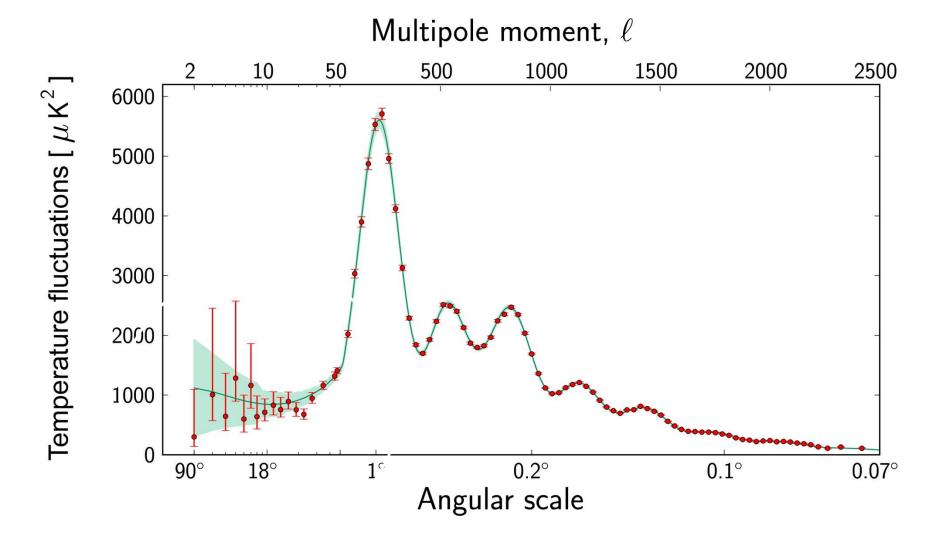
Challenges for theoretical cosmology

Andreas Albrecht

Fundamental Questions in Cosmology UC Davis May 2013

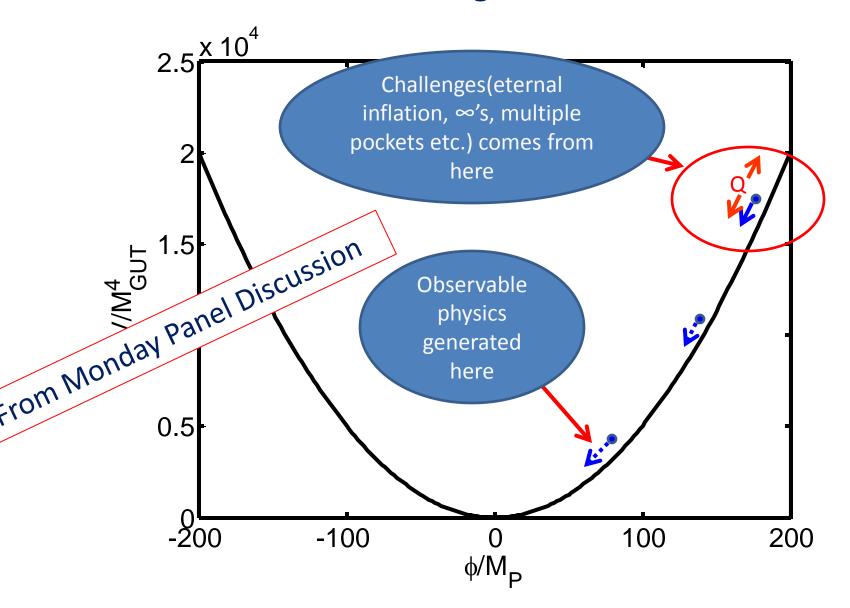


Challenges for Cosmic Inflation (eternal inflation)

"Anything that can happen will happen infinitely many times" (A. Guth)

- 2) Problems defining probabiliting probabili problem claiming generic predictions about state cannot claim "solution to cosmological
 - problems"
 - → Related to 2nd law, low S start
 - 4) **Yet**, Successful fits to data

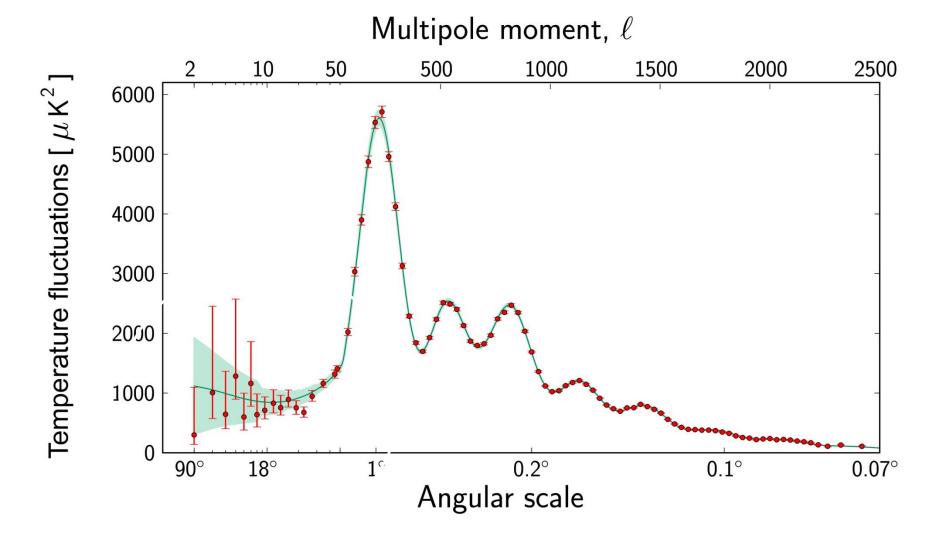
Slow rolling of inflaton



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Challenges: Answer these questions re your theories & beliefs:

- 1) Do you predict the observed state of the universe to be likely or natural? (And do you care?)
- 2) Do you treat infinities rigorously?
- 3) Do you require a probability tooth fairy?

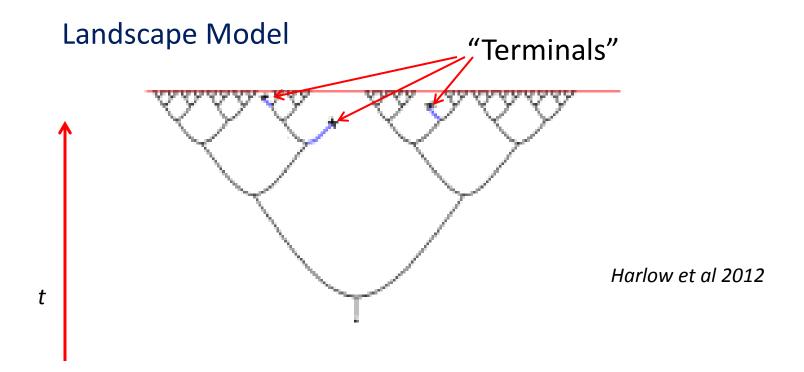
• Beware hidden assumptions about initial conditions (often related to 2^{nd} law: $\dot{S} > 0 \implies S$ initially small \implies starting in limited part of phase space)

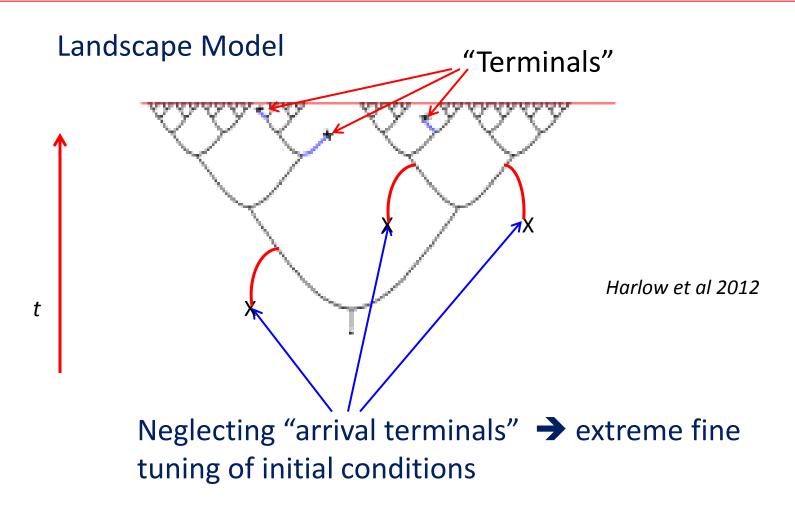
Gibbons & Turok
Carroll & Tam
Shiffren & Wald
Penrose

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Carroll & Tam
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(just as true of cyclic models)





In general: Need a quantitative theory for your starting point (inflation, cyclic, whatever) to make this claim.

Attempts I know to create this rigor have led to surprises.

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X	Υ
Volume of inflated	Probability for starting
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Entropy	Probability of starting a cyclic universe

X	Υ
Volume of inflated regions	Probability for starting inflation
Entropy	Probability of starting a cyclic universe
Number of observers (in my theory) who see a universe like ours	The infinitely many other observers who see something totally different

"Property X is infinite, so I don't need to worry about issue Y"

Need more rigor:

- Hernley, AA & Dray (2013) ←→ Guth toy model
- AA & Sorbo (2004)

Increasing the level of rigor usually reveals significant hidden assumptions that amount to tuning of initial conditions.

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Non-Quantum probabilities in a toy model:

$$U = A \otimes B \qquad A: \{|1\rangle^{A}, |2\rangle^{A}\} \qquad B: \{|1\rangle^{B}, |2\rangle^{B}\}$$

$$U: \{|11\rangle, |12\rangle, |21\rangle, |22\rangle\} \qquad |ij\rangle \equiv |i\rangle^{A} |j\rangle^{B}$$

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Possible Measurements Projection operators:

Measure A only:
$$\hat{P}_i^A = (|i\rangle^A \langle i|) \otimes \mathbf{1}^B = [|i1\rangle\langle i1| + |i2\rangle\langle i2|]$$

Measure *B* only:
$$\hat{P}_{i}^{B} = (|i\rangle^{B} \langle i|) \otimes \mathbf{1}^{A} = [|1i\rangle\langle 1i| + |2i\rangle\langle 2i|]$$

Measure entire *U*:
$$\hat{P}_{ii} \equiv |ij\rangle\langle ij|$$

BUT: It is impossible to construct a projection operator for the case where you do not know whether it is A or B that is being measured.

Non-Quan

$$U = A$$

Could Write

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$$|j\rangle^{\!\scriptscriptstyle B}$$

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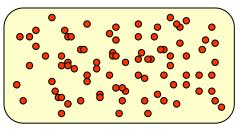
Classica

Pro abilities

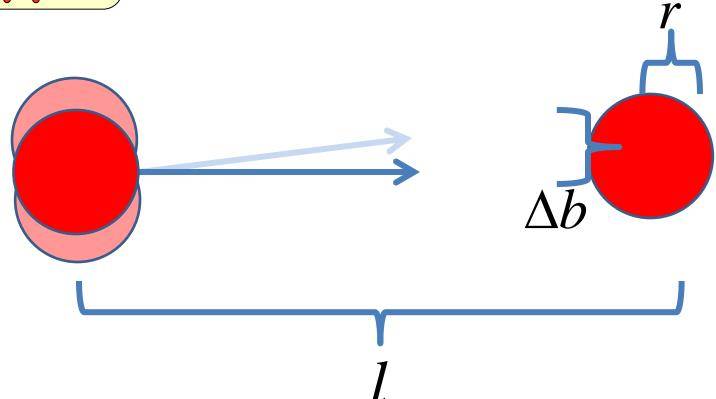
A, B

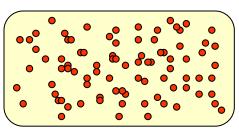
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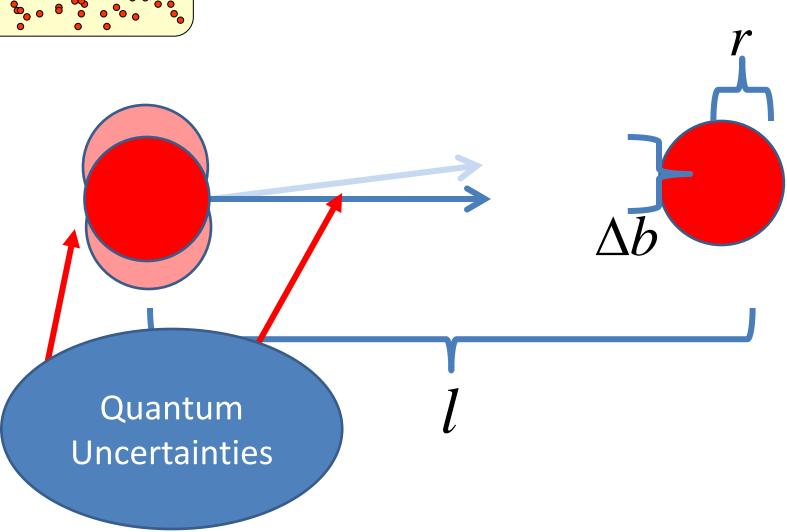


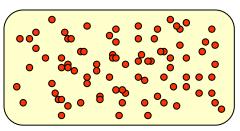
Quantum effects in a billiard gas



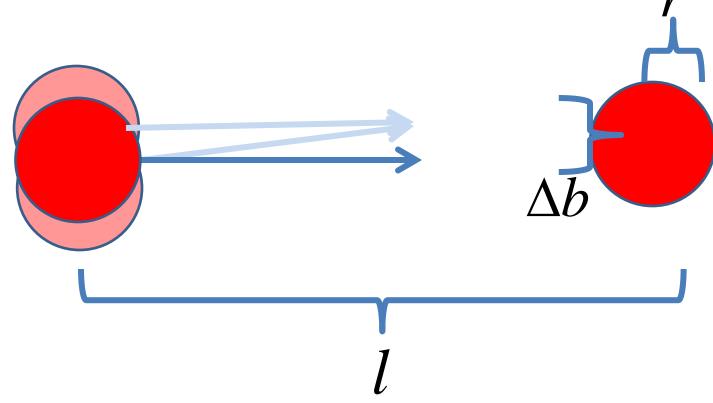


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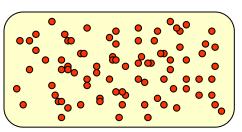


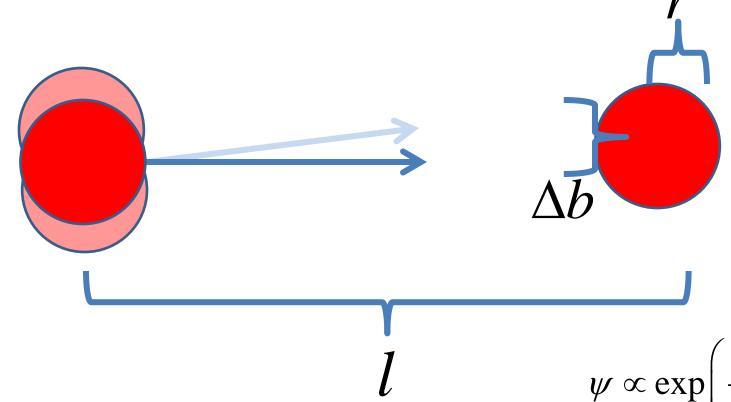


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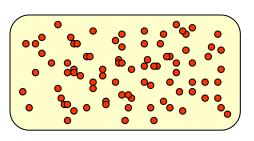


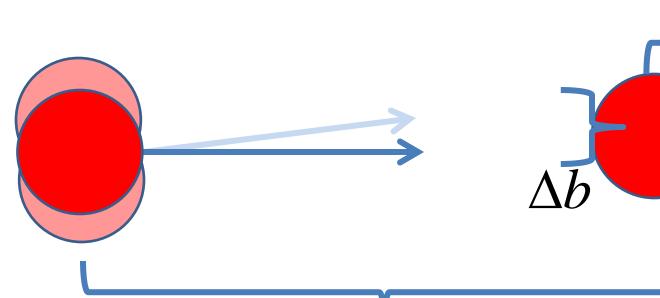
$$\Delta b = \delta x_{\perp} + \frac{\delta p_{\perp}}{m} \Delta t$$





 $\Delta b = \delta x_{\perp} + \frac{\delta p_{\perp}}{m} \Delta t = \sqrt{2} \left(a + \frac{\hbar}{2a} \frac{l}{m\overline{v}} \right)$



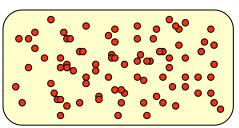


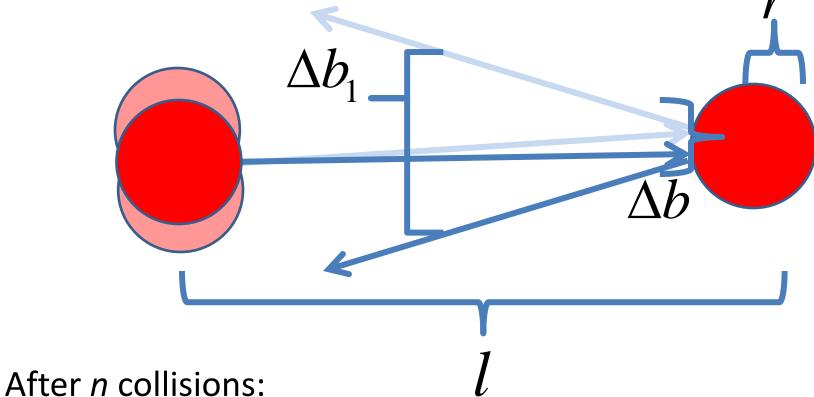
$$\Delta b = \delta x_{\perp} + \frac{\delta p_{\perp}}{m} \Delta t = \sqrt{2} \left(a + \frac{\hbar}{2a} \frac{l}{m\overline{v}} \right) \qquad \psi \propto \exp\left(\frac{-x^2}{2a^2} \right)$$

$$\xrightarrow{\min} 2^{3/2} \left(\frac{\hbar l}{2m\overline{v}} \right) \equiv \sqrt{l \mathcal{A}_{dB} / 2}$$

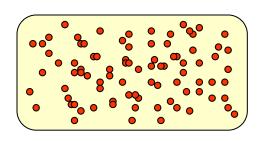
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38

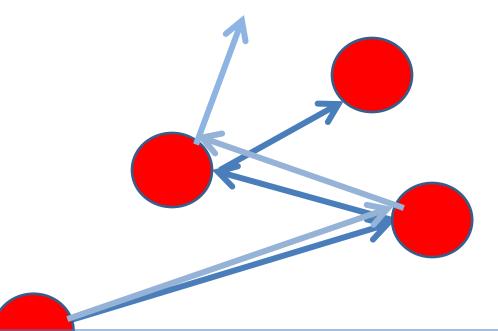
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$$\Delta b_n = \Delta b \left(1 + 2l / r \right)^n$$





 ${\it n_{\scriptscriptstyle O}}$ is the number of collisions so that

$$\Delta b_{n_O} = r$$

(full quantum uncertainty as to which is the next collision)

$$n_Q = -\frac{\log\left(\frac{\Delta b}{r}\right)}{\log\left(1 + \frac{2l}{r}\right)}$$
vis 5/22/13

	r	l	m	\overline{v}	λ_{dB}	Δb	n_{Q}
Air							,
Water							
Billiards							
Bumper Car							

	r	l	m	\overline{v}	$\frac{\lambda_{dB}}{2}$	Δb	n_{Q}
Air							
Water							
Billiards							
Bumper Car	1	2	150	0.5	1.4×10^{-36}	3.4×10^{-18}	25



	r	l	m	\overline{v}	$\frac{\lambda_{dB}}{\partial B}$	Δb	n_{Q}
Air							Q
Water							
Billiards	0.029	1	0.16	1	6.6×10^{-34}	5.1×10^{-17}	8
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Albrecht @ FQC, Davis 5/22/13

	r	l	m	\overline{v}	λ_{dB}	Δb	$ n_{Q} $
Air							
Water	3.0×10^{-10}	5.4×10^{-10}	3×10^{-26}	460	7.6×10^{-12}	1.3×10^{-10}	0.6
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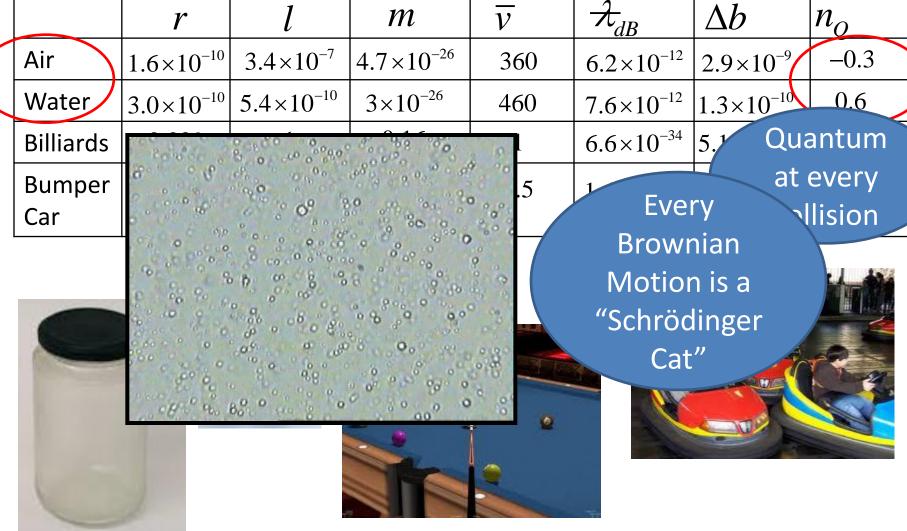


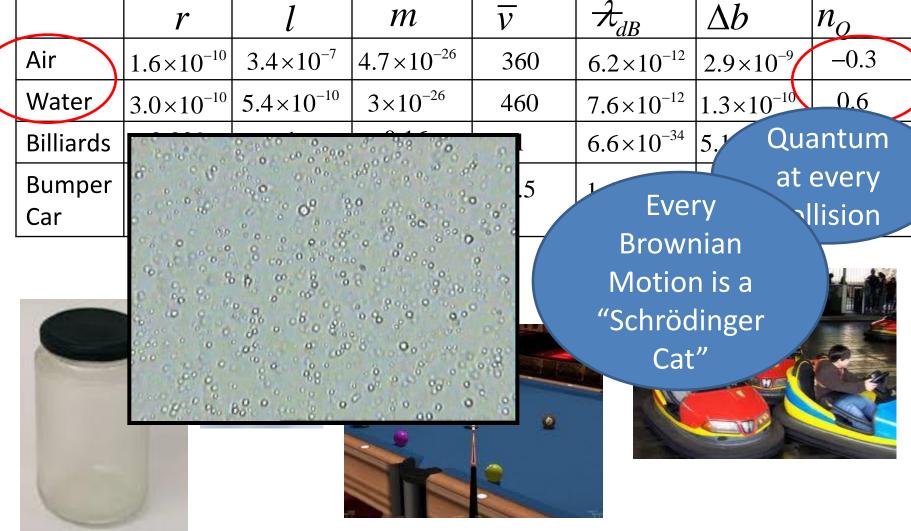




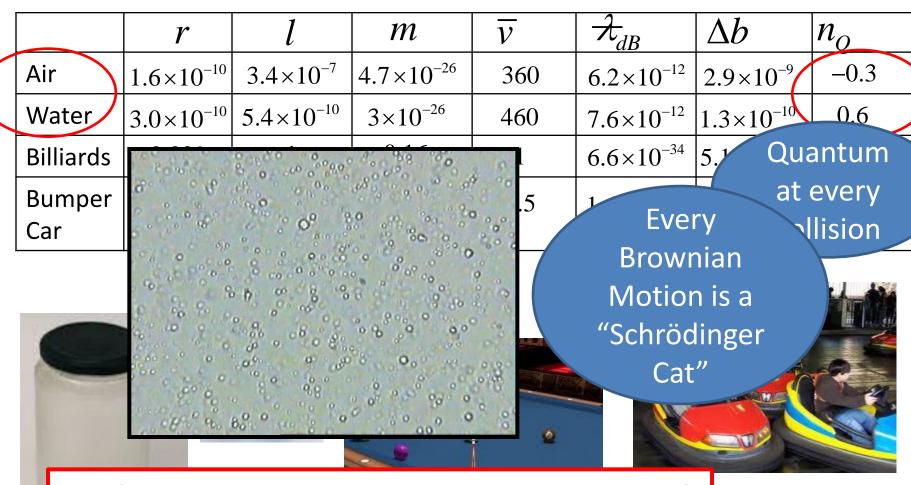


Albrecht @ FQC, Davis 5/22/13





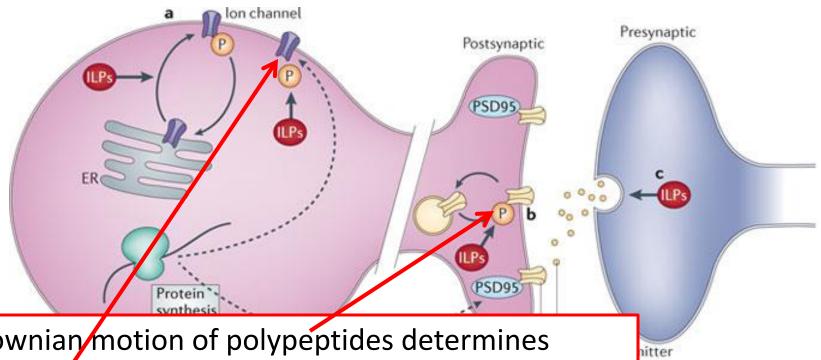
(all units MKS)



(independent of "interpretation")

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An important role for Brownian motion: Uncertainty in neuron transmission times



Brownian motion of polypeptides determines exactly how many of them are blocking ion channels in neurons at any given time. This is believed to be the dominant source of neuron transmission time uncertainties $\delta t_n \approx 1 ms$

re Reviews | Neuroscience

$$\delta t_f = \delta t_n \times \left(\frac{v_h}{v_h + v_f}\right)$$

$$\delta t_t = \sqrt{2}\delta t_f$$

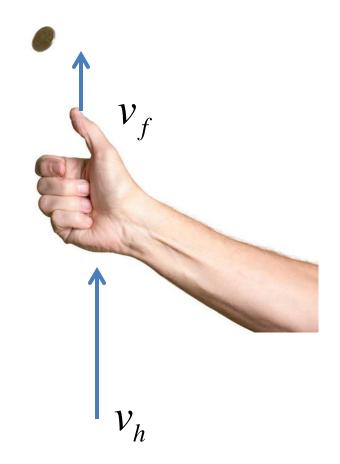
$$f = \frac{4v_f}{\pi d}$$

$$\delta N = f \delta t_t = 0.5$$



$$\delta t_n \approx 1 ms$$
 $v_h = v_f = 5 m / s$

$$d = 0.01m$$



Coin diameter = d

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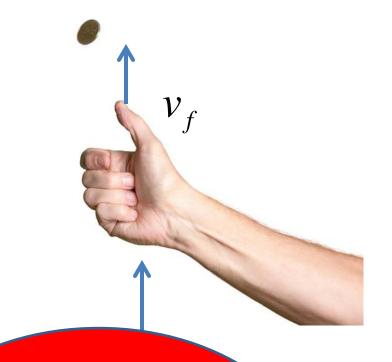
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Using:

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50-50 coin flip probabilities are a derivable quantum result

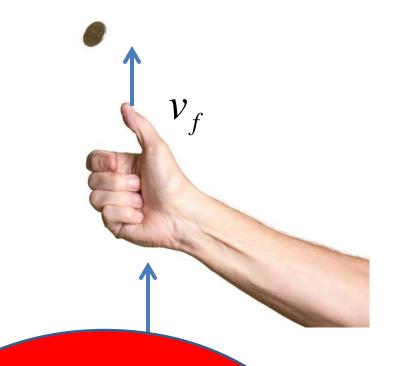
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Usin Without reference to "principle of indifference" etc.



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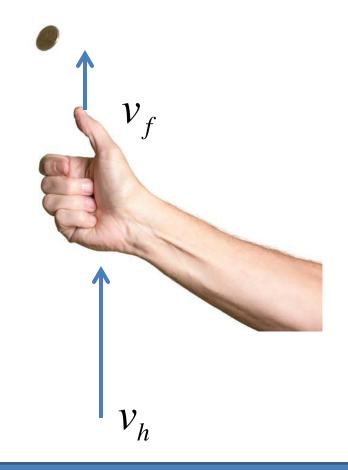
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NB: Coin flip is "at the margin" of classical vs quantum control: Increasing d or deceasing v_h can reduce δN substantially

Albrecht @ FQC, D

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- 2) Only then can we claim to resolve the famous cosmological puzzles (Monopoles already OK).
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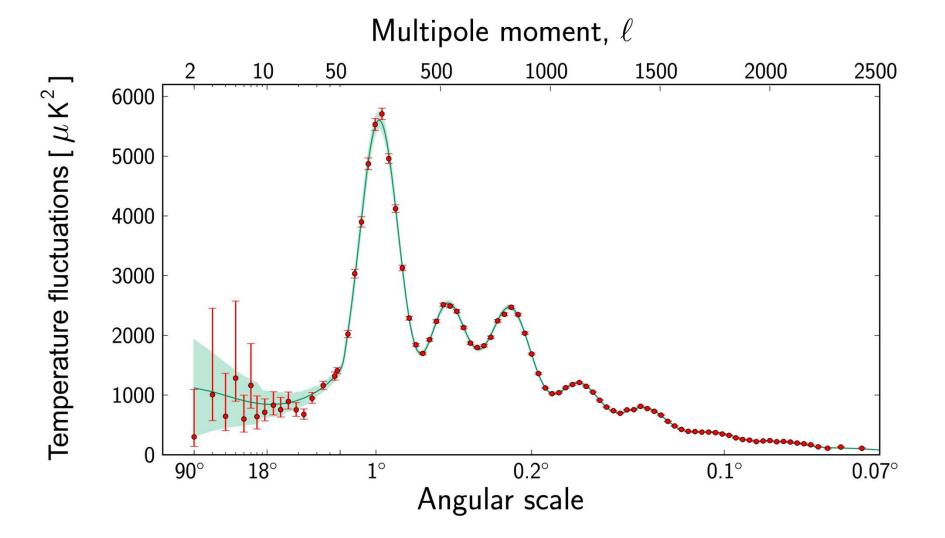
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- 2) Do you treat infinities rigorously?
- 3) Do you require a probability tooth fairy?



Challenges: Answer these questions re your theories & beliefs:

- 1) Do you predict the observed state of the universe to be likely or natural? (And do you care?)
- 2) Do you treat infinities rigorously?

Finite

3) Do you require a probability tooth fairy

No