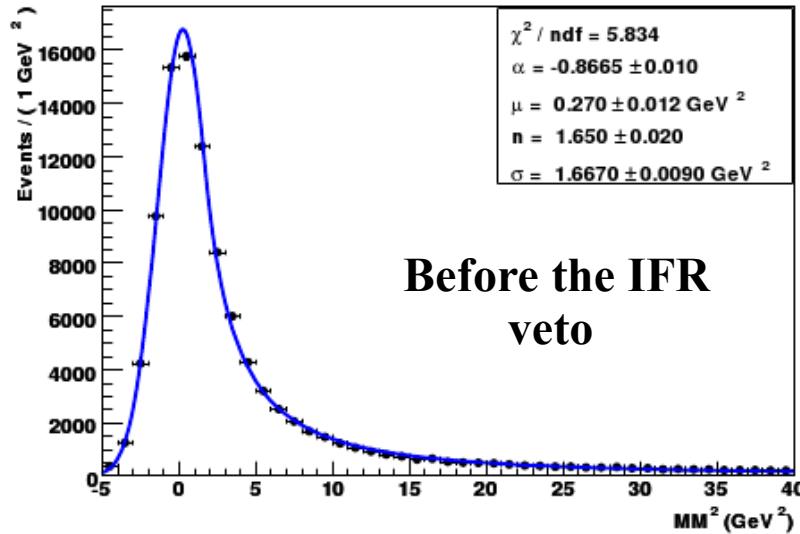
- 1-D fit to the missing mass-squared:

$$m_X^2 = M_{Y(3S)}^2 - 2 E_\gamma^* M_{Y(3S)}$$

- Signal model
 - parameterized using same detector resolution function as η_b search (zero-width Higgs)
 - parameters vary with assumed Higgs mass, due to calorimeter response
- Background models
 - determined from data control samples
 - Major backgrounds: $e^+e^- \rightarrow \gamma\gamma, \gamma\gamma\gamma, e^+e^-\gamma$

$e^+e^- \rightarrow \gamma\gamma$ background

Off-Resonance Y(3S) Data

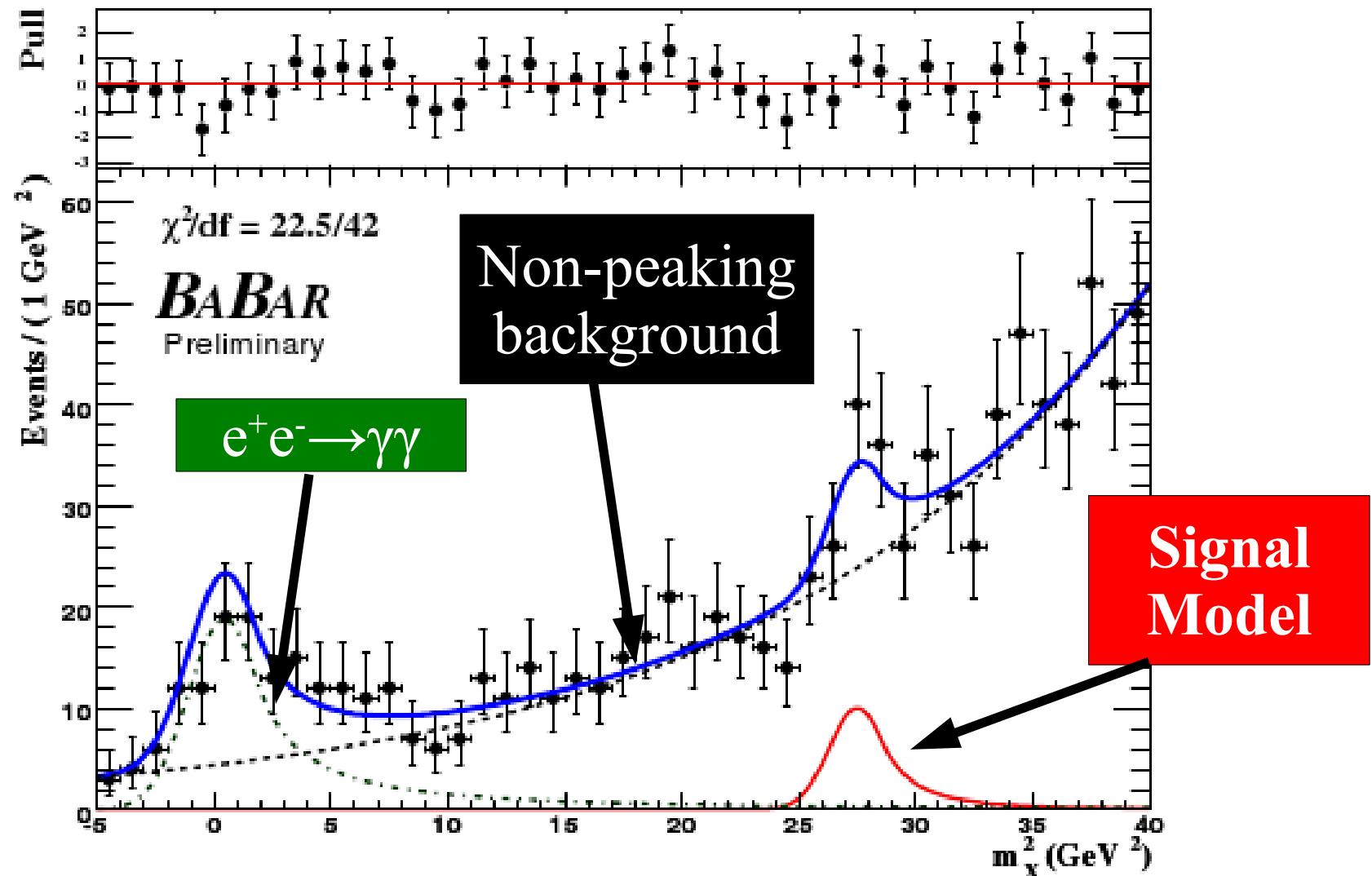


The $\gamma\gamma$ background model is taken from data before the IFR veto.

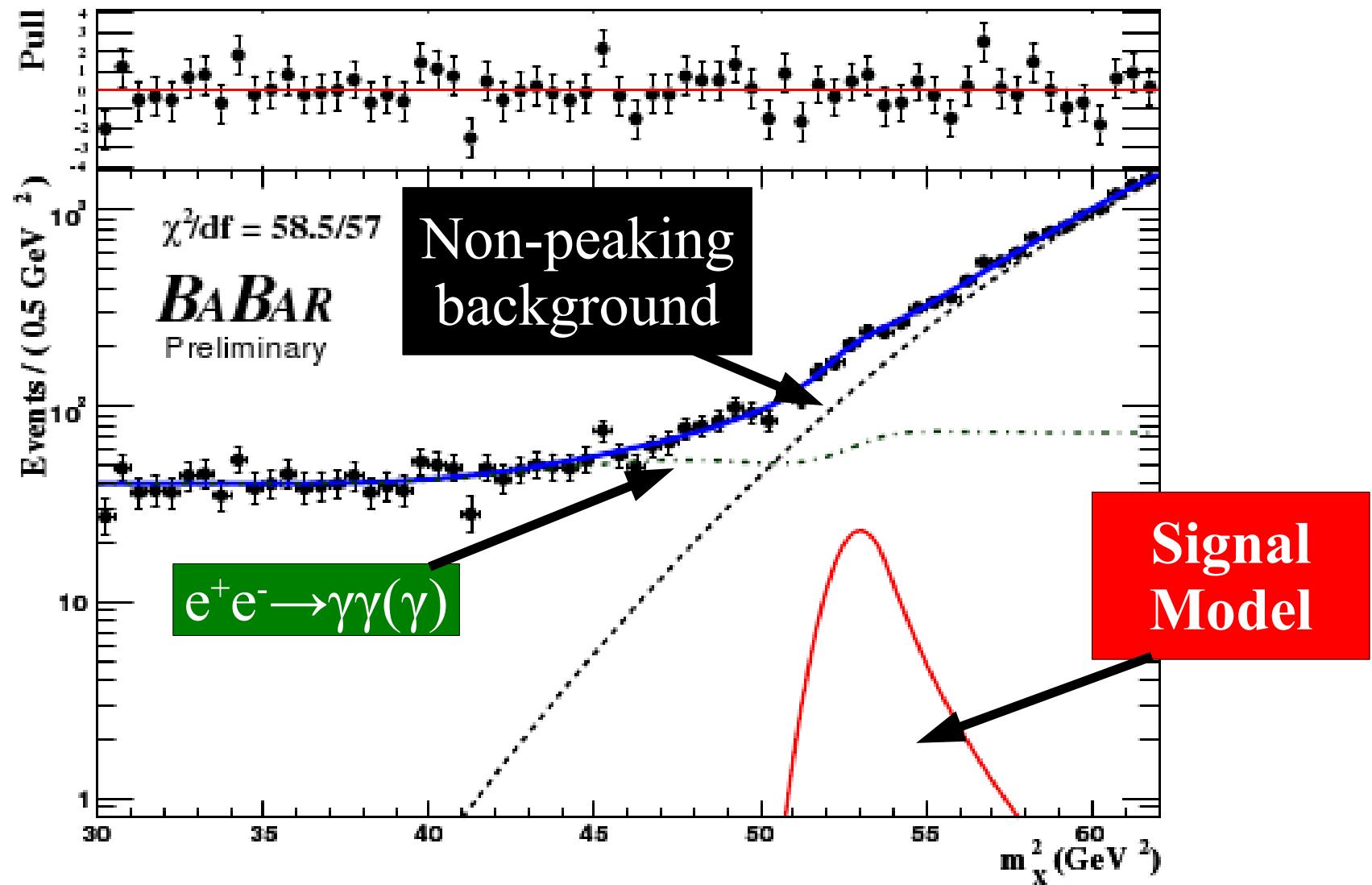
The tail of this function extends into the high-mass region, and is treated as part of the background there.

The remaining non-peaking background in the low-mass region is empirically modeled as a single exponential function.

A Snapshot: Fits to the Spectrum Low-Mass Region



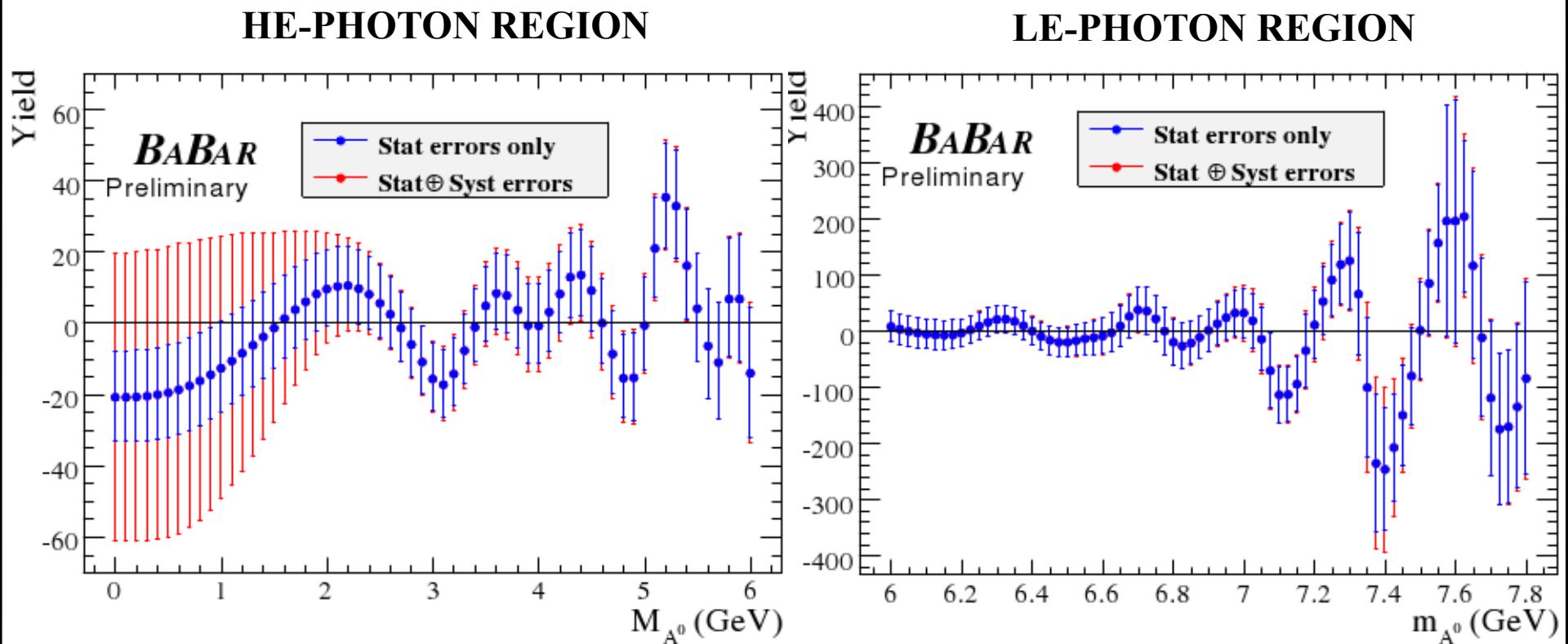
A Snapshot: Fits to the Spectrum High-Mass Region



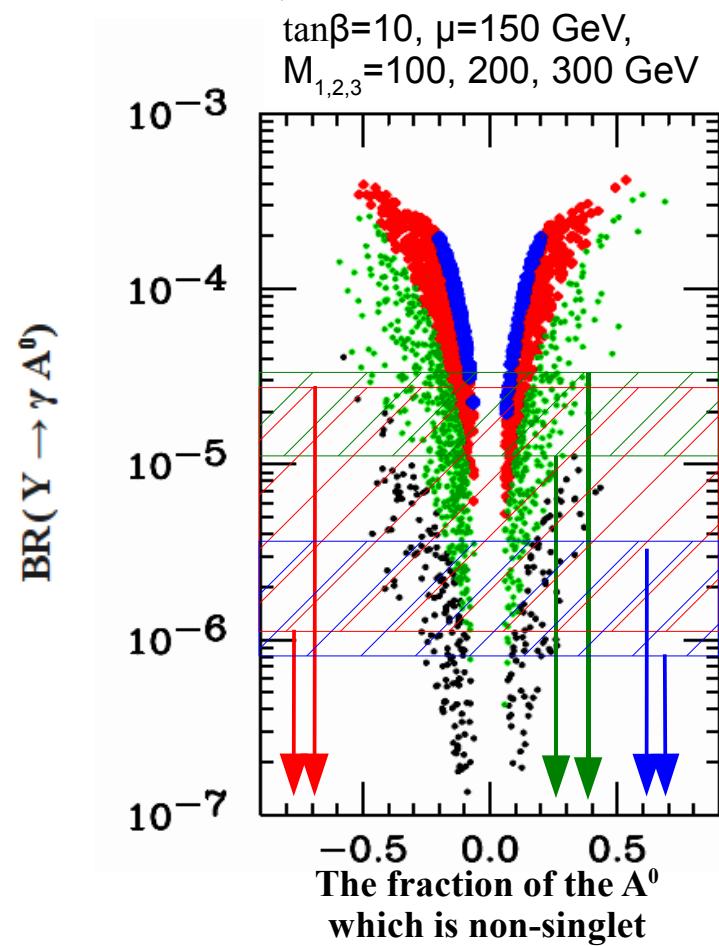
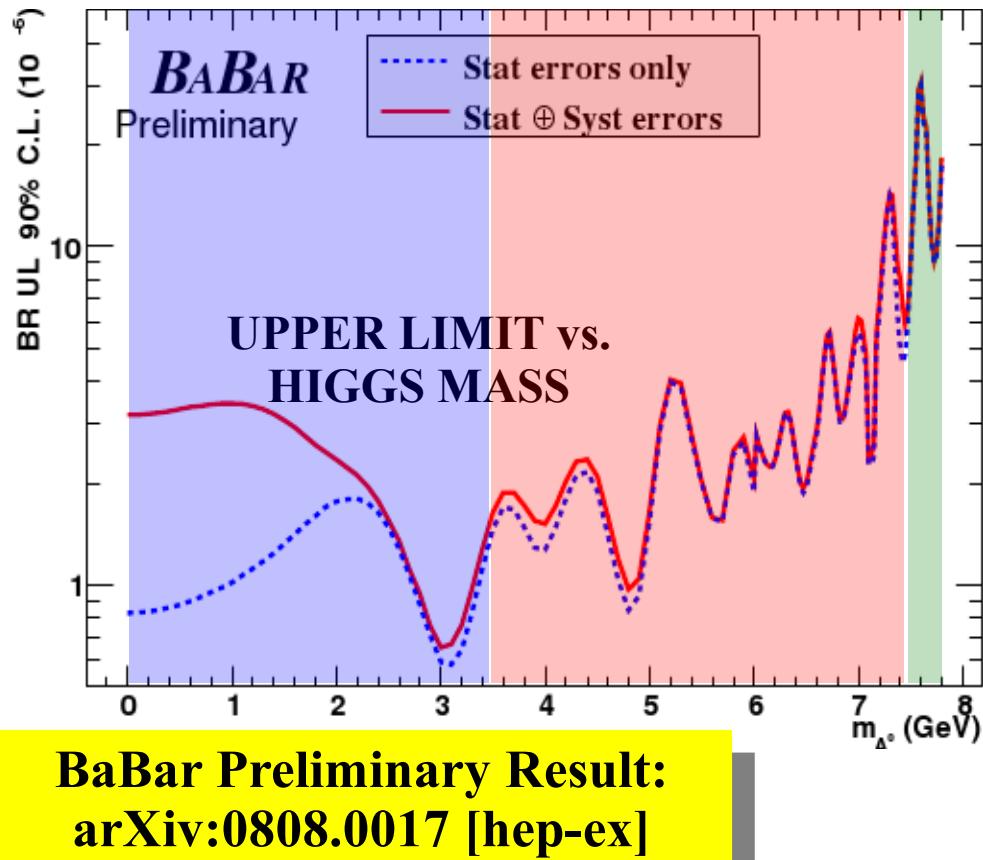
Results

Most significant yields:

- low-mass region: 37 ± 15 (2.6σ , stat. only)
- high-mass region: 119 ± 71 (1.7σ , stat. only)



Results (continued)



Concluding Thoughts: Prospects for Further Discovery

First Results from BaBar Upsilon Sample

- Unmatched samples of Upsilon mesons below threshold open up new doors of exploration
 - Standard Model – discovery and further study of the η_b

Mass: $9388.9^{+3.1}_{-2.3} \pm 2.7 \text{ MeV}/c^2$

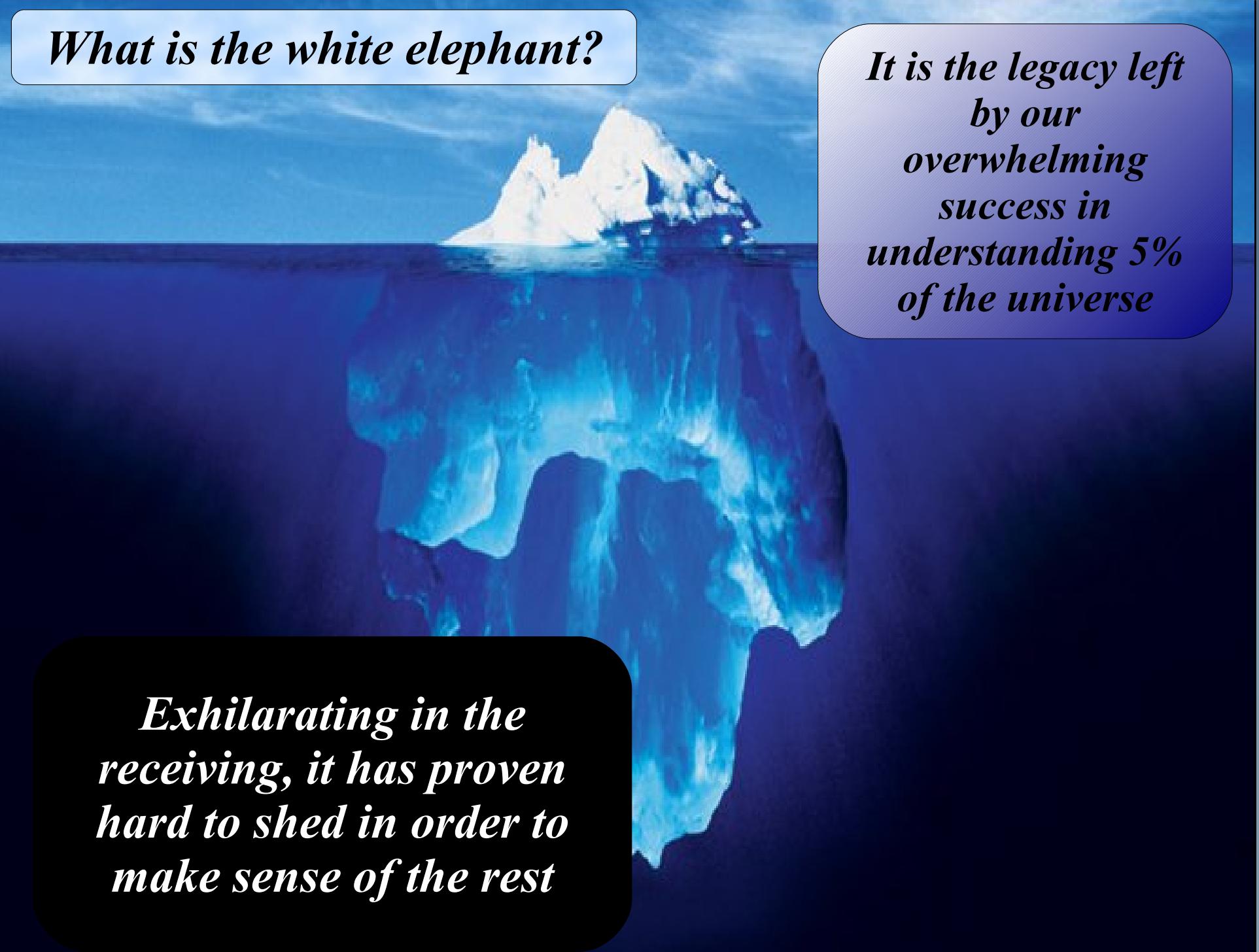
Hyperfine Splitting: $71.4^{+2.3}_{-3.1} \pm 2.7 \text{ MeV}/c^2$

Branching Fraction:

$$BR(Y(3S) \rightarrow \gamma \eta_b) = (4.8 \pm 0.5 \pm 1.2) \times 10^{-4}$$

- New Physics – searches for low-mass Higgs and light dark matter
 - We exclude an invisibly decaying light Higgs up to $7.8 \text{ GeV}/c^2$ at the 90% CL at the level of $10^{-5} \text{ -- } 10^{-6}$

What is the white elephant?



*It is the legacy left
by our
overwhelming
success in
understanding 5%
of the universe*

*Exhilarating in the
receiving, it has proven
hard to shed in order to
make sense of the rest*

Backup Slides: Reference and Details

References

QCD Calculations of the η_b mass and branching fraction

Recksiegel and Sumino, Phys. Lett. B 578, 369 (2004) [hep-ph/0305178]

Kniehl et al., PRL 92 242001 (2004) [hep-ph/0312086]

Godfrey and Isgur, PRD 32, 189 (1985)

Fulcher, PRD 44, 2079 (1991)

Eichten and Quigg, PRD 49, 5845 (1994) [hep-ph/9402210]

Gupta and Johnson, PRD 53, 312 (1996) [hep-ph/9511267]

Ebert et al., PRD 67, 014027 (2003) [hep-ph/0210381]

Zeng et al., PRD 52, 5229 (1995) [hep-ph/9412269]

Spectroscopy

N. Brambilla et al., “Heavy Quarkonium Physics,” hep-ph/0412158 (December 13, 2004),
<http://arxiv.org/abs/hep-ph/0412158>.

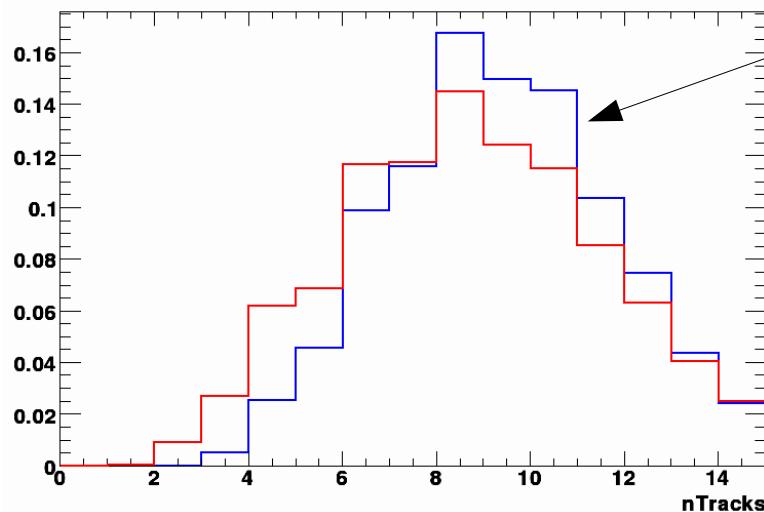
M. Artuso et al., “Photon Transitions in Upsilon(2S) and Upsilon(3S) Decays,” Physical Review Letters 94, no. 3 (January 28, 2005): 032001-5

η_b Event Pre-selection

- Selection chosen to have high signal efficiency
 - Dominant η_b decay expected to be $\eta_b \rightarrow gg$
 - require ≥ 4 charged tracks in an event
 - exclude “jetty” events (e.g. $e^+e^- \rightarrow qq$) using Fox-Wolfram moment ratio, $H_2/H_0 < 0.98$
 - Select high-quality photons:
 - lateral moment of EMC shower < 0.55
 - EMC barrel-only photons ($-0.762 < \cos\theta_\gamma < 0.890$)
 - Spin-0 η_b leaves a small correlation between the photon and event thrust axis, in contrast to $e^+e^- \rightarrow qq$: $|\cos\theta_T| < 0.7$
 - Veto photons consistent with a π^0 decay

η_b – track multiplicity

Track multiplicity after all other cuts, compared between signal MC (BLUE) and the test data (RED)



According to MC simulation, the ≥ 4 track multiplicity is 99.5% efficient on signal events: check signal simulation against $\chi_{bJ}(2P)$ data!

Despite the expected higher multiplicity of the $\chi_{bJ}(2P) \rightarrow \gamma Y(1S)$ events (due to $Y(1S) \rightarrow ggg$), the difference in the efficiencies due to the track multiplicity cut is only about 10%. We conservatively assign this as part of the selection efficiency systematic

Cut	S/\sqrt{B}	Eff. (from χ_b peak)	Eff. (signal MC)
No cut	101.5	-	0.629
BGFMultiHadron	109.8	0.973	0.977
≥ 4 Charged Tracks	107.2	0.903	0.995
$ \Delta T < 0.55$	113.2	0.997	0.991
$-0.762 < \cos(\theta_{\gamma, LAB}) < 0.890$	109.6	0.928	0.901
$ \cos(\theta_T) < 0.7$	135.2	0.672	0.690
π^0 -50 MeV cut	164.7	0.849	0.899

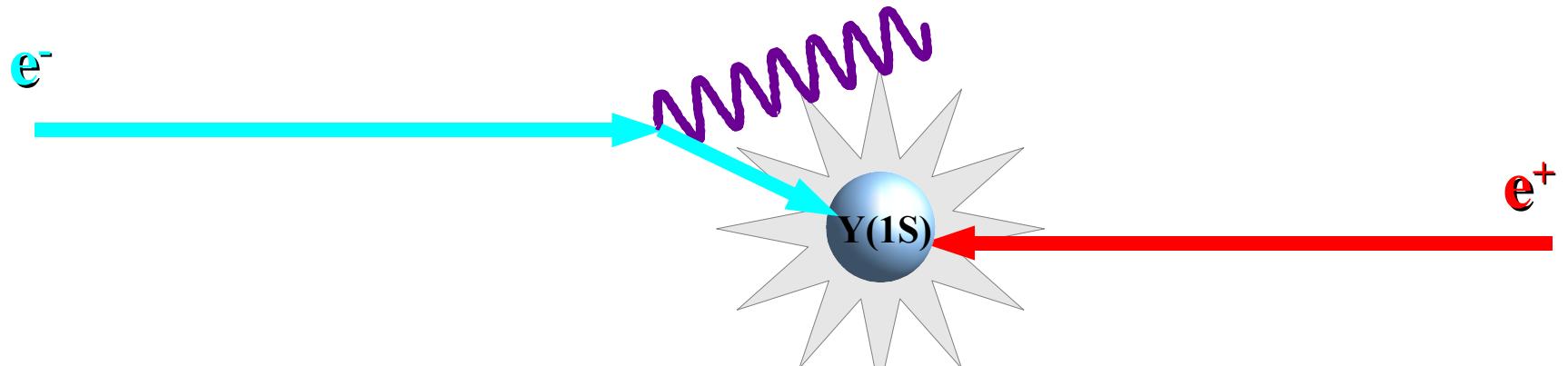
The η_b width

- Predictions of the width:
 - based on the ratio of $\Gamma(\eta_b \rightarrow \gamma\gamma)$ and $\Gamma(\eta_b \rightarrow gg)$, predictions range from 4-20 MeV/c²
 - c.f. *W. Kwong et al., Phys. Rev. D 37, 3210 (1988); C. S. Kim, T. Lee, and G. L. Wang, Phys. Lett. B 606, 323 (2005); J. P. Lansberg and T. N. Pham, Phys. Rev. D 75, 017501 (2007).*
- Systematic variations:
 - fit with width floated won't converge
 - variations from 5-20 MeV/c² lead to largest single systematic uncertainty on yield (10%)

The Details of the η_b Fit

- The fit is done using a maximum likelihood function on the binned data, $0.5 < E_\gamma < 1.1$ GeV
- bin size: 5 MeV
- Fit models
 - non-peaking parameters floated, with initial values set from the peaking-region-blinded fit
 - $\chi_{bJ}(2P)$ shape fixed, yield floated
 - ISR shape fixed, yield fixed
 - signal shape fixed, except the peak position; yield floated

$e^+e^- \rightarrow \gamma_{ISR} Y(1S)$: Expectation

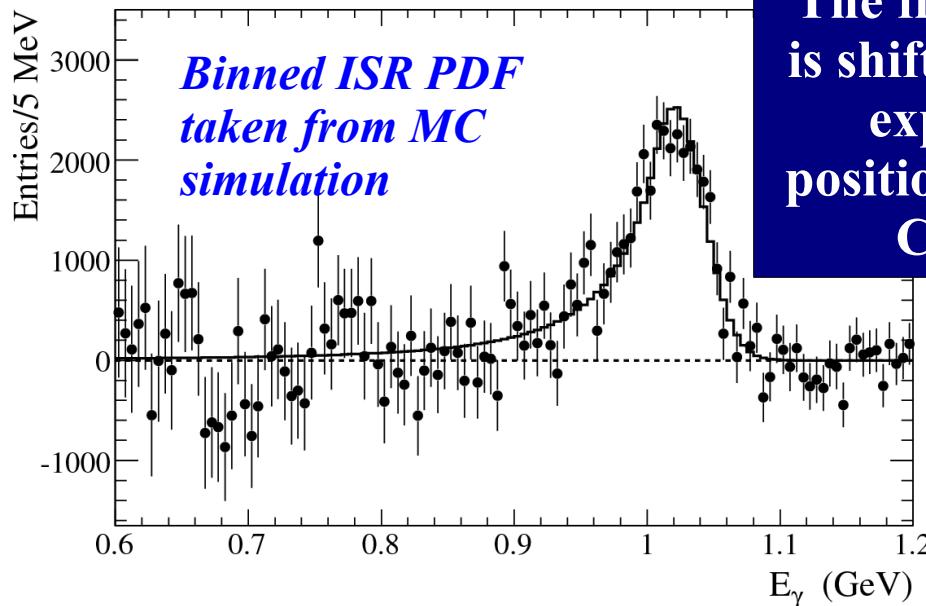
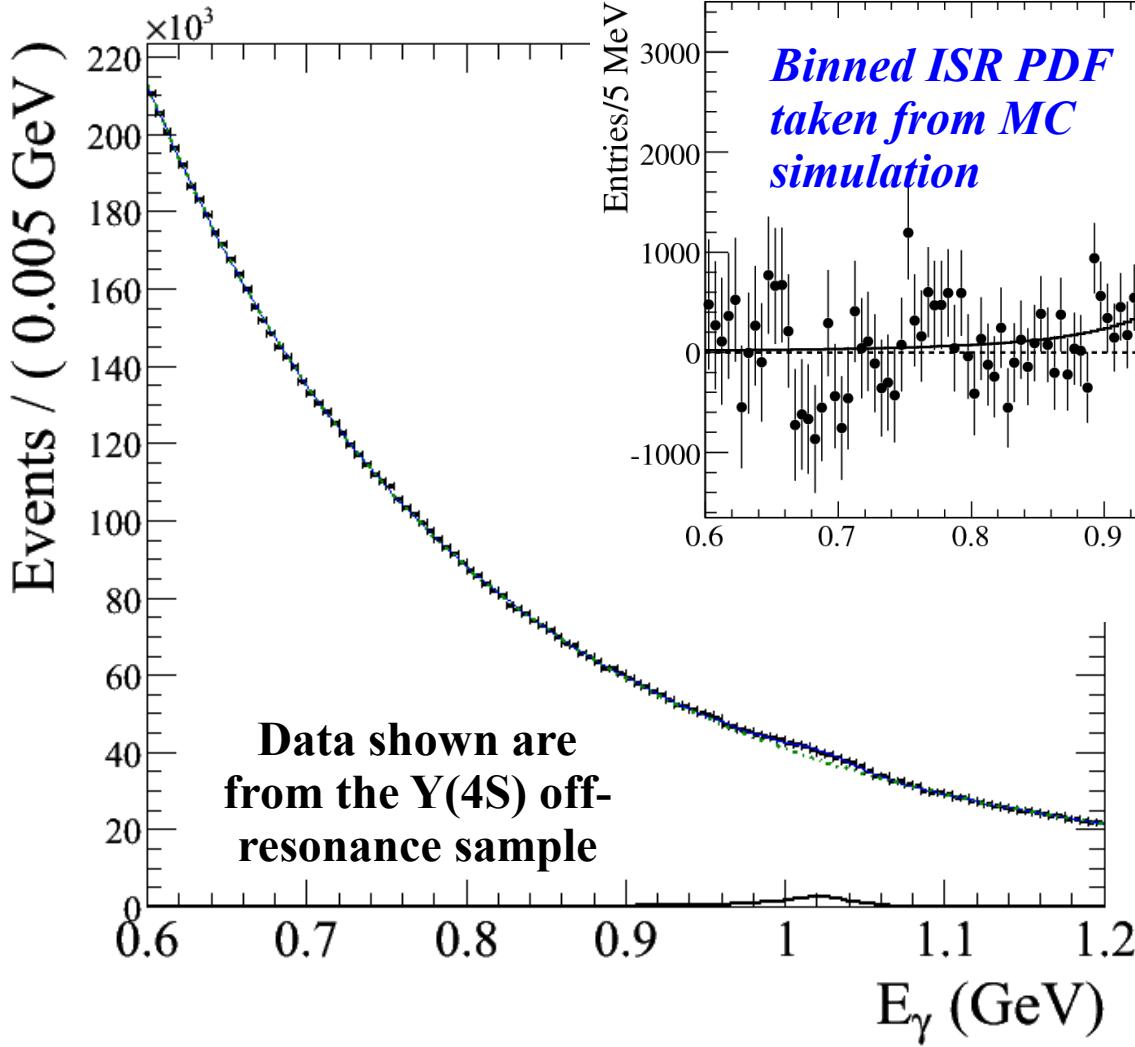


$$\sigma(e^+e^- \rightarrow \gamma_{ISR} Y(1S)) = \frac{12\pi}{M_{Y(1S)} s} \cdot \Gamma_{ee} \cdot W(s, 2E_\gamma/\sqrt{s})$$

$$N_{\sqrt{s}=M_{Y(3S)}} = N_{\sqrt{s}'} \frac{\sigma_{\sqrt{s}=M_{Y(3S)}} \epsilon_{\sqrt{s}=M_{Y(3S)}}}{\sigma_{\sqrt{s}'} \epsilon_{\sqrt{s}'}}$$

Use the ratio of cross-sections and efficiencies to cancel most of the uncertainties from either source.

Sample	Lumi [fb ⁻¹]	Cross-Section [pb]	Reconstruction Efficiency	Yield	Extrapolation to $Y(3S)$ On-Peak
$Y(3S)$ Off-Peak	2.415	25.4	5.78 ± 0.09	2773 ± 473	29393 ± 5014
$Y(4S)$ Off-Peak	43.9	19.8	6.16 ± 0.12	35759 ± 1576	25153 ± 1677

$$e^+ e^- \rightarrow \gamma_{\text{ISR}} Y(1S)$$


The fitted ISR shape is shifted down to the expected peak position for the Y(3S) CM energy.

We use 40/fb of data taken 40 MeV below the Y(4S) resonance to study ISR production of the Y(1S).

The data are fitted with the same non-peaking model and a Gaussian + Power-Law Tail (ISR peak).

Systematic Uncertainties - η_b

Signal Yield:

ISR Background:

- fit with ISR yield floated – consistent with the fixed yield of ISR, and has no effect on η_b yield or peak position

- fixed value varied by $\pm 1\sigma$ to get systematic on signal yield
 η_b width varied in fit (5, 15, 20 MeV), yielding largest single systematic effect: 10%

PDF parameters – varied by $\pm 1\sigma$

TOTAL UNCERTAINTY: 11%

Mass:

χ_{bJ} (2P) peak shift: (3.8 ± 2.0) MeV

Branching Fraction:

Selection efficiency: compare data yield to expectation from PDG branching fractions (18%) and MC efficiency – 22% uncertainty

TOTAL UNCERTAINTY: 25%

Data Samples

- Data with single-photon triggers:
 - 28 fb^{-1} taken at the Y(3S)
 - signal analysis sample
 - 4.7 fb^{-1} taken at the Y(4S)
 - used HE trigger, can be used to tune cuts on photons
 - “Off-resonance” data
 - 2.6 fb^{-1} taken 40 MeV below the Y(3S)
 - 0.97 fb^{-1} taken 30 MeV below the Y(2S)
 - 4.5 fb^{-1} taken in a scan above the Y(4S)

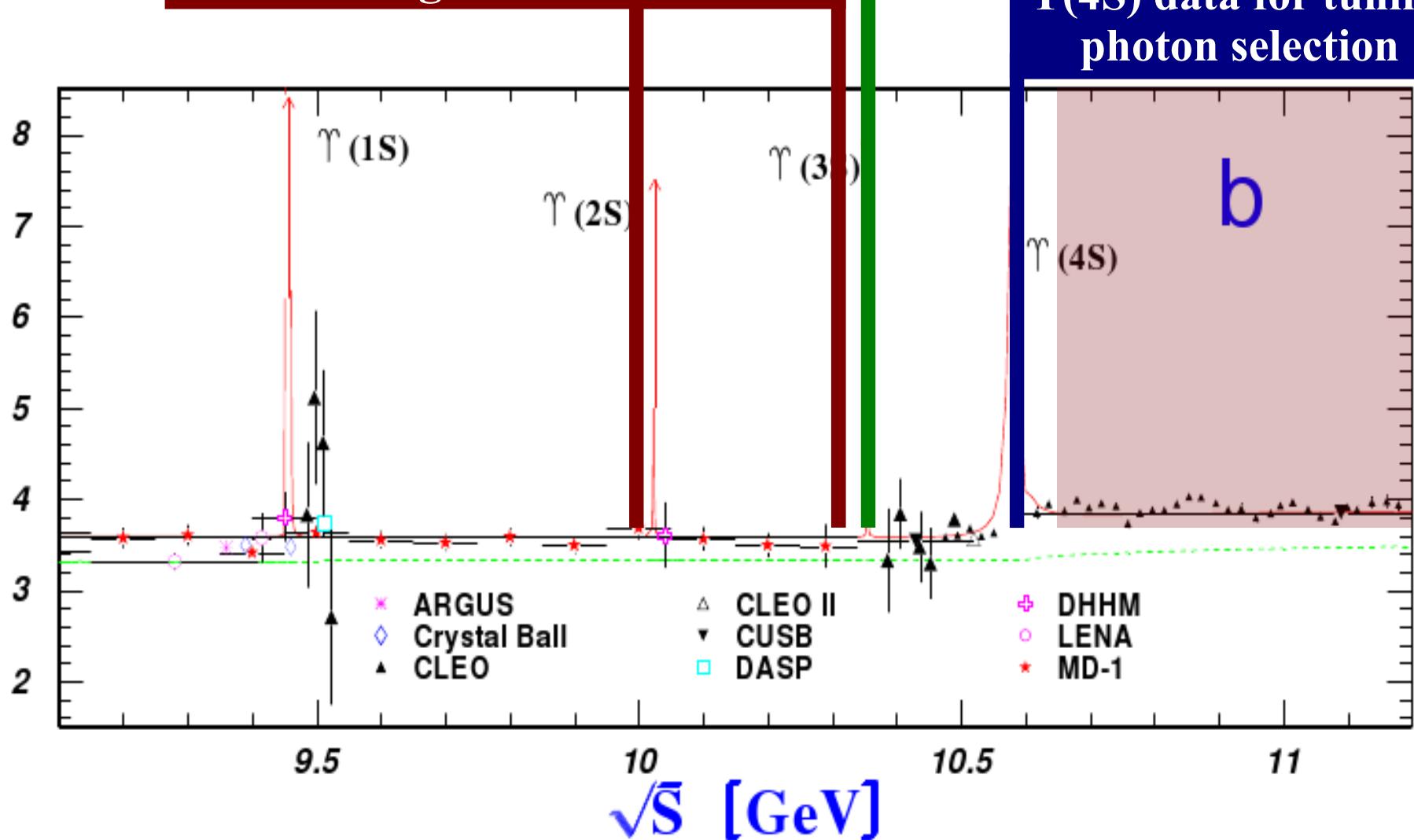
**Data For
Tuning Cuts
and Studying
Backgrounds**

The plot below is taken from arXiv:hep-ph/0312114v2 and is meant to illustrate the $e^+e^- \rightarrow$ hadrons spectrum between 9.1 GeV and 11.2 GeV

**Data taken away from
resonances or above the Y(4S)
– background studies**

122M Y(3S)

**Y(4S) data for tuning
photon selection**



Triggering on Single Photons + E

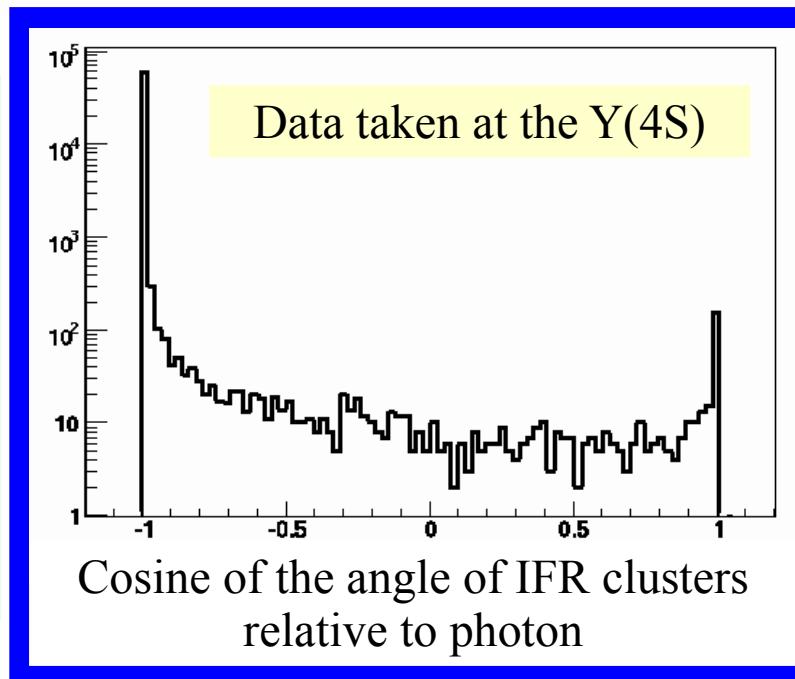
The ability to trigger on events with a single photon and significant missing energy is critical to this analysis

- Dedicated online triggered and filtering were developed
 - Level 1 (hardware trigger): require at least one EMC cluster with energy > 800 MeV (lab frame)
 - Level 3 (software trigger): two lines developed
 - High-energy (HE) line: require isolated EMC cluster with CM-frame energy > 2 GeV
 - Low-energy (LE) line: developed later (only 82 million Y(3S) taken), requires cluster energy > 1 GeV and no tracks from the IP
- 100 Hz

Event Selection

Variable	$3.2 < E_{\gamma}^* < 5.5 \text{ GeV}$	$2.2 < E_{\gamma}^* < 3.7 \text{ GeV}$
Number of crystals in EMC cluster	$20 < N_{\text{crys}} < 48$	$12 < N_{\text{crys}} < 36$
LAT shower shape	$0.24 < LAT < 0.51$	$0.15 < LAT < 0.49$
a_{42} shower shape	$a_{42} < 0.07$	$a_{42} < 0.07$
Polar angle acceptance	$-0.31 < \cos \theta_{\gamma}^* < 0.6$	$-0.46 < \cos \theta_{\gamma}^* < 0.46$
2nd highest cluster energy (CMS)	$E_2^* < 0.2 \text{ GeV}$	$E_2^* < 0.14 \text{ GeV}$
Extra photon correlation	$\cos(\phi_2^* - \phi_1^*) > -0.95$	$\cos(\phi_2^* - \phi_1^*) > -0.95$
Extra EMC energy (Lab)	$E_{\text{extra}} < 0.1 \text{ GeV}$	$E_{\text{extra}} < 0.22 \text{ GeV}$
IFR veto	$\cos(\Delta\phi_{\text{NH}}^*) > -0.9$	$\cos(\Delta\phi_{\text{NH}}^*) > -0.95$
IFR fiducial	$\cos(6\phi_{\gamma}^*) < 0.96$...

Selection of high-quality photons, with tighter criteria for lower photon energies (increasing backgrounds)



Total Efficiency:

High Energy Region: 10-11%

Low Energy Region: 20%

Systematic Uncertainties - Higgs

$e^+e^- \rightarrow \gamma\gamma$ background (dominant effect)

varying the yield gives a ± 38 event uncertainty for $m_{A_0} = 0 \text{ GeV}/c^2$,
with a decreasing effect for larger masses.

varying the shape gives a ± 70 event uncertainty at $m_{A_0} = 7.4 \text{ GeV}/c^2$

Signal PDF

corrected using data vs. simulation comparison of $e^+e^- \rightarrow \gamma\gamma$ events,
taking half the correction as the systematic uncertainty

- The largest impact is at $m_{A_0} = 7.4 \text{ GeV}/c^2$, where the signal yield varies by ± 64 events

Signal Efficiency

trigger/event filter efficiency checked with $e^+e^- \rightarrow \gamma\gamma$ and $e^+e^- \gamma$ (0.4%)

Photon selection checked using $e^+e^- \rightarrow \mu\mu\gamma$, $\tau\tau\gamma$, and $\omega\gamma$ (2%)

Neutral reconstruction: 2%

